

Positron-Electron Pair Creation with Intense Lasers and Applications

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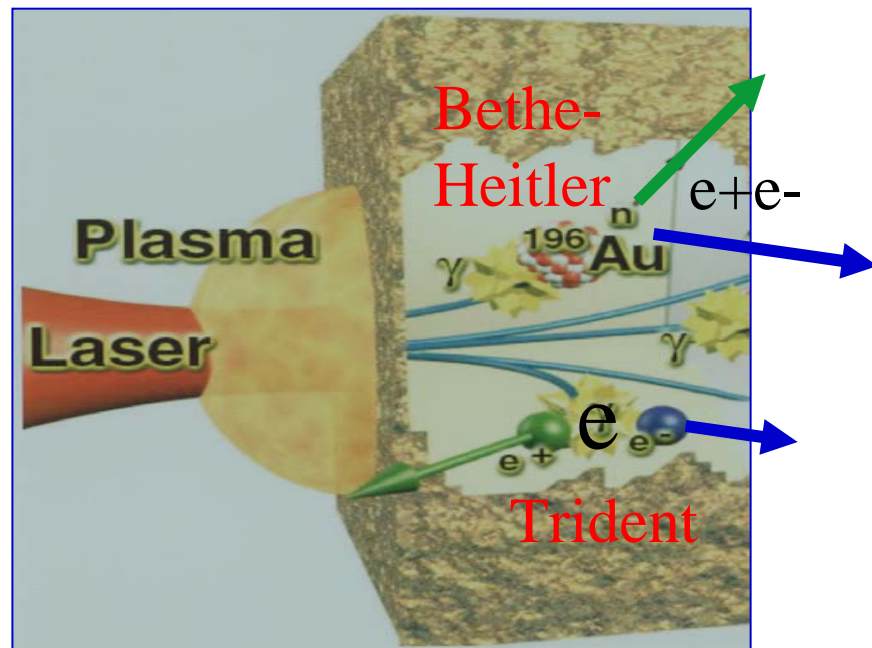
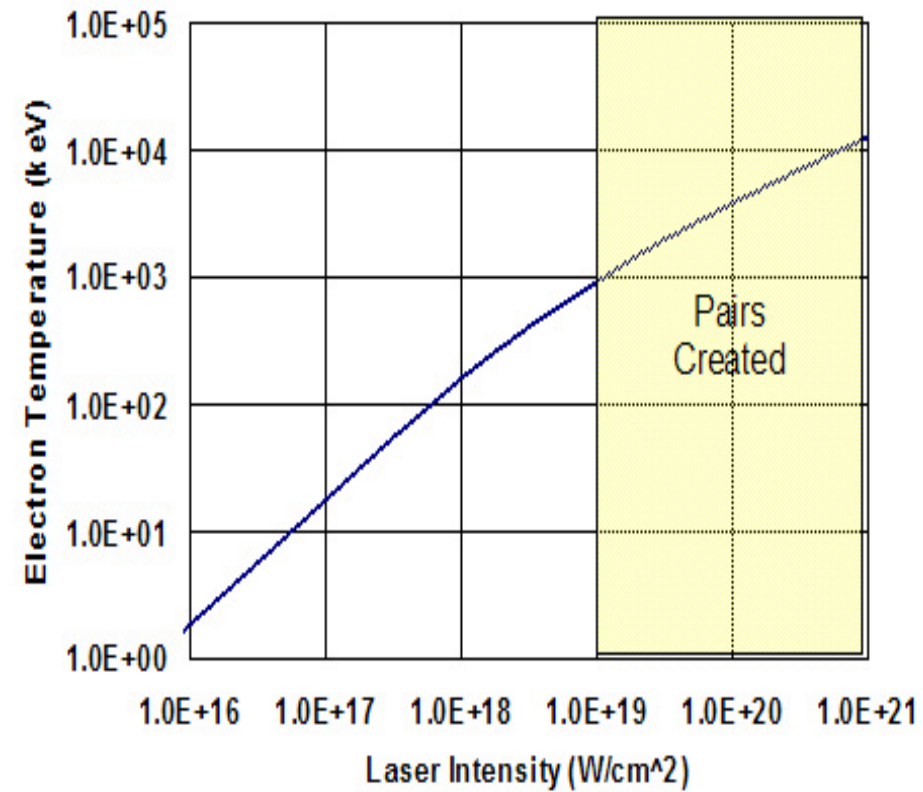
Antimatter has many potentially revolutionary applications. But current technique of creating positrons with accelerators and trapping them with electrostatic traps is very inefficient and can only produce low density e^+ (Surko and Greaves 2004):

Maximum Beam Intensity $< 10^{10} e^+/s$

Maximum Density $< 10^{14} e^+/cm^3$

Ultra-intense lasers opens up exciting new approach to produce high intensity, high density e^+e^- pairs with high efficiency

It is very desirable to reach pair density $> 10^{18}/cm^3$
for applications



Sample Laser Numbers

$$1 \text{ PW} = 1 \text{ kJ} / 1 \text{ ps}$$

$$1 \text{ PW} / (30 \text{ } \mu\text{m})^2 = 10^{20} \text{ W/cm}^2$$

$$10^{20} \text{ W/cm}^2/c \sim 3 \cdot 10^{16} \text{ erg/cm}^3 \sim 2 \cdot 10^{22} \text{ e}^+\text{e}^- / \text{cm}^3$$

$$\text{Solid Au ion density} \sim 6 \cdot 10^{22} / \text{cm}^3$$

$$n_+/n_e \sim 4 \cdot 10^{-3}$$

In reality, the max. achievable pair density is

probably around $10^{19} - 10^{20} \text{ cm}^{-3}$

PAIR PRODUCTION BY SUPERHERMALS ON HIGH-Z TARGET:

$$dN_{+}/dt = \underbrace{(dN_{+}/dt)_{eion}}_{\text{trident 1}} + \underbrace{(dN_{+}/dt)_{\gamma ion}}_2 + \underbrace{(dN_{+}/dt)_{\gamma\gamma}}_{\text{Bethe-Heitler 3}}$$

for thin ($\ll 20 \mu\text{m}$) laser targets. Hence

$$dN_{+}/dt = (N_{+} + N_{-}) \langle N_{ion} (f(\gamma) v \sigma_{eion}) \rangle$$

$f(\gamma)$ is normalized superthermal distribution function and

$$\sigma_{eion} \sim 1.4 \times 10^{20} \text{ cm}^2 Z^2 (\ln \gamma)^3 \quad \text{for } \gamma \gg 1$$

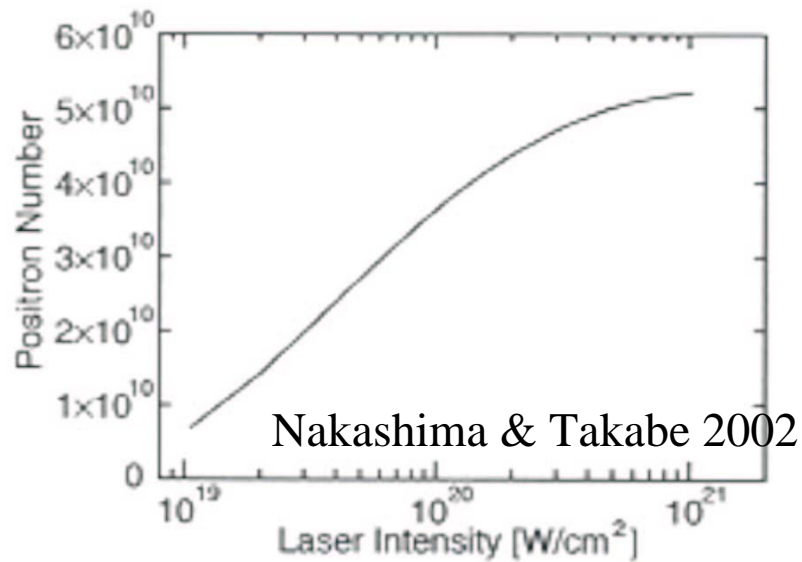
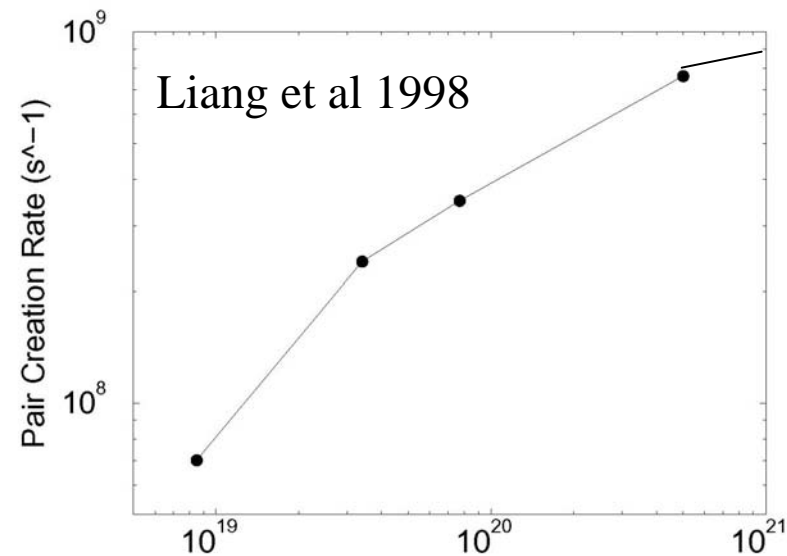
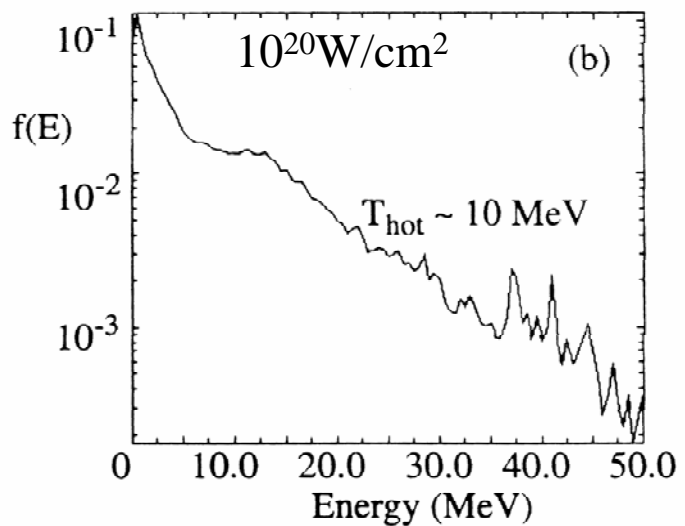
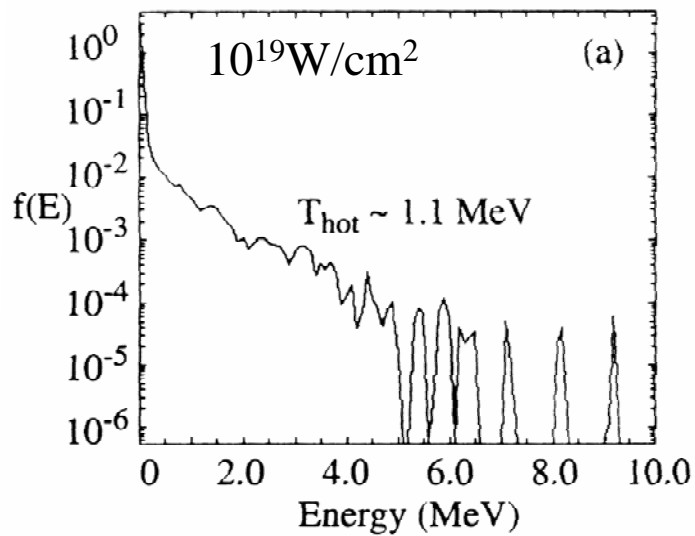
is trident pair production cross section ($e+ion \rightarrow e+ion+\gamma\gamma$):

