The First Blaise Pascal Lecture Ecole Polytechnique 10/22/09

Laser Acceleration and High Field Science: 1979-2009

Toshi Tajima Blaise Pascal Chair, ENS, Paris and LMU,MPQ, Garching

Acknowledgments for Advice and Collaboration: G. Mourou, Iate-J. Dawson, N. Rostoker, F. Krausz, D. Habs, S. Karsch, L. Veisz, F. Gruener, T. Esirkepov, M. Kando, K. Nakajima, A. Chao, A. Suzuki, F. Takasaki, S. Bulanov, A. Giullietti, F. Mako, X. Yan, J. Meyer-ter-Vehn, W. Leemans, T. Raubenheimer, A. Ogata, A. Caldwell, P. Chen, Y. Kato, Iate-A. Salam, M. Downer, S. Ichimaru, M. Tigner, V. Malka, A. Henig, H.C. Wu, K. Kondo, Y. Sano, M. Abe, S. Kawanishi, M. Hegelich, D. Jung, P. Shukla

Can the society continue to support ever escalating accelerators?



Accelerator = crown of 20th C science



LHC at CERN

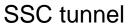


hadron therapy accelerator and gantry

supermagnets quench



Terminated Texas tunnel. The SSC was abandoned after about 25% of the tunnel for the 87-kilometercircumference large collider ring had been bored.



Demise of SSC (Super collider)





Terminated Texas tunnel. The SSC was abandoned after about 25% of the tunnel for the 87-kilometercircumference large collider ring had been bored. By largest machine to probe smallest of structure of matter

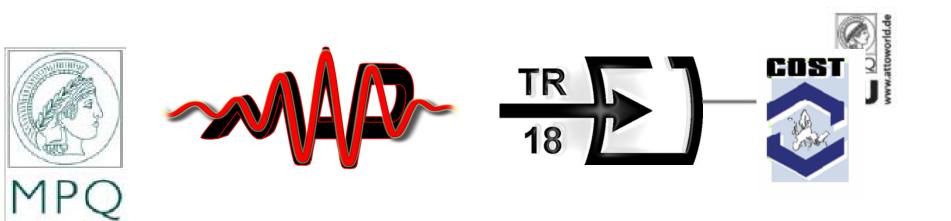
| size | 10^2 km |
|--------|---------------|
| energy | 20TeV |
| cost | \$10 B |

US:

Texas site decided (1989)

US Government decided to terminate its work: 1993

Tajima: 'Tamura Symposium' on <u>the Future of Accelerator</u> <u>Physics</u> @ UT Austin (1995)





Dream Beams Symposium

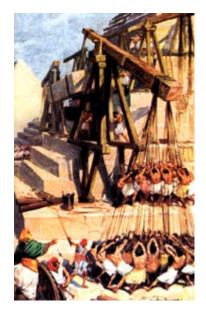
MPQ Garching Feb. 26 – 28, 2007

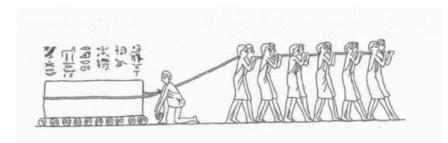
(given by F. Krausz and J. Meyer-ter-Vehn)

What is *collective force* ?



How can a Pyramid have been built?





Individual particle dynamics \rightarrow <u>Coherent</u> and <u>collective</u> movement

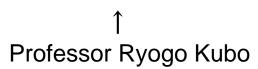
Collective acceleration (Veksler, 1956; Tajima & Dawson, 1979) Collective radiation (N² radiation) Collective ionization (N² ionization) Collective deceleration (Tajima & Chao, 2008; Ogata, 2009)

Tutelage by giants of collective phenomena

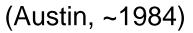


Physics of individual particles; Physics of collection of particles---collective phenomena





Professor Iliya Prigogine



Advent of collective acceleration (1956)



CERN Symposium

ON HIGH ENERGY ACCELERATORS AND PION PHYSICS

Geneva, 11th - 23rd June 1956

Proceedings

COHERENT PRINCIPLE OF ACCELERATION OF CHARGED PARTICLES

V. I. VEKSLER

Electrophysical Laboratory, Academy of Sciences, Moscow

This paper will include a very brief description of a new principle regarding the acceleration of charged particles.

In all existing accelerators of charged particles, the constant and varying electric field accelerating them is created by a powerful external source, and hence the strength of the field is independent, in the first approximation, of the number of particles which are being accelerated.

In resonance accelerators, the electromagnetic field has to be synchronized with the movement of the particles (this is of particular importance in linear accelerators). Finally, none of the existing methods permits the acceleration of neutral bunches of particles.

A new principle of particle acceleration is set forth below. Its distinctive feature lies in the fact that the particle-accelerating electric field is produced by the interaction of a geometrically small group of accelerated particles with another group of charges, plasma or an electromagnetic wave. This method has a number of important features. It appears, in the first place, that the magnitude of the accelerating field produced by this interaction act acting on each particle dependence in the sume to Theoretical studies of various aspects of the coherent acceleration method have been made by M. S. Rabinovich, A. A. Kolomenski, B. M. Bolotovski, L. V. Kovrizhnikh and I. V. Iankov, as well as by A. I. Akhiezer, Ia. Fainberg and their collaborators. The calculations made by these theoretical workers shed light on a number of complicated problems connected with the development of the different variants of this new acceleration principle, and it therefore seems appropriate to describe the new method despite the fact that a great many problems involved still await solution.

1. Acceleration of charged bunches by means of the medium

It was pointed out in a paper by Tamm that the loss of energy by particles due to Čerenkov radiation could be reversed, i.e. the medium travelling at a great velocity past charged particles should be able to convey energy to the latter. Up to now, however, no attention has been paid to the possibility of developing an acceleration process of this kind by using a high density electron beam (plasma) as the moving medium. Of course, if a single charge e is

Prehistoric activities (1973-75,....84)



Professor N. Rostoker

Collective ion acceleration by a reflexing electron beam: Model and scaling

F. Mako Naval Research Laboratory, Washington, D. C. 20375 T. Tajima

Institute for Fusion Studies, University of Texas, Austin, Texas 78712

(Received 21 June 1983; accepted 2 April 1984)

Analytical and numerical calculations are presented for a reflexing electron beam type of collective ion accelerator. These results are then compared to those obtained through experiment. By constraining one free parameter to experimental conditions, the self-similar solution of the ion energy distribution agrees closely with the experimental distribution. Hence the reflexing beam model appears to be a valid model for explaining the experimental data. Simulation shows in addition to the agreement with the experimental ion distribution that synchronization between accelerated ions and electric field is phase unstable. This instability seems to further restrict the maximum ion energy to several times the electron energy.

