

The Review Talk  
PIF (Physics in Intense Fields) 2010  
Friday, Nov. 26, 2010  
KEK, Tsukuba

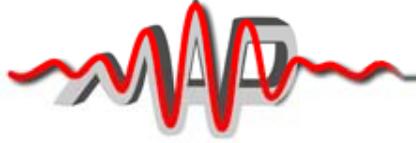
# High Field Science

Toshiki Tajima\*

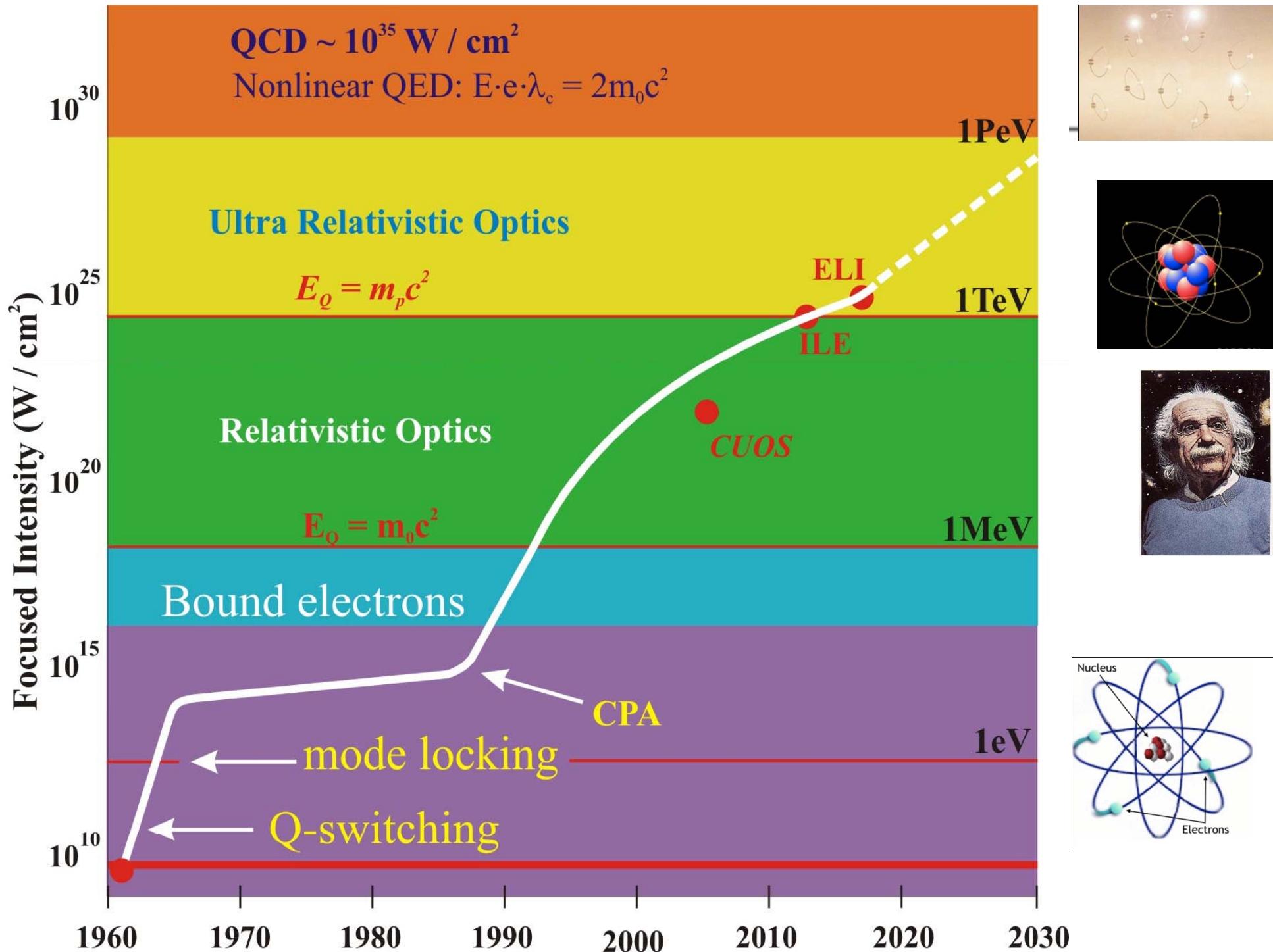
Ludwig Maximilian University Munich,  
and MPQ,  
Garching, Germany

\*Blaise Pascal Chair , Ecole Normale Supérieure

Acknowledgments for collaboration and advice: G. Mourou, F. Krausz, G. Tsakiris, D. Habs, K. Homma, Esirkepov, S. Bulanov, M. Kando, W. Sandner, X. Q. Yan, S. Karsch, F. Gruener, M. Zepf, N. Naumova, H. Gies, E. Moses, M. Downer, P. Corkum, Y. Kato, C. Barty, G. Dunne, R. Schuetzhold, K. Homma, A. Suzuki, F. Takasaki, M. Nozaki, K. Nakajima, E. Goulielmakis, M. Teshima, Y. Fujii, A. Caldwell, W. Leemans, L. Veisz