Solving above equation:

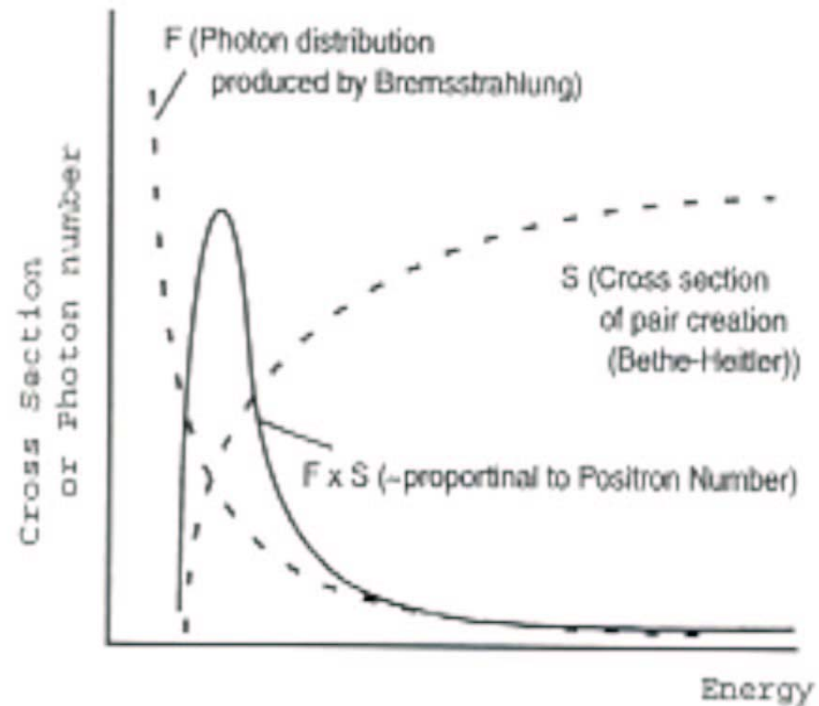
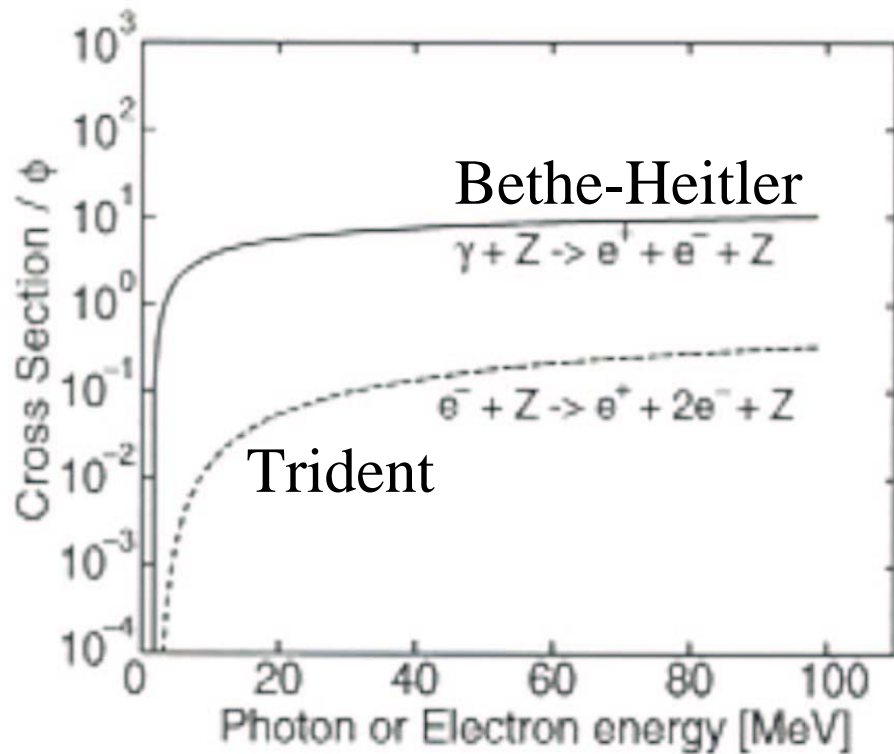
$$N_{+} = Z N_{ion} \{ \exp(\Gamma t) - 1 \} / 2 \sim Z N_{ion} \Gamma t / 2 \quad \text{for } \Gamma t \ll 1$$

$$\rightarrow N_{+}/N_e \sim \Gamma t / 2 \sim 2 \times 10^{-5} \quad \text{for } t \sim 1 \text{ ps, } I = 10^{20} \text{ Wcm}^{-2}$$