I. INTRODUCTION

Experiments on collectively accelerating ions utilizing a reflexing intense relativistic electron beam in a plasma have been carried out.^{1,2} These experiments began to reveal severchronous fashion. Thus, energetic ions would be expected. The ion energy would, of course, be bounded above by the ion to electron mass ratio times the initial electron energy; that is, the energy is bounded when the ions reach the initial Collective acceleration suggested: Veksler (1956) (ion energy)~ (M/m)(electron energy)

Many experimental attempts (~'70s):

led to no such amplification (ion energy)~ (several)x(electron)

Mako-Tajima analysis (1978;1984)

sudden acceleration, ions untrapped, electrons return

 \rightarrow #1 gradual acceleration necessary

→ #2 electron acceleration possible with trapping (with Tajima-Dawson field), more tolerant for sudden process

Path once trodden

Collective acceleration of ions by electron beam

F.Mako / T. Tajima

lons <u>left out</u>, while electrons shoot backward

- \rightarrow laser electron acceleration (1979)
- → laser ion acceleration of limited ion mass (2009)

The electric field is

$$\epsilon = \frac{\phi_0}{n_f} \frac{5}{36} \left(\frac{6}{\sqrt{3}} - \frac{z}{n_f} \right)$$

where the conservation of energy was used as a boundary condition, i.e.,

 $U^2/2 + \psi = 0$ at $\zeta = 0$.

The maximum ion energy can now be obtained by setting $n_i = 0$, i.e,

$$E_{imax} = 6q\phi_0$$
 at $\zeta = 6/\sqrt{3}$

In the experiment the diode voltage was 0.8 MV and the ions were doubly ionized helium,⁶ thus the maximum ion energy predicted by theory is

 $E_{\text{imax}} = 9.6 \text{ MeV}$.

The experimental result⁶ for the maximum helium ion energy was 9.6 MeV and therefore is in good agreement with the theory.

The ion number as a function of energy is calculated to be

$$N_{i}(E_{i}) = \frac{n_{0}A}{\beta} \left[\left(\frac{6}{5}\right)^{1/2} - \left(\frac{E_{i}}{5q\phi_{0}}\right)^{1/2} \right]^{6}, \quad (15)$$
here

$$n_0 A = \frac{16}{5} \frac{J_0 A}{\epsilon} \left(\frac{2m}{e\phi_0}\right)^{1/2},$$

 $\beta = (i)^{1/2} (1/v_{ef})$

 $A = \pi r_b^2$, $r_b = \text{electron beam radius}$,



and

Equation (15) is our main result. The natural logarithm of Eq. (15) is plotted in Fig. 2 along with the experimental

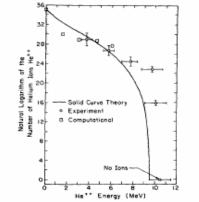


FIG. 2. Comparison between theory, experiment, and simulation, of the natural logarithm of the ion number versus energy.



data. The following experimental values were used: $J_{\phi} = 40$ kA, $\phi_{\phi} = 0.8$ MV, q = 2e (doubly ionized helium), t = 100 ns and $r_{\phi} = 2.5$ cm. The agreement between Eq. (15) and the experiment⁶ is reasonable. The relation in Ref. 3 does not provide such a good fit: it has too weak a slope.

III. SCALING AND ACCESSIBILITY OF THE MODEL

In the preceding section, the analysis assumed that a self-similar state could be reached. To address the question of whether a self-similar state can be attained, a detailed analysis of the initial value problem is required. This detailed analysis should include a self-consistent treatment of the dy-

ELECTRON PHASE SPACE (TIME 1.5/was)

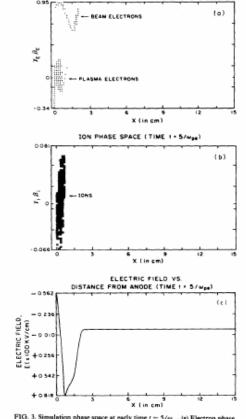


FIG. 3. Simulation phase space at early time $t = 5/\omega_{\mu\nu}$, (a) Electron phase space (beam and plasma), (b) ion phase space, (c) electron field versus position.

Laser Acceleration of **Electrons**

← Lesson #2 trapping of electrons easier



Gradient limit : breakdown threshold for microwave (< 100MeV/m)

E. Lawrence: cyclotron (c. 1932)

SSC: 10^2 km circumference († 1993); Linear Collider: > 10km (~2020?)

Plasma : already 'broken' matter. No breakdown threshold.

'collective ion acceleration' (Veksler, 1956): ion trapping difficult ($v_{tr,ion} << c$) Introduction of laser acceleration (Tajima and Dawson, 1979)

Linear EM field: cannot accelerate: *Woodward-Lawson Theorem* <u>Strong nonlinear fields</u>

longitudinal acceleration (rectification of laser fields; v x B/c ~ O(E)) laser plays master, plasma slaves----- provides <u>hard structure</u>

electron trapping possible (revisit of ion acceleration now) $(v_{tr,e} \sim c)$

→ High Field Science

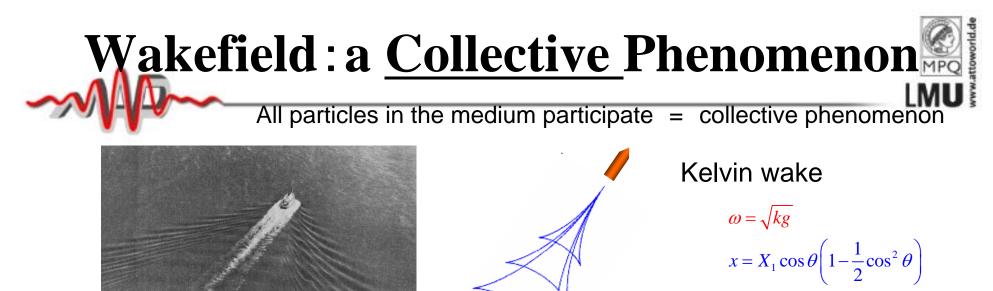
<u>Ultrafast pulses</u>

fs regime: ions immobile; enhanced with <u>collective</u> electron resonance absence of 'notorious' hydrodynamical plasma instabilities; controllability;

relatively small laser energy (e.g. ELI)