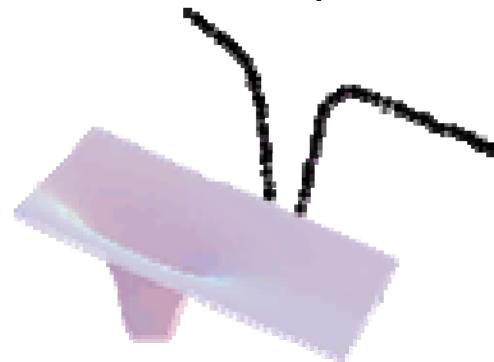
1. Laser intensity and nonlinearities
2. Relativistic Nonlinearity = Laser Wakefield Acceleration
3. Radiation Pressure Acceleration x LWFA toward TeV over cm
4. ICAN: toward fiber laser with 100MW @ 50%
5. Prof. Suzuki's Challenge:  
    1000 times higher energies, 1000 times shorter time scales
6. PeV acceleration to probe quantum gravity fluctuations
7. Ultrafast science: toward attosecond, zeptosecond,  
    yoctosecond(?)
8. Integration of high field science and ultrafast science
9. Exploration of vacuum
10. High field science as a tool to explore light-mass new fields (dark matter such as ALP and dark energy)



# Nonlinearities in atom, plasma, and vacuum



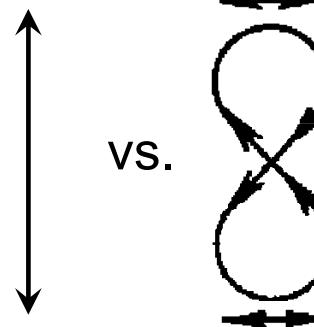
**Atomic**  
nonlinear potential



**Keldysh** field for  
laser atomic  
ionization

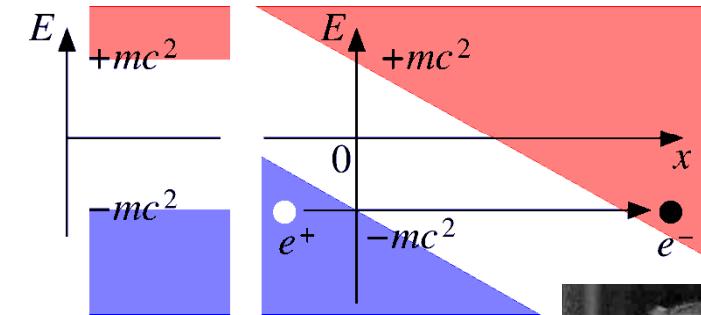
Compact high energy colliders  
Compact accelerator applications  
PeV acceleration for quantum gravity →

**Plasma** electron  
nonlinear  
relativistic motion

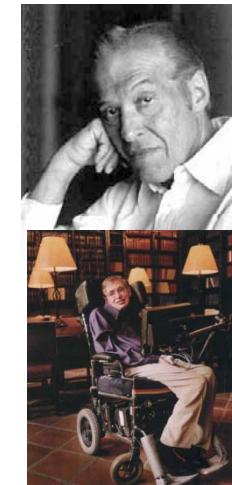


Laser wakefield

**Vacuum** nonlinearity

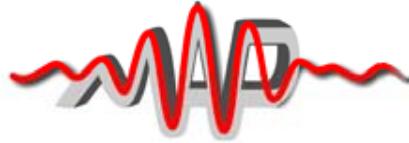


**Schwinger** field for  
vacuum breakdown



Nonlinear QED fields  
General relativistic effects  
Vacuum probe (s.a. Dark energy)