$$\text{For Au: } N_{+} \sim 10^{20} \text{ cm}^{-3}$$

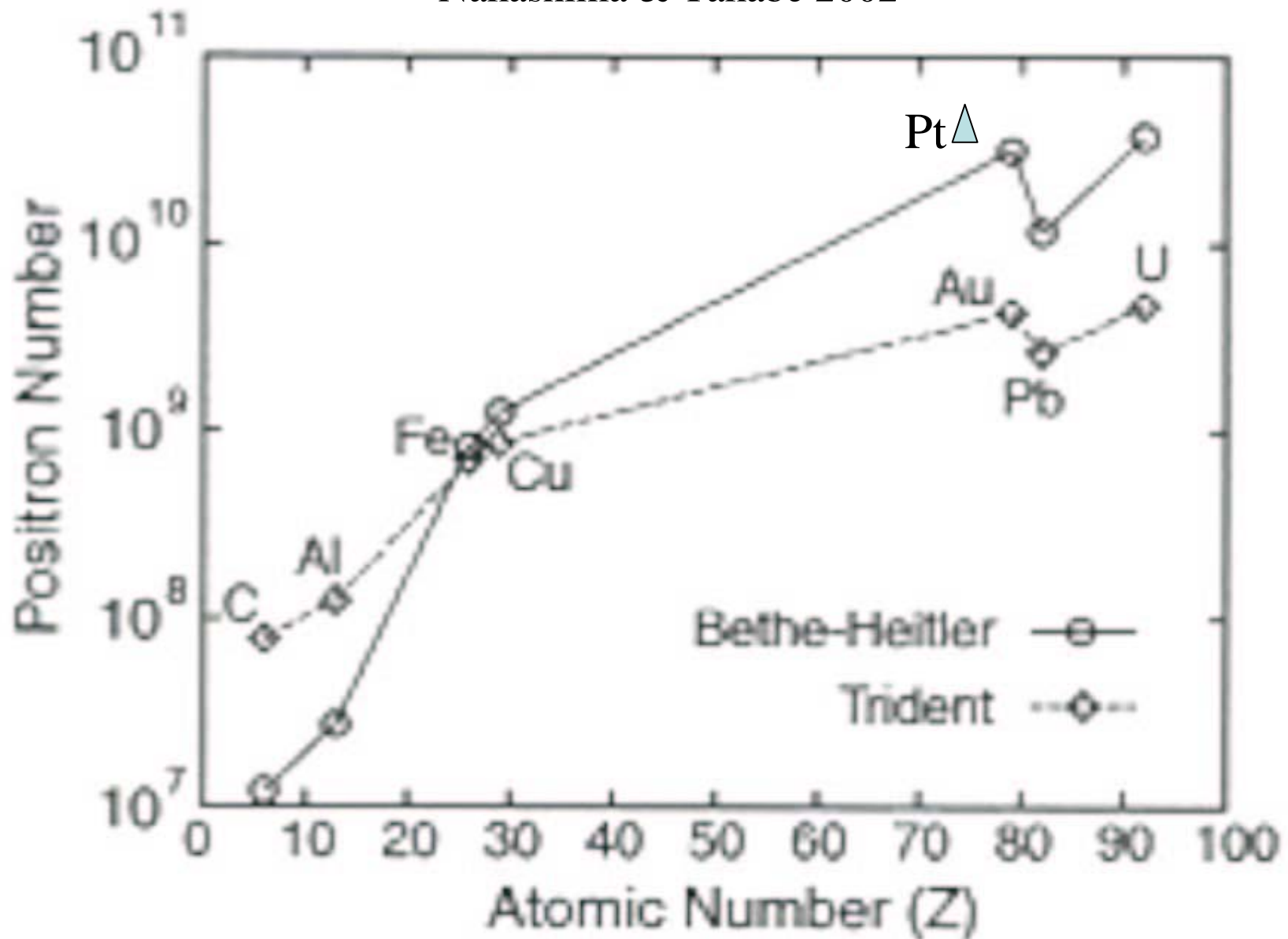


Pair Creation Rate Rises Rapidly with Laser Intensity to $\sim 10^{20} \text{ Wcm}^{-2}$,
but levels off after 10^{21} Wcm^{-2} .

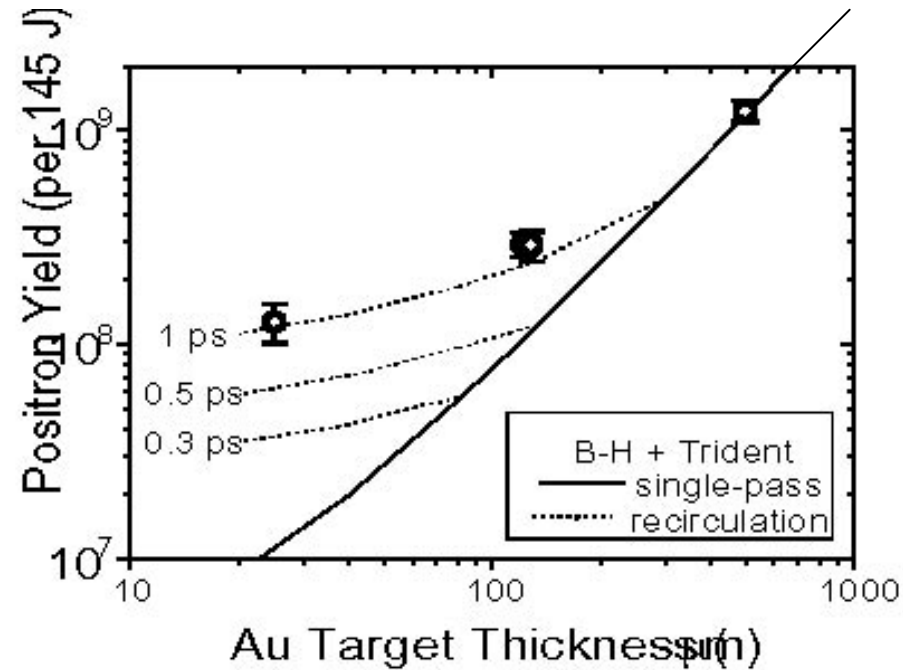
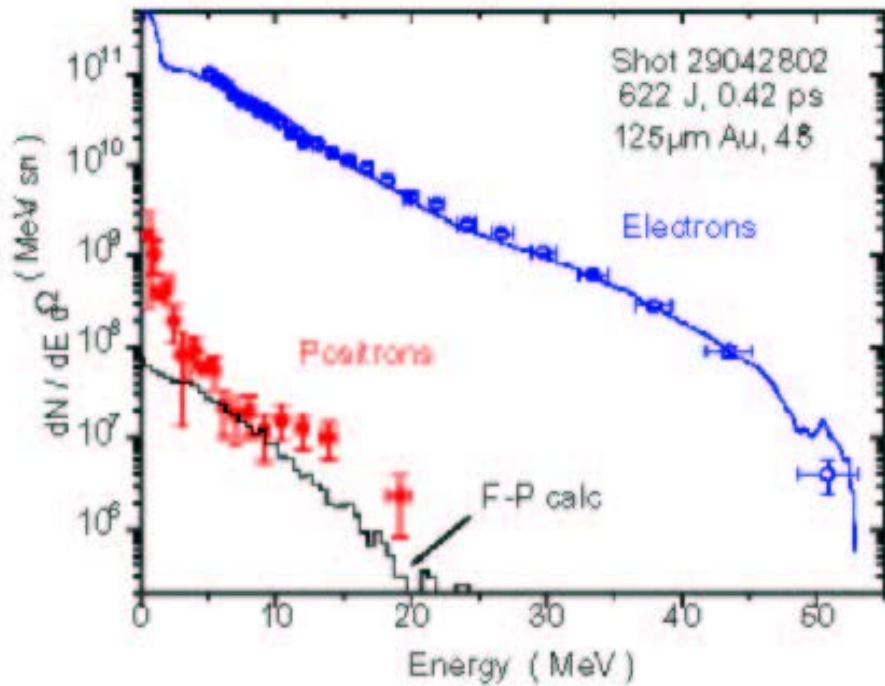


Nakashima & Takabe 2002

Bethe-Heitler pair-production has larger cross-section than Trident, but it depends on photon density, Z and optical depth of the high- Z target



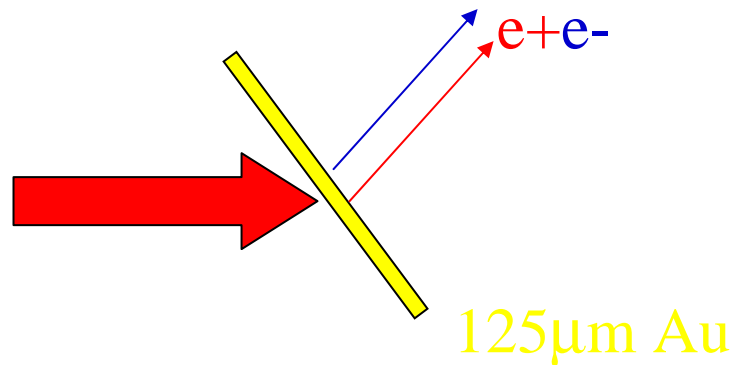
LLNL PW laser experiments first demonstrated e+e- production with Au foils. But the flux was low due to off-axis measurements