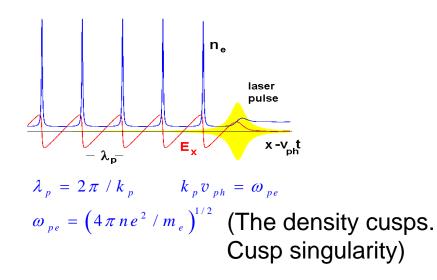
<u>Large gradient</u> (> 10 GeV/m, leap by > 3 orders of magnitude)

Low emittance (< mm mrad regime)



 $y = X_1 \cos^2 \theta \sin \theta$ $-\pi / 2 < \theta < \pi / 2$

No wave breaks and wake peaks at v≈c

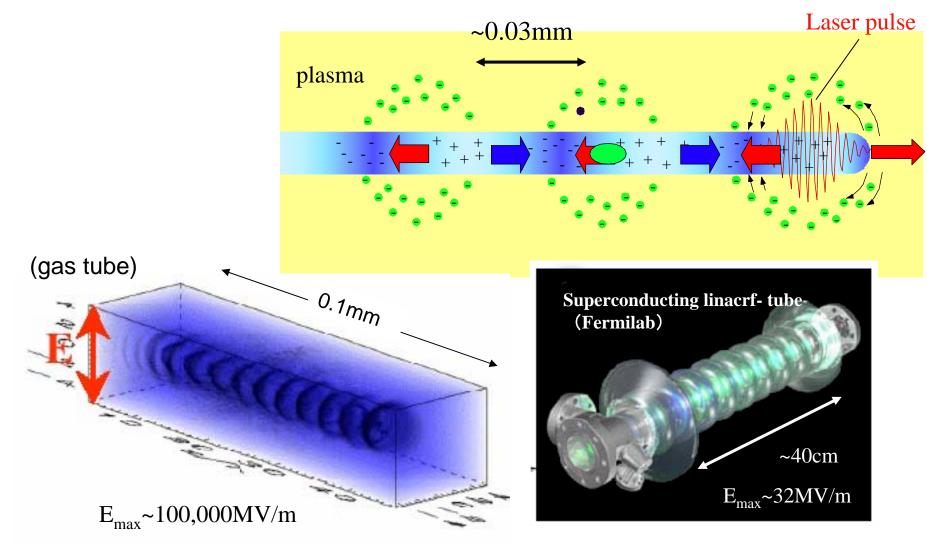


Wave **breaks** at v<c



Thousand-fold Compactification

Laser wakefield: thousand folds gradient (and emittance reduction?)



The late Prof. Abdus Salam





At ICTP Summer School (1981), Prof. Salam summoned me and discussed about laser wakefield acceleration.

Salam: 'Scientists like me began feeling that we had less means to test our theory. However, with your laser acceleration, I am encouraged'. (1981)

He organized the Oxford Workshop on laser wakefield accelerator in 1982.

Effort: many scientists over many years to realize his vision / dream **High field science**: spawned

Laser technology invented (1985)





Chirped pulse amplification (CPA) invented: to overcome the gain medium nonlinearities in spatially expanded amplification to temporal expansion: smaller, shorter pulse, more intense, higher reprate, all simultaneous.

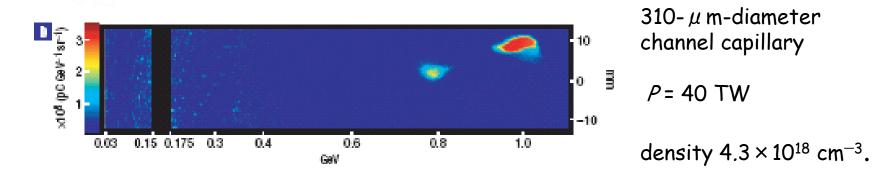
(Professor Gerard Mourou)

→ many table-top TW and PW lasers world-wide first Chair, ICUIL (International Committee for Ultra Intense Lasers) toward EW laser (*Extreme Light Infrastructure*)

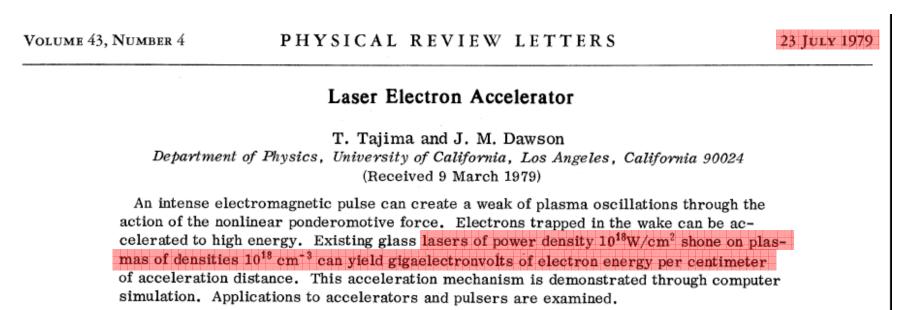
 →First LWFA experiments (Nakajima et al 1994; Modena et al1995)
 →drives High Field Science

GeV electrons from a centimeter accelerator (a slide given by S. Karsch)





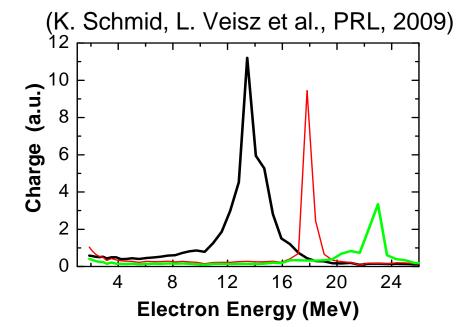
Leemans et al., Nature Physics, september 2006