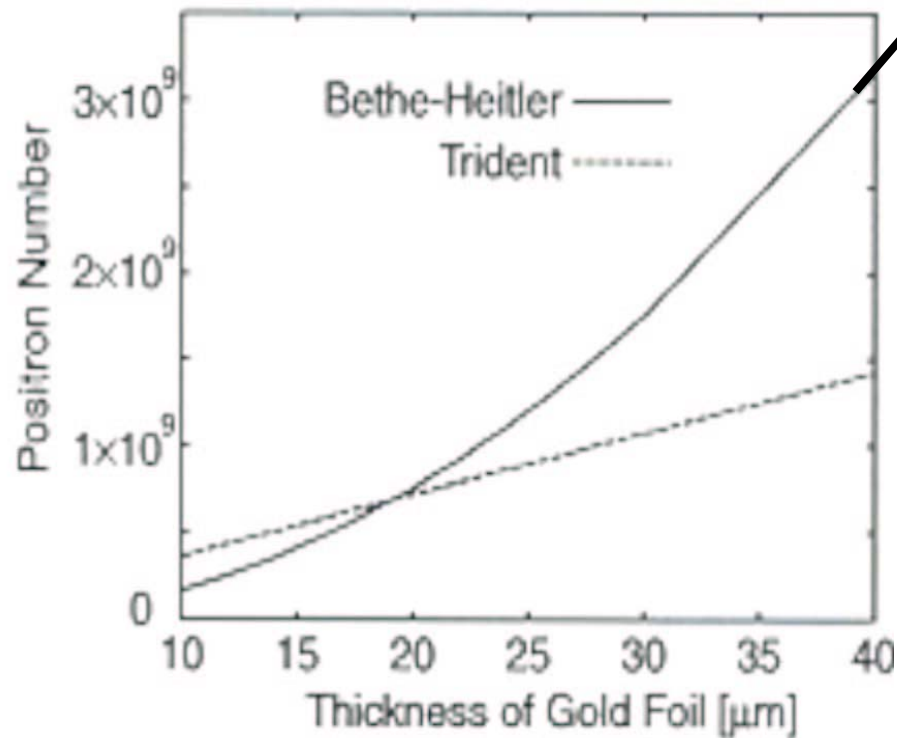


Cowan et al 2002

$2.10^{20} \text{W.cm}^{-2}$

0.42 ps



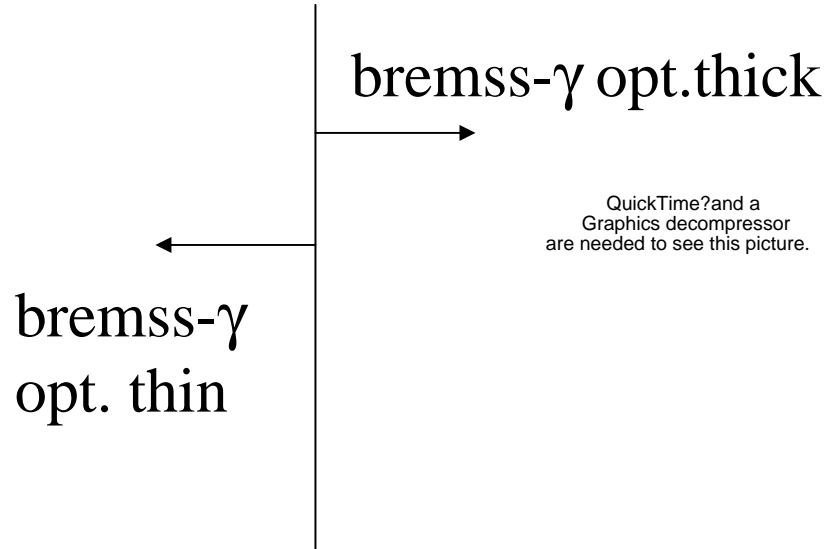


$$I=10^{20}\text{Wcm}^{-2}$$

Nakashima & Takabe 2002

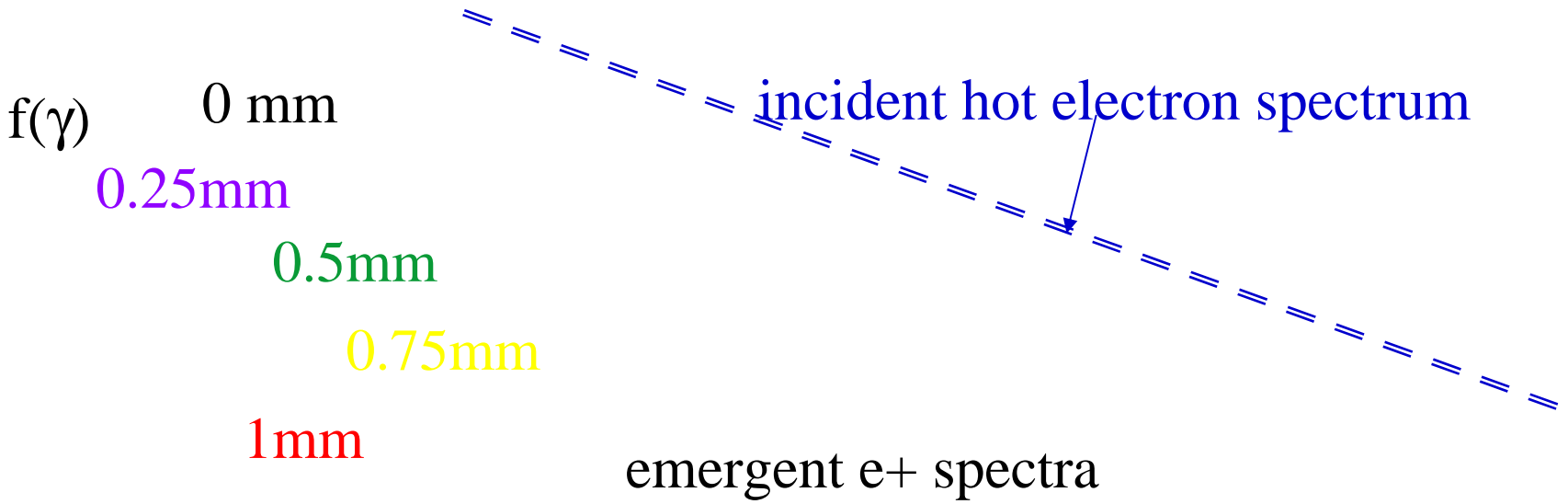
Trident process dominates for thin targets.
But Bethe-Heitler dominates for thick targets.
How high can the e^+ yield go if we use very thick targets?

Positron yield per incident hot electron increases with Au thickness:
transition from quadratic to linear occurs around 2-3 mm.



thickness (mm)

But emergent positrons are attenuated by cold absorption inside target :
low energy e+ are absorbed by thick target due to ionization losses



QuickTime?and a Graphics decompressor are needed to see this picture.

γ

Allowing for cold attenuation , emergent e+ yield per incident hot electron peaks at ~ 2.6 % where Au thickness ~ 5 - 6 mm

QuickTime?and a
Graphics decompressor
are needed to see this picture.

thickness (mm)

Assuming that the conversion of laser energy to hot electrons is
~ 30 % and the hot electron temperature is ~ 5 -10 MeV, the
above analysis suggests that the optimal positron yield is

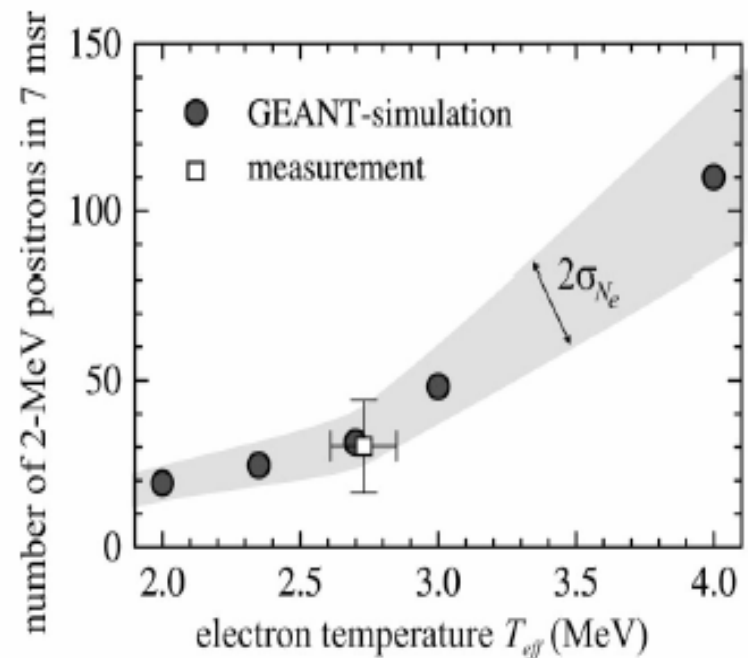
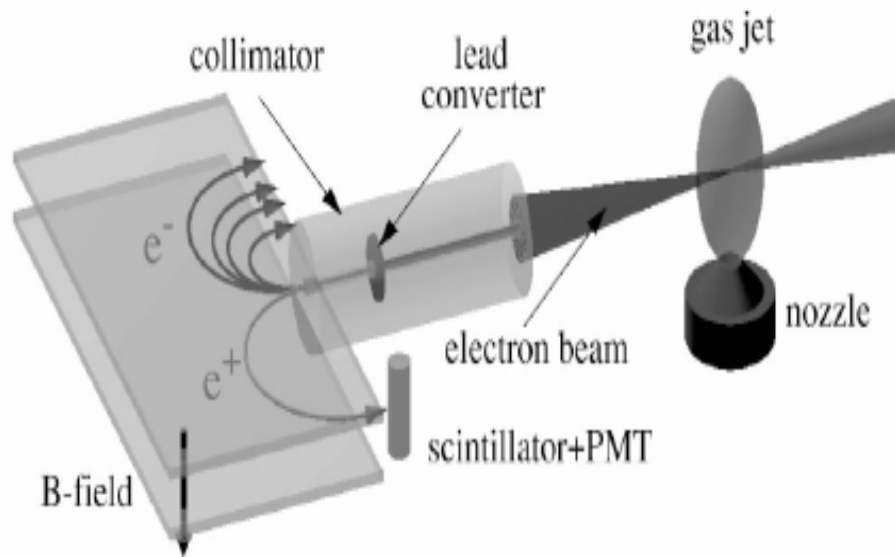
~ few x 10^{12} e+ per kJ of laser energy
when the Au target ~ 5 - 6 mm

The in-situ e+ density should reach > $10^{17}/\text{cm}^3$

For a 1 mm target, the e+ yield is ~ 10^{12} per kJ of laser energy.

These results are roughly consistent with the latest Titan
experiments (*Chen et al* PRL 2008 submitted and talk at the
upcoming DPP November 2008 meeting in Dallas)

Above direct irradiation of solid high-Z solid target requires PW-class lasers (“1-step” process). Alternative approach demonstrated by Gahn et al (2000, 2002) uses smaller lasers to first accelerate an MeV e-beam in a gas target (LWFA or bubble accelerator). The MeV electrons then create pairs by hitting a second high-Z converter. This “2-step” process produces lower e⁺ yield per J of laser energy, and lower e⁺ density, but it has higher rep rate so it can produce a semi-continuous e⁺ beam.



What are the most exciting applications of a dense e+ source?

High density e+ source makes it feasible to create a BEC of Ps at cryogenic temperature

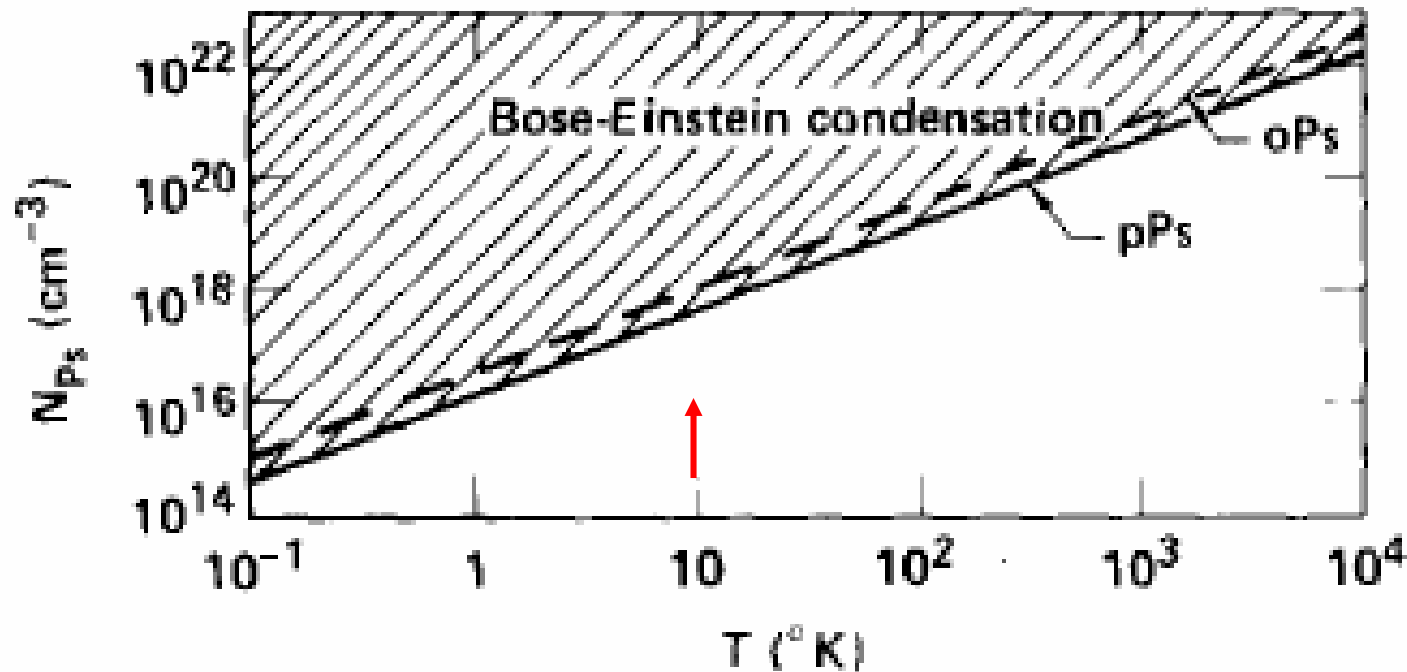
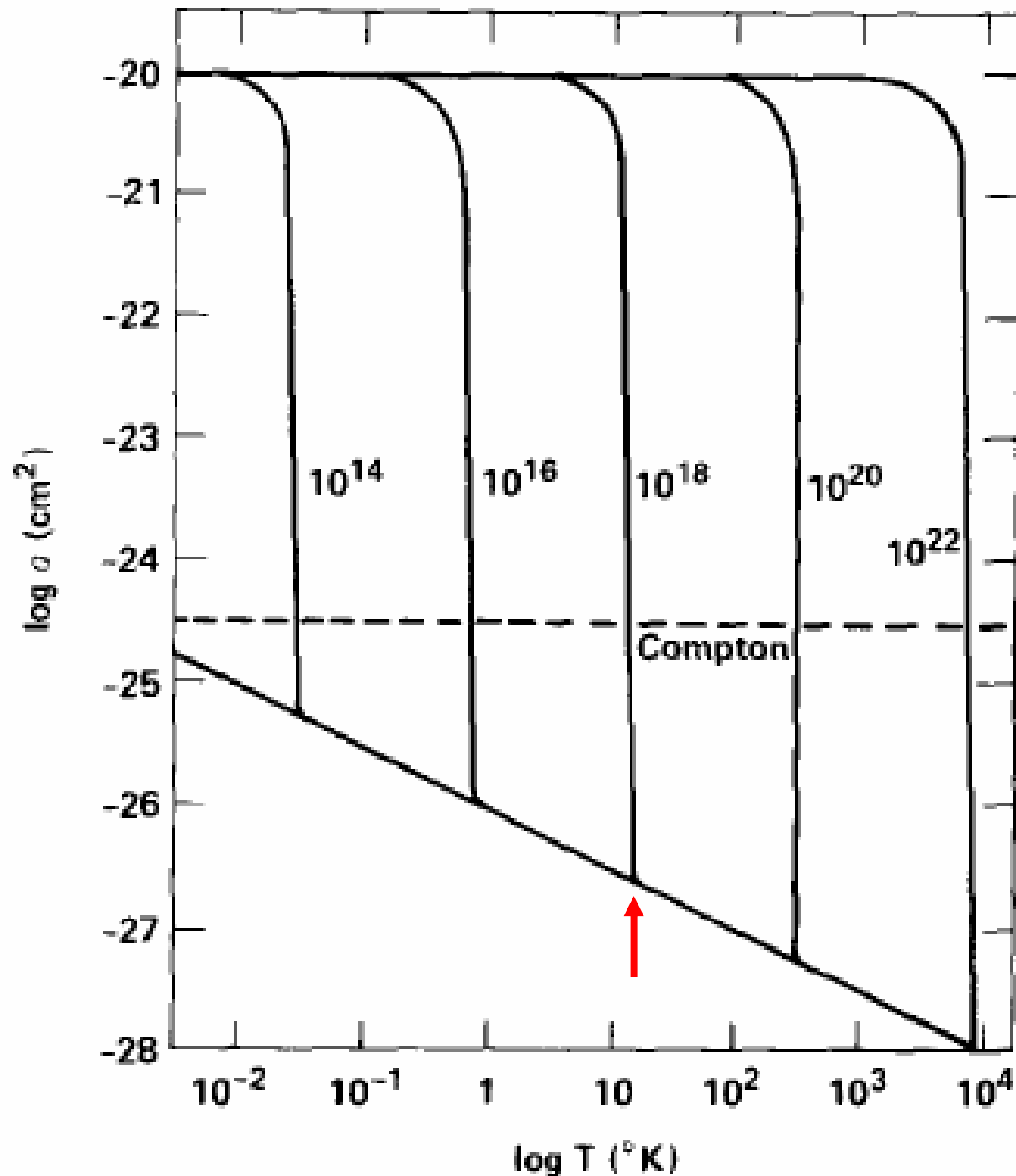
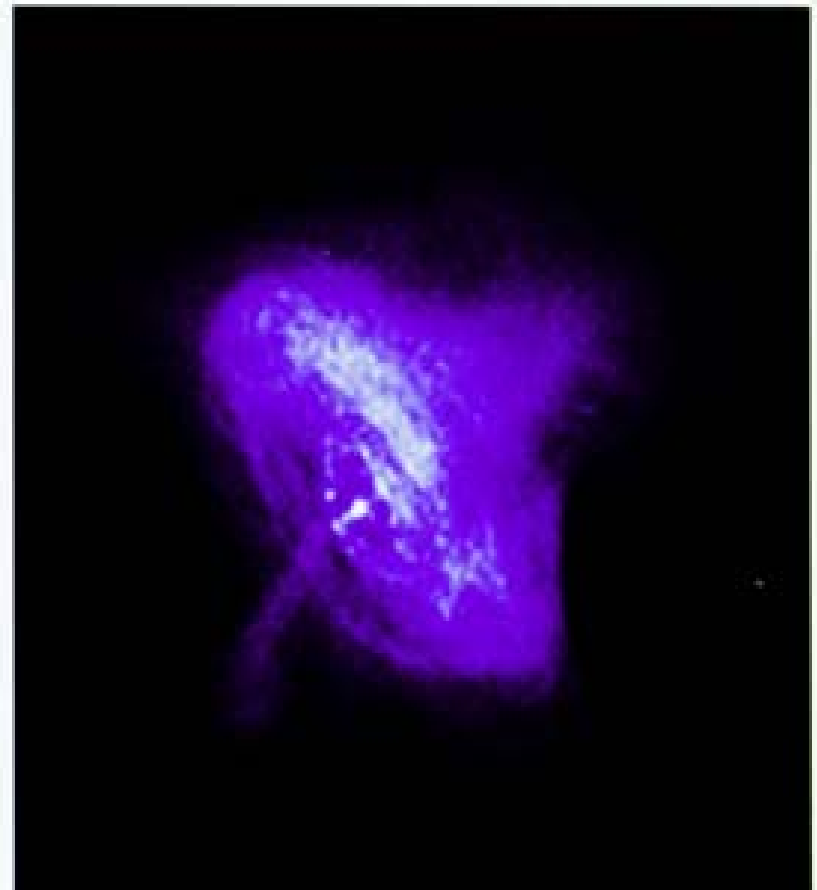
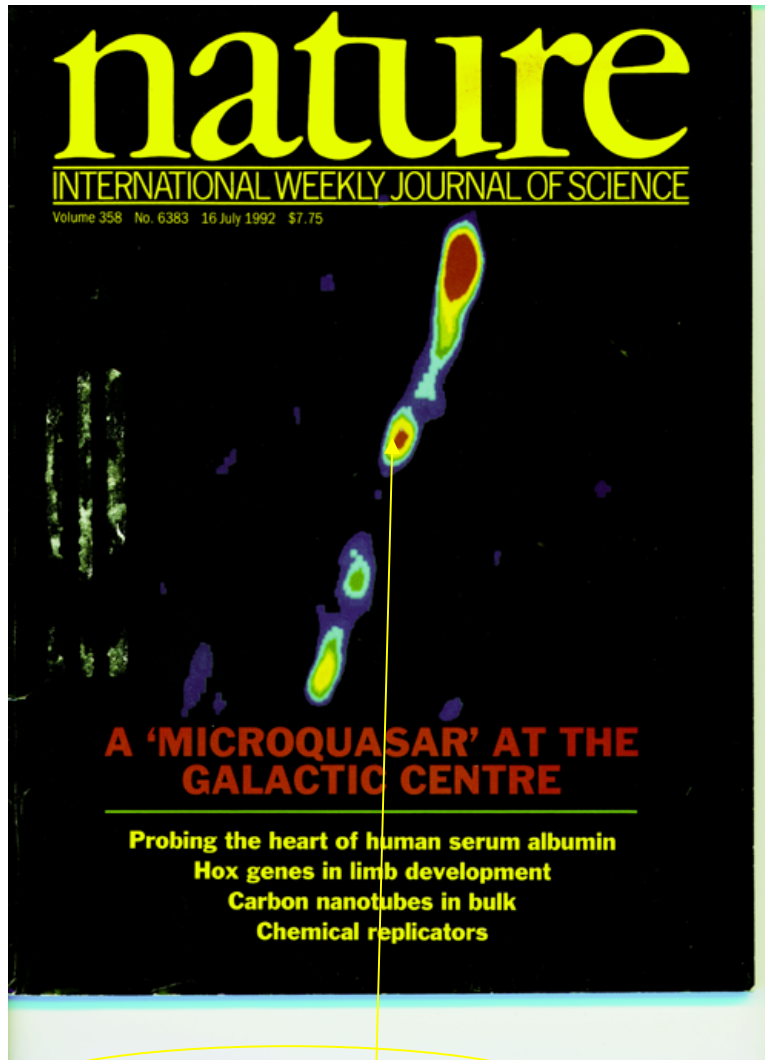