MPQ Laser Acceleration Effort (1)



Monoenergy electron spectra: from few-cycle laser (LWS-10)

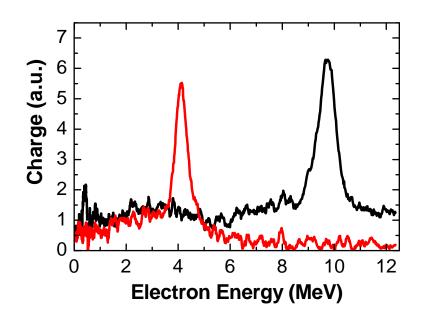


Large electron spectrometer 2 – 400 MeV

- No thermal background !
- Energies: 13.4 MeV, 17.8 MeV, 23 MeV
- FWHM energy spread: 11%, 4.3%, 5.7 %
- ~ 10 pC charge

Small electron spectrometer:

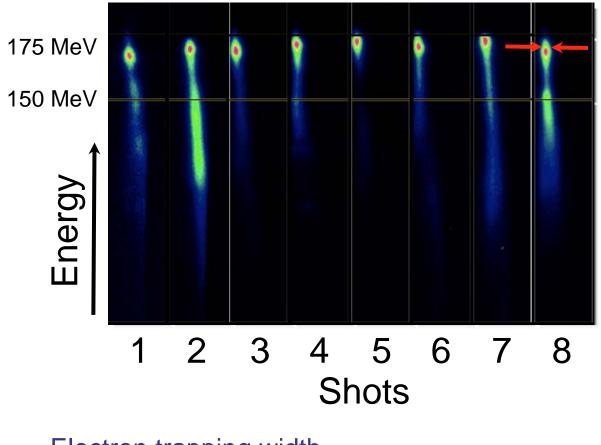
- Electron energies below 500keV
- No thermal background !
- 4.1 MeV (14%); 9.7 MeV (9.5%)



MPQ Laser Acceleration Effort (2)



Reproducible acceleration conditions



E ≈ 169.7 ± 2.0 MeV

1.1% peak energy fluctuation !

 $\Delta E/E \approx 1.76 \pm 0.26\%$ RMS \rightarrow Essential property for future table-top FEL operation

Source size image: provides emittance measurement, given the resolution can be improved

Electron trapping width $V_{tr,e} \sim c\sqrt{a_0}$

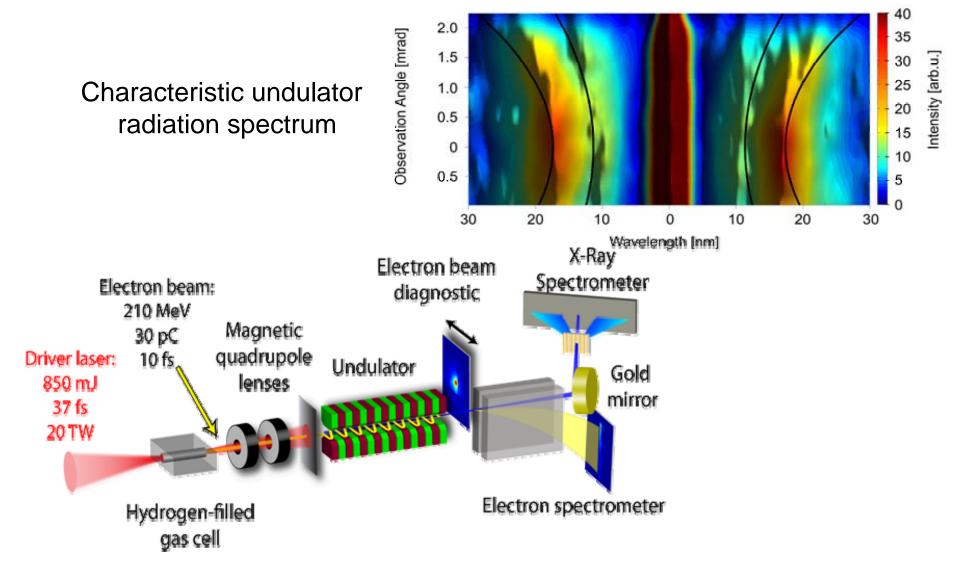
(J. Osterhoff,...S. Karsch, et al., PRL 2008)

MPQ Laser Acceleration Effort (3)



Laser-driven Soft-X-Ray Undulator Radiation

(F. Gruener, S. Karsch, et al., Nature Phys., 2009)



Intra-Operatory Radiation Therapy (IORT)

VS.



LWFA electron sources: technology transferred to company

NOVAC7 (HITESYS SpA) RF-based

El. Energy < 10 MeV (3, 5, 7, 9 MeV)

Peak curr. 1.5 mA Bunch dur. 4 µs Bunch char. 6 nC

> Rep. rate 5 Hz Mean curr. 30 nA

Releas. energy (1 min) @9 MeV (≈dose) 18 J Releas. energy (1 min) @20 MeV (≈dose) 21 J





(A. Giulietti et al., Phys. Rev. Lett.,2008)

Peak curr.> 1.6 KABunch dur.< 1 ps</td>Bunch char.1.6 nC

Rep. rate10 HzMean curr.16 nA

CEA-Saclay

experim. source

Laser-based

El. Energy > 10 MeV

(10 - 45 MeV)





Beam dump: <u>harder to stop</u> and more <u>hazardous radioactivation</u> ↓

Gas (plasma) collective force to shortstop the HE beams

- the shorter the bunch is, the easier to stop
 - (ideally suited for laser wakefield accelerated beams)
- little radioactivation (good for environment)
 - example of <u>'Toilet Science</u>' that tends impact of own produce (as opposed to 'Kitchen Science' of 20th C)
- possible energy recovery

Beam Stopping and its Energy Recovery Using Plasma

> February 25, 2008 Toshiki Tajima and Alexander W. Chao

Tajima and Chao, (2008 applied for patent) H. C. Wu et al. (2009)

1 Motivations

1.1 Beam Stopping

In the effort to make a high energy accelerator system as compact as possible, it is necessary not only to make the accelerator compact, but also to make the beam stopping system compact. With this motivation, we introduce the concept of passive plasma decelerator at the end of the use of the high energy beam by immersing the beams to be decelerated into an appropriately designed plasma.