Fig. 4. Solid curve denotes the Bose-Einstein critical density as a function of temperature for pPs. Dashed curve denotes that of oPs. (See also ref. [7]). *(from Liang and Dermer 1988).*

A Ps column density of 10^{21} cm^{-2} could in principle achieve a gain-length of 10 for gamma-ray amplification via stimulated annihilation radiation (GRASAR).
(from Liang and Dermer 1988)



relativistic e^+e^- plasmas are ubiquitous in the universe

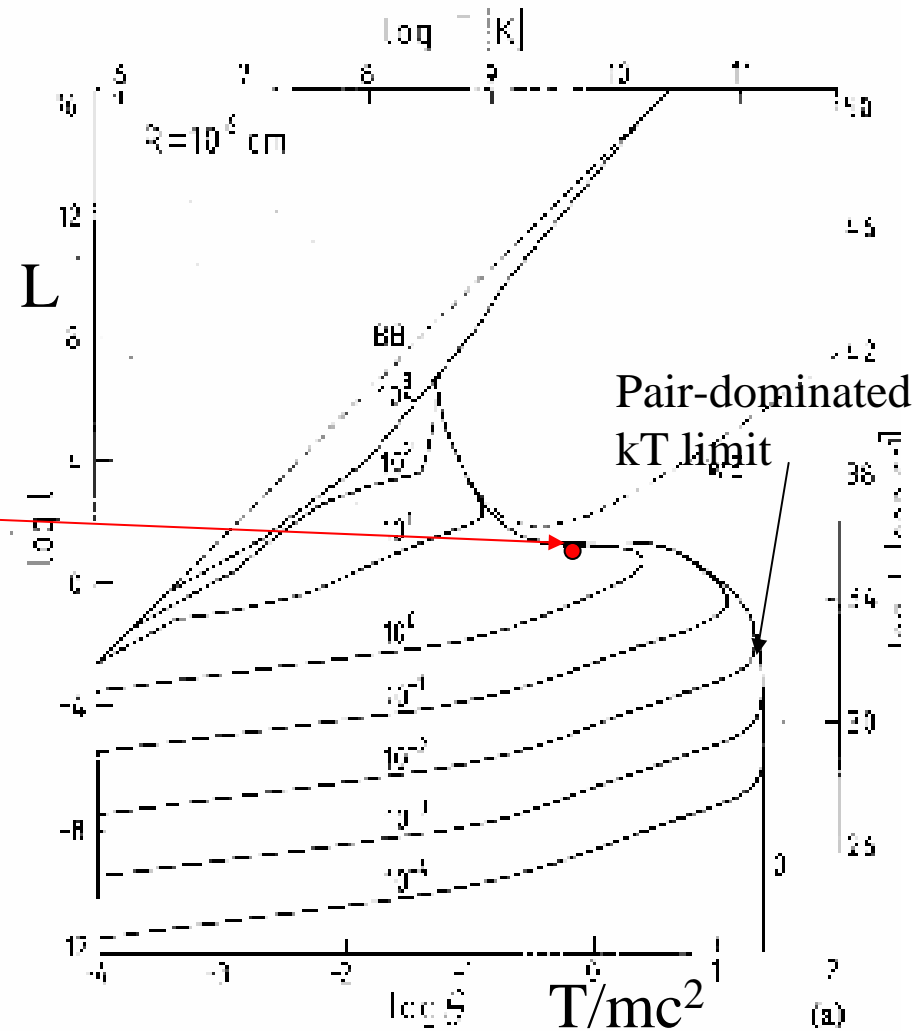
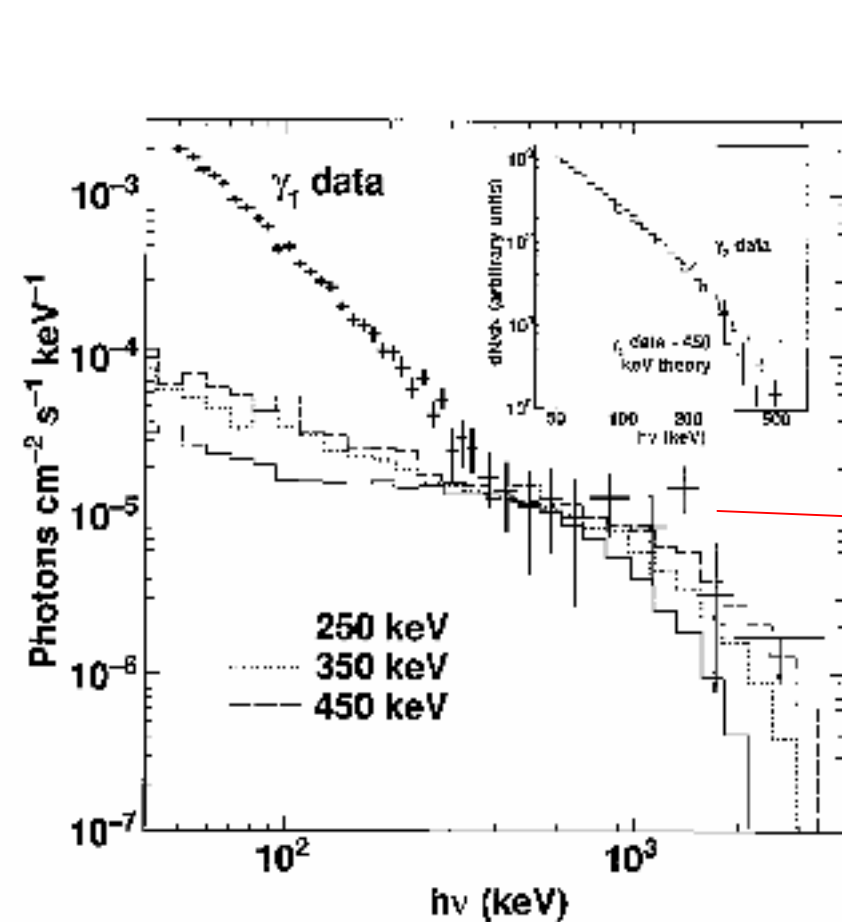