Stopping power due to collective force

Bethe-Bloch stopping power in <u>matter</u> <u>Plasma</u> stopping power due to individual force

$$-(dE/dx)_{ind} = (F/\beta^2)\ln(m_e v^2/e^2 k_D)$$

That due to collective force (perturbative regime) $-(dE/dx)_{coll} = (F/\beta^2)\ln(k_D v/\omega_{pe})$ $F = 4\pi e^4 n_{e,m}/m_e c^2 = e^2 k_{pe,m}^2$ (Ichimaru, 1973)

Professor Setsuo Ichimaru

Plasma stopping power due to short-bunch <u>wakefield</u> (wavebreak regime) $-(dE/dx)_{C} = m_{e}C\omega_{pe}(n_{b}/n_{e})$

(Wu et al, 2009)

Greater by several orders in gas over Bethe-Bloch in solid



Key issues of future colliders



(T. Raubenheimer, SLAC, 2008)

Beam Acceleration

- * Largest cost driver for a linear collider is the acceleration
 - − ILC geometric gradient is ~20 MV/m \rightarrow 50km for 1 TeV
- * Size of facility is costly \rightarrow higher acceleration gradients
 - High gradient acceleration requires high peak power and structures that can sustain high fields
 - · Beams and lasers can be generated with high peak power
 - Dielectrics and plasmas can withstand high fields
- Many paths towards high gradient acceleration
 - High gradient microwave acceleration ~ ~100 MV/m
 - Acceleration with laser driven structures
 - Acceleration with beam driven structures
 - Acceleration with laser driven plasmas
 - Acceleration with beam driven plasmas
- SLAC

13th AAC Workshop July 27 - August 2, 2008 Page 11

~1 GV/m

~10 GV/m

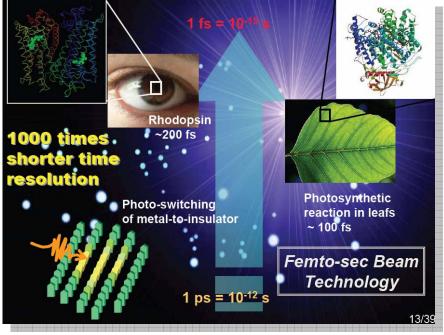


Challenge Posed by DG Suzuki



Frontier science driven by advanced accelerator



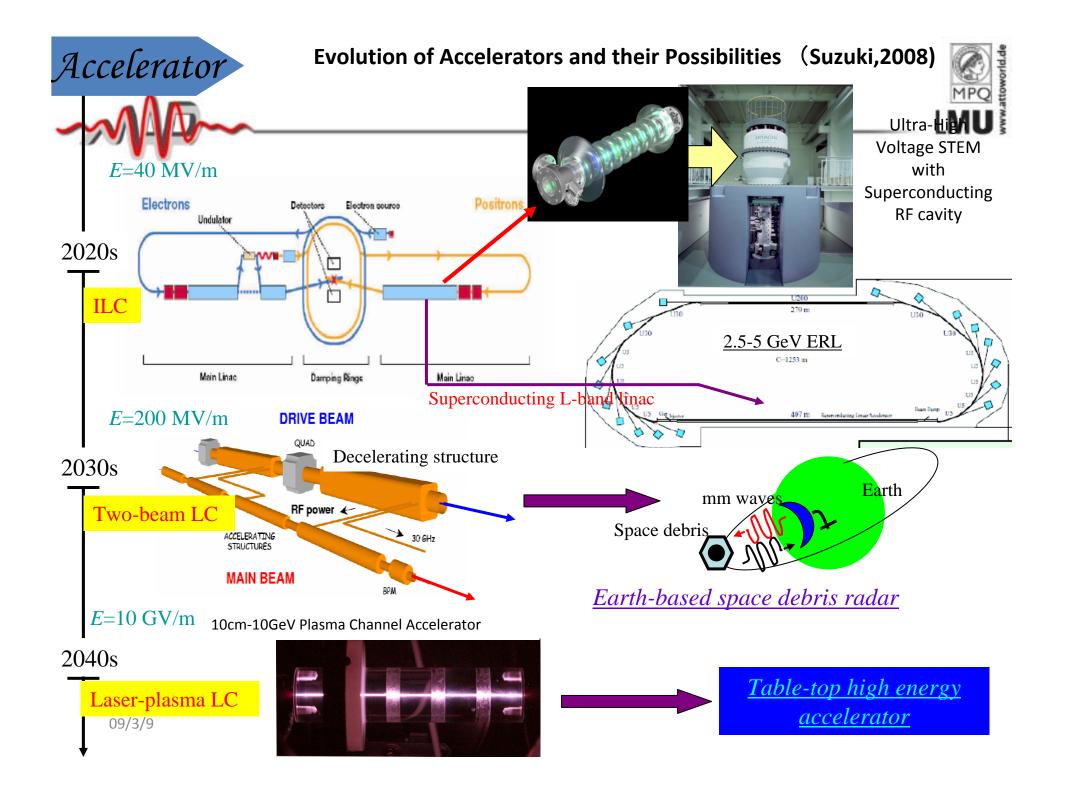


compact, ultrastrong a

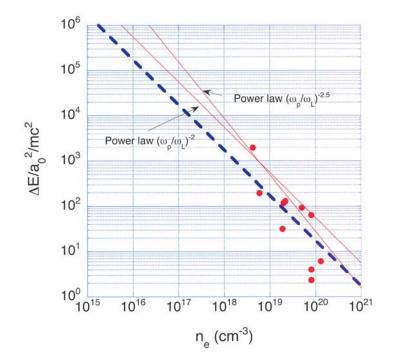
Can we meet the challenge?

atto-, zeptosecond

A. Suzuki @KEK(2008)



Meeting Suzuki's Challenge: aser acceleration toward ultrahigh energies MU $\Delta E \approx 2m_0 c^2 a_0^2 \gamma_{ph}^2 = 2m_0 c^2 a_0^2 \left(\frac{n_{cr}}{n_e}\right), \text{ (when 1D theory applies)}$



| $L_d = \frac{2}{\pi} \lambda_p c$ | $a_0^2 \left(\frac{n_{cr}}{n_e} \right),$ | | $a_p = \frac{1}{3\pi} \lambda_p a_0$ | $\left(\frac{n_{cr}}{n_e}\right),$ |
|-----------------------------------|--|---|--------------------------------------|------------------------------------|
| | | _ | | |

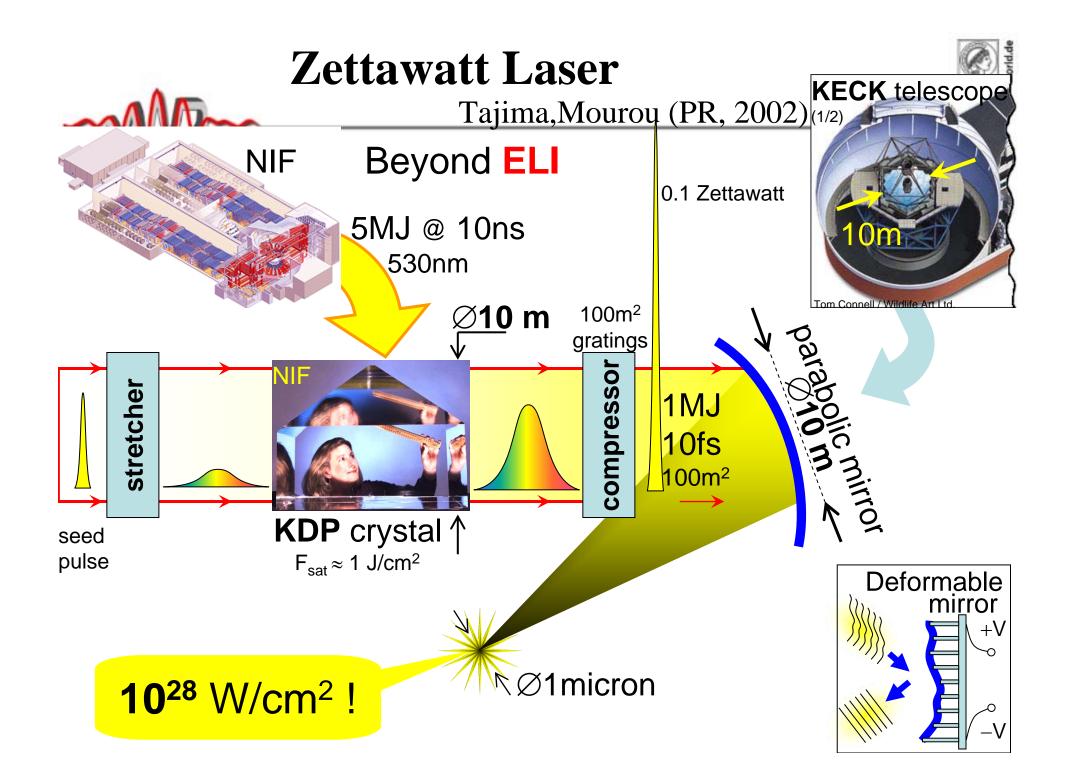
 \mathbf{r}

| | | case I | case II | case III |
|---------------------|------------------|----------------------|----------------------|----------------------|
| | | 10 | 3.2 | 1 |
| energy gain | GeV | 1000 | 1000 | 1000 |
| plasma density | cm ⁻³ | 5.7x10 ¹⁶ | 5.7×10^{15} | 5.7×10^{14} |
| acceleration length | m | 2.9 | 29 | 290 |
| spot radius | μm | 32 | 100 | 320 |
| peak power | PW | 2.2 | 2.2 | 2.2 |
| pulse duration | ps | 0.23 | 0.74 | 2.3 |
| laser pulse energy | kJ | 0.5 | 1.6 | 5 |

Even 1PeV electrons (and γ s) are possible, albeit with lesser amount

→ exploration of new physics such as the reach of relativity and quantum gravity (correlating with primordial gamma-ray burst [GRB] observation)?

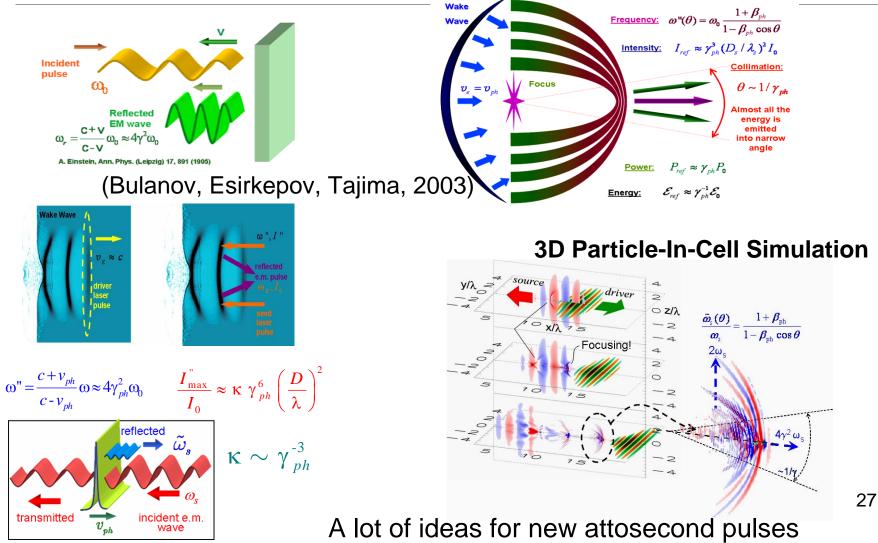
(laser energy of 10MJ@plasma density of 10¹⁶/cc; maybe reduced with index 5/4)



Relativistic Engineering: relativity as the guiding principle (cf. quantum engineering)