Nonthermal TeV pairs

Thermal MeV pairs

Laser-produced pair plasmas can be used to study lab astrophysics

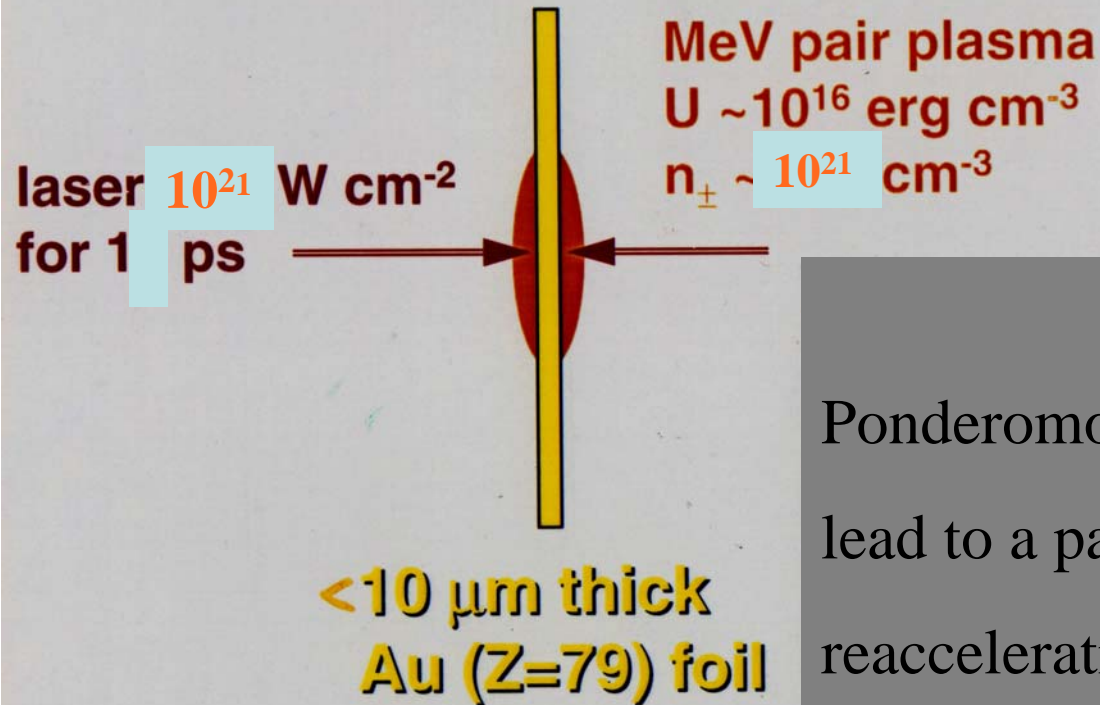
The Black Hole gamma-ray-bump can be interpreted as emissions from a pair-dominated MeV plasma with $n_+ \sim 10^{17} \text{ cm}^{-3}$



Can laser-produced pair plasmas probe the pair-dominated temperature limit?

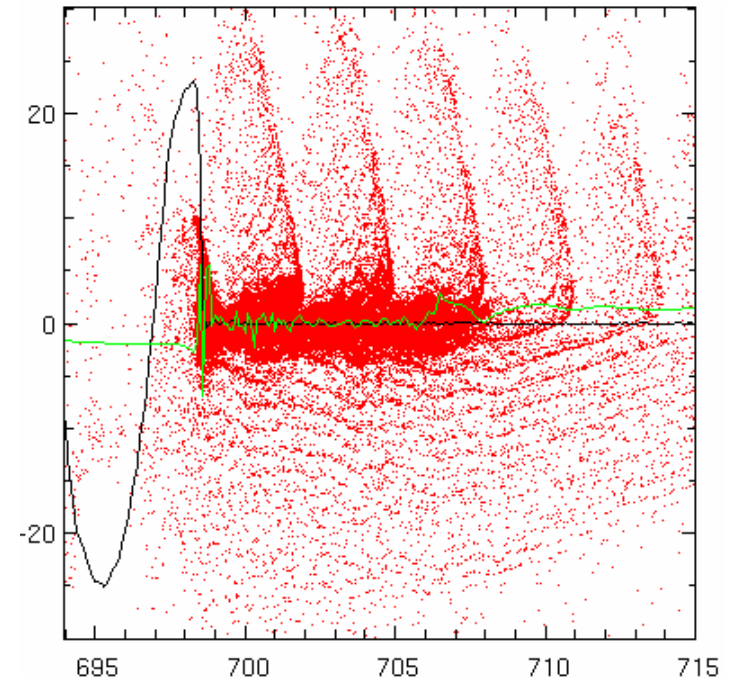
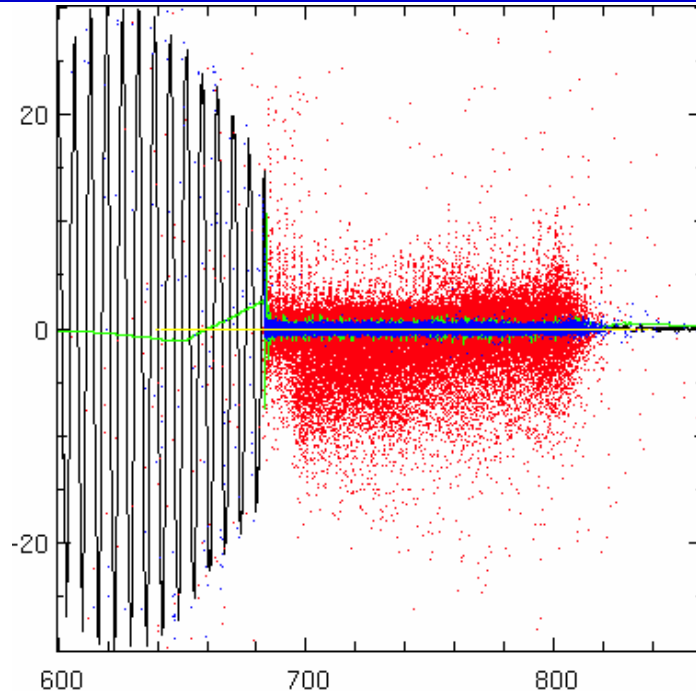
Another concept: two-sided irradiation may create more pairs, due to hotter electrons and longer confinement time

Creating a Pair Fireball

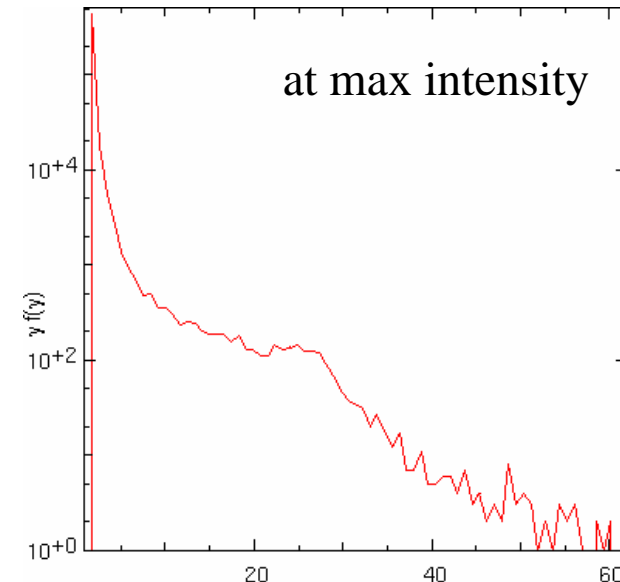
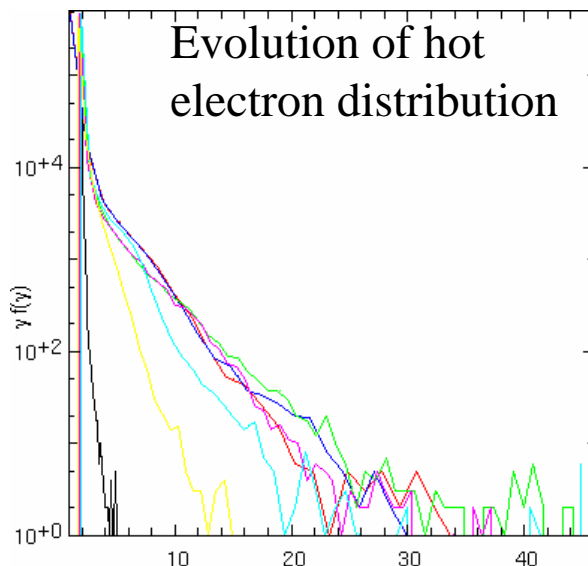


Ponderomotive forces can lead to a pair cascade by reaccelerating the primary pairs in the foil

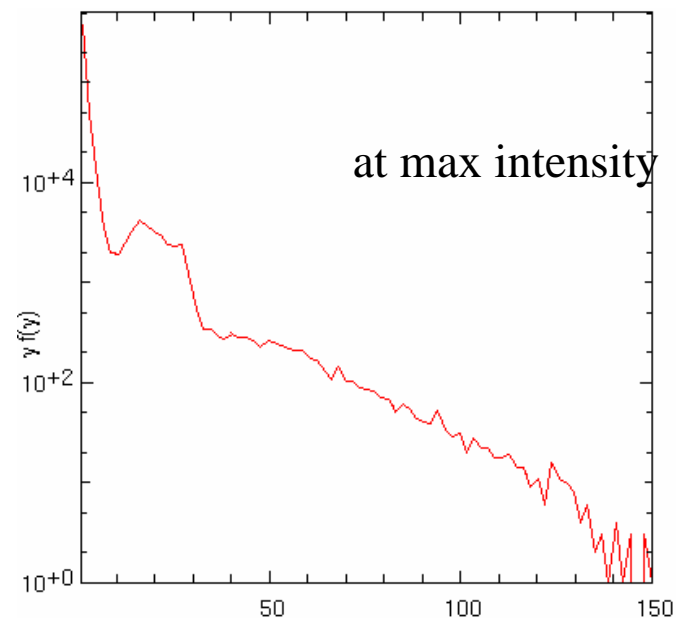
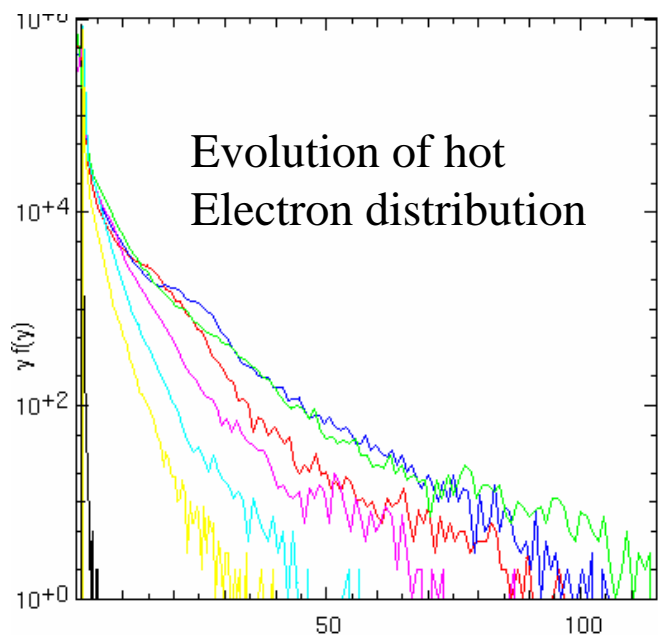
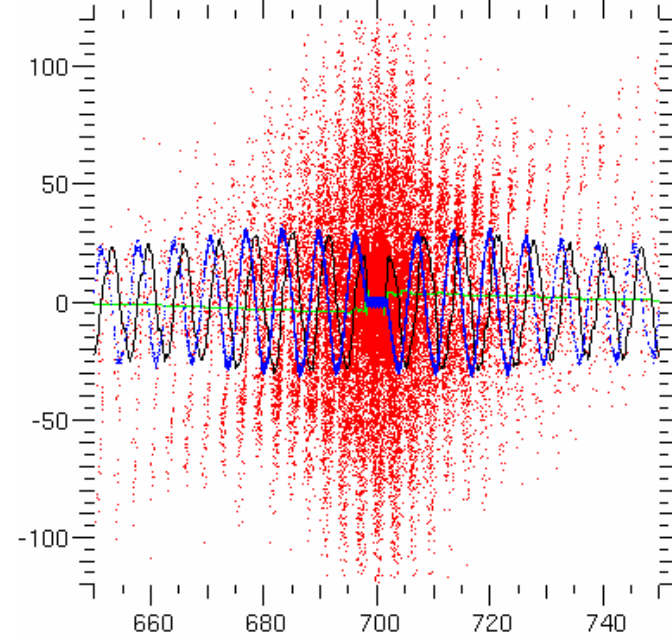
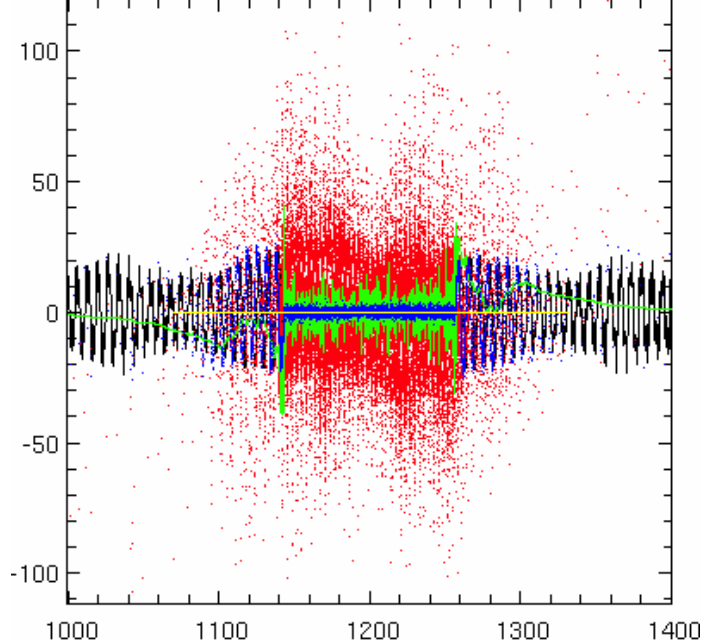
Future Work: Needs to study the detailed interactions of laser induced EM fields with hot electrons and pairs inside target



20 μm
Au foil
 10^{21} Wcm^2
150 fs
laser



2 μm
foil

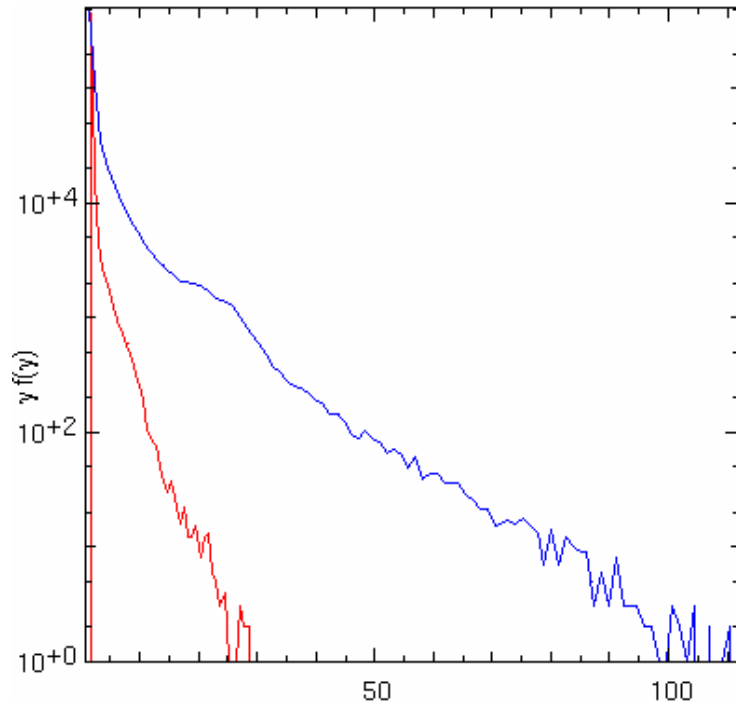


20 μm foil

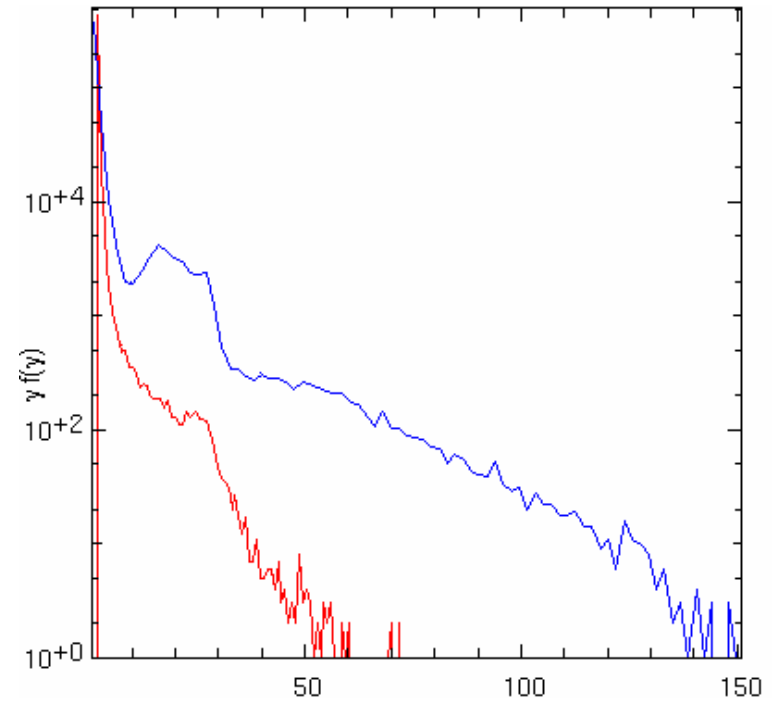
2 μm foil

blue=2-sided irradiation,

red=1-sided irradiation



20 μm foil



2 μm foil

2-sided irradiation of a thin foil seems to produce much hotter electrons for pair production

Summary

1. Both experiments and numerical simulations point towards copious production of e^+e^- pairs using lasers with $I > 10^{20} \text{ Wcm}^{-2}$.
2. We estimate max. e^+ yield $> 10^{12}$ per kJ of laser energy (conversion efficiency \sim few %).
3. The in-situ e^+ density can exceed 10^{17} cm^{-3} .
4. Such a dense and intense e^+ source makes it feasible to think about creating a BEC of Ps at cryogenic temperatures.
5. Such a Ps BEC would be a strong candidate for a GRASAR