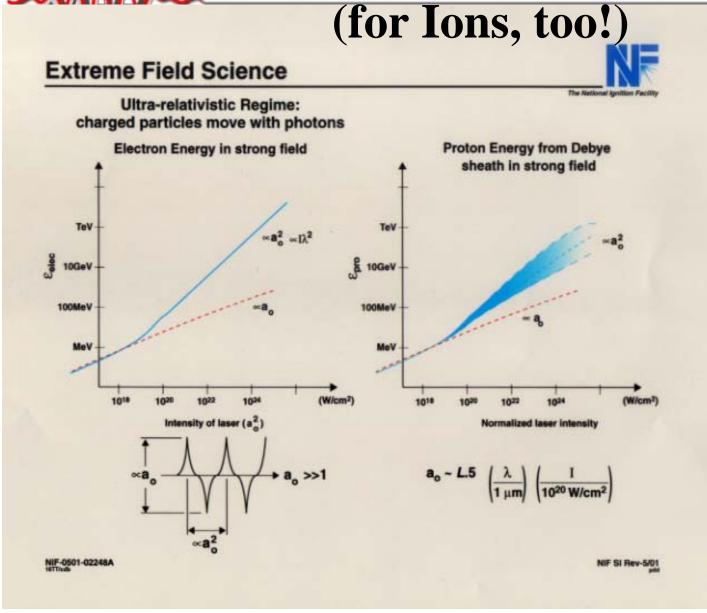


EM Pulse Intensification and Shortening by the Flying Mirror



Relativity Helps Acceleration

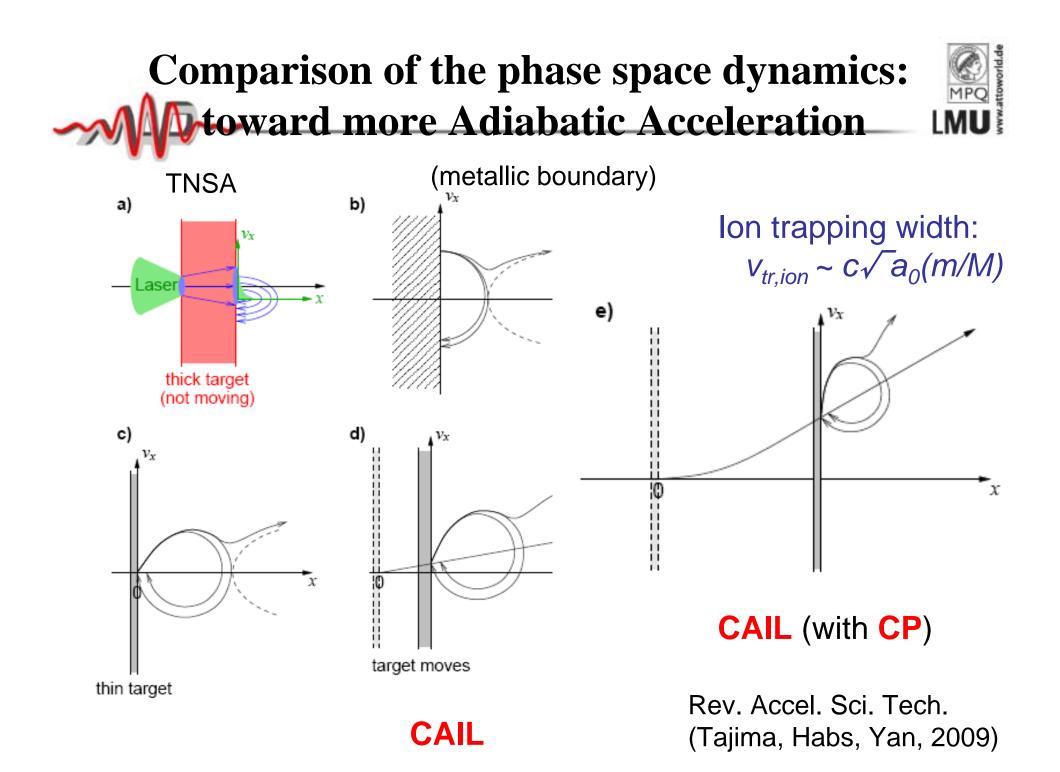


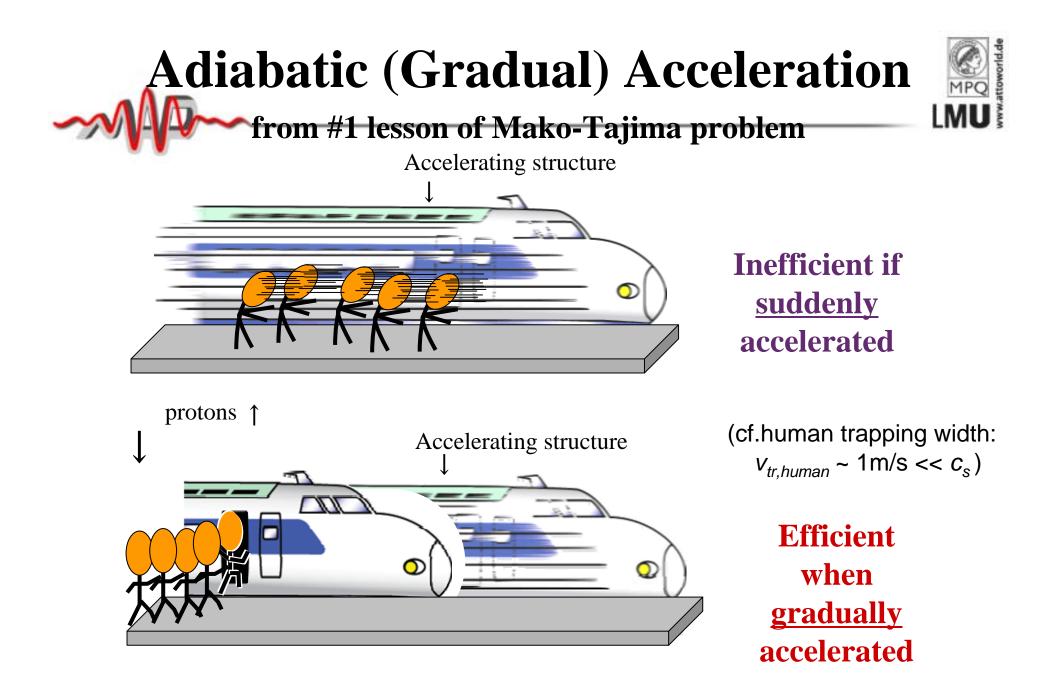


Strong fields: rectifies laser to longitudinal fields

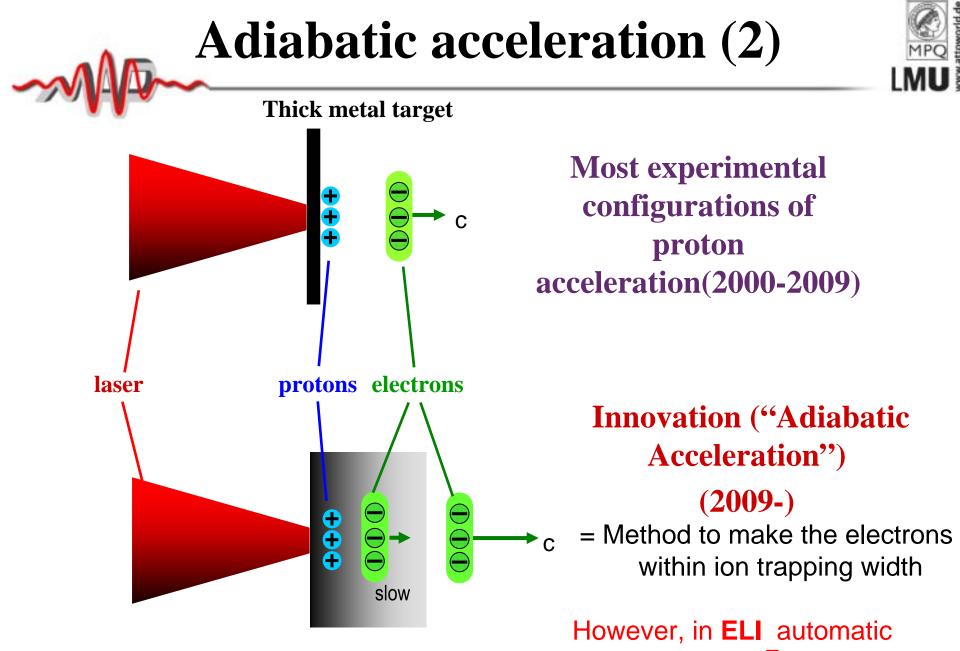
In <u>relativistic</u> regime, photon x electrons and even protons couple stronger.

> (Tajima, 1999 @LLNL; Esirkepov et al., PRL,2004)





Lesson #1: gradual acceleration \rightarrow Relevant for ions

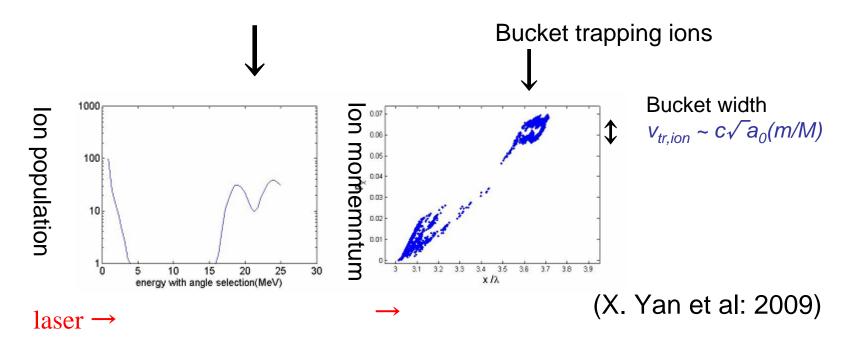


Graded, thin (nm), or clustered target and/or circular polarization

 $v_{tr, ion} \sim c \sqrt{a_0(m/M)} \sim c$ (ultrarelativistic $a_0 \sim M/m$)

Good quality ion beams

Circularly Polarized Laser drives ions out of ultrathin (nm) foil **adiabatically Monoenergy** peak emerges; energy **more rapidly** increases as $\sim a_0^2$



Ponderomotive force drives electrons, Electrostatic force nearly cancels Slowly accelerating bucket formed

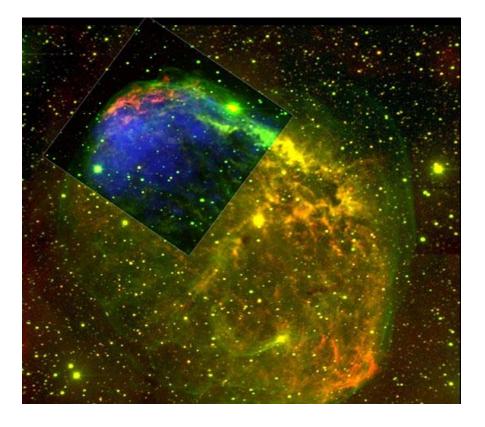
good-quality and efficient acceleration of ions

Conclusions



- Collective acceleration: hard birth / long way and near maturation(electron→ion; laser→electron; laser→photon; electron→electron; ion→electron); unexpected 'homecoming' (laser →ion), too
- Leap by many orders (≥ 3) in many respects; equally more demanding by many orders : N² vs. N.
- Laser has come around to match the condition set 30 years ago; Still some ways to go to realize the dream (such as ELI)
- GeV electrons; 10 GeV soon; 100GeV considered; TeV laser collider contemplated; PeV ?
- Societal obligations and applications: already <u>beginning</u>, soon to flourish (e.g., cancer therapy, radiolysis, bunch decelerator, nuclear detection, compact FEL source, compact radiation sources, ultrafast diagnosis,...)

Cosmic Acceleration in the Wake of Intense Radiation



UHECR (ultra high energy cosmic rays): beyond **Fermi** acceleration necessary, wakefiled acceleration?

Merci Beaucoup et a la Prochaine Fois!



In dedication to the late-Professor John Dawson

I plan to give Pascal Lectures approximately once a month from now on. Look forward to hearing your opinions and feedbacks. Toshi Tajima