# Relativistic nonlinearity under intense laser



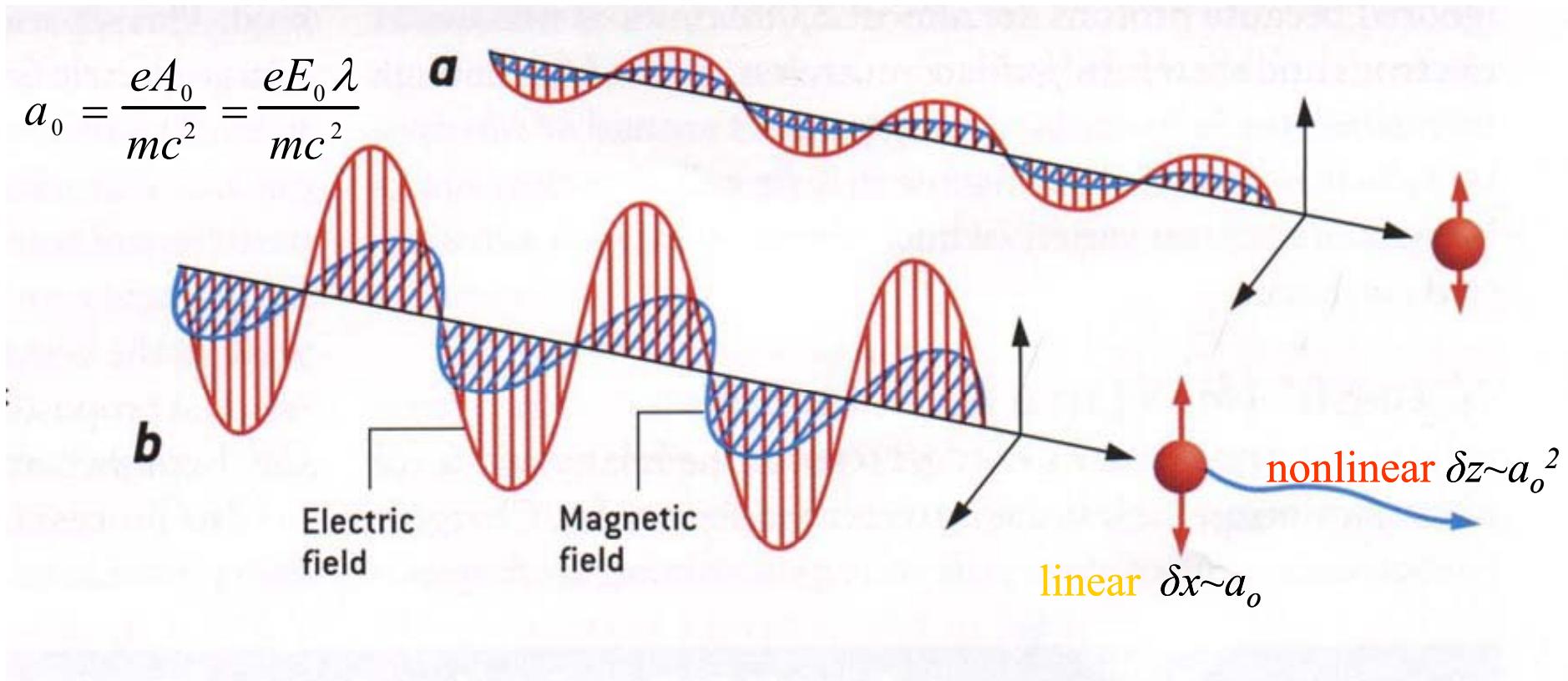
Plasma free of binding potential , but its electron responses:

a) Classical optics :  $v \ll c$ ,

$a_0 \ll 1$ :  $\delta x$  only

b) Relativistic optics:  $v \sim c$

$a_0 \gg 1$ :  $\delta z \gg \delta x$



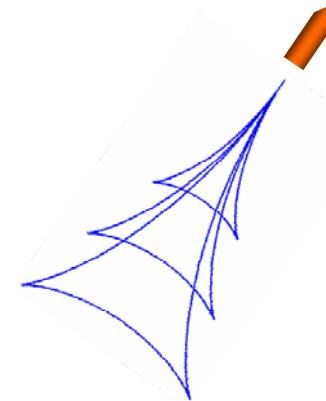


# Wakefield : a Collective Phenomenon

All particles in the medium participate = collective phenomenon



Рис. 71. Наблюдаемая картина корабельных волн. [Любезно предоставлено Aerofilms Ltd.]



Kelvin wake

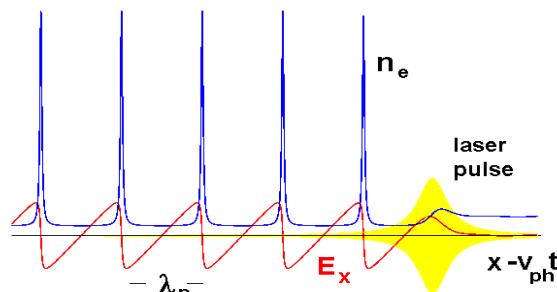
$$\omega = \sqrt{kg}$$

$$x = X_1 \cos \theta \left( 1 - \frac{1}{2} \cos^2 \theta \right)$$

$$y = X_1 \cos^2 \theta \sin \theta$$

$$-\pi/2 < \theta < \pi/2$$

No wave breaks and wake **peaks** at  $v \approx c$



$$\lambda_p = 2\pi / k_p \quad k_p v_{ph} = \omega_{pe}$$

$$\omega_{pe} = (4\pi n e^2 / m_e)^{1/2}$$

(The density cusps.  
Cusp singularity)

← relativity  
regularize

Wave **breaks** at  $v < c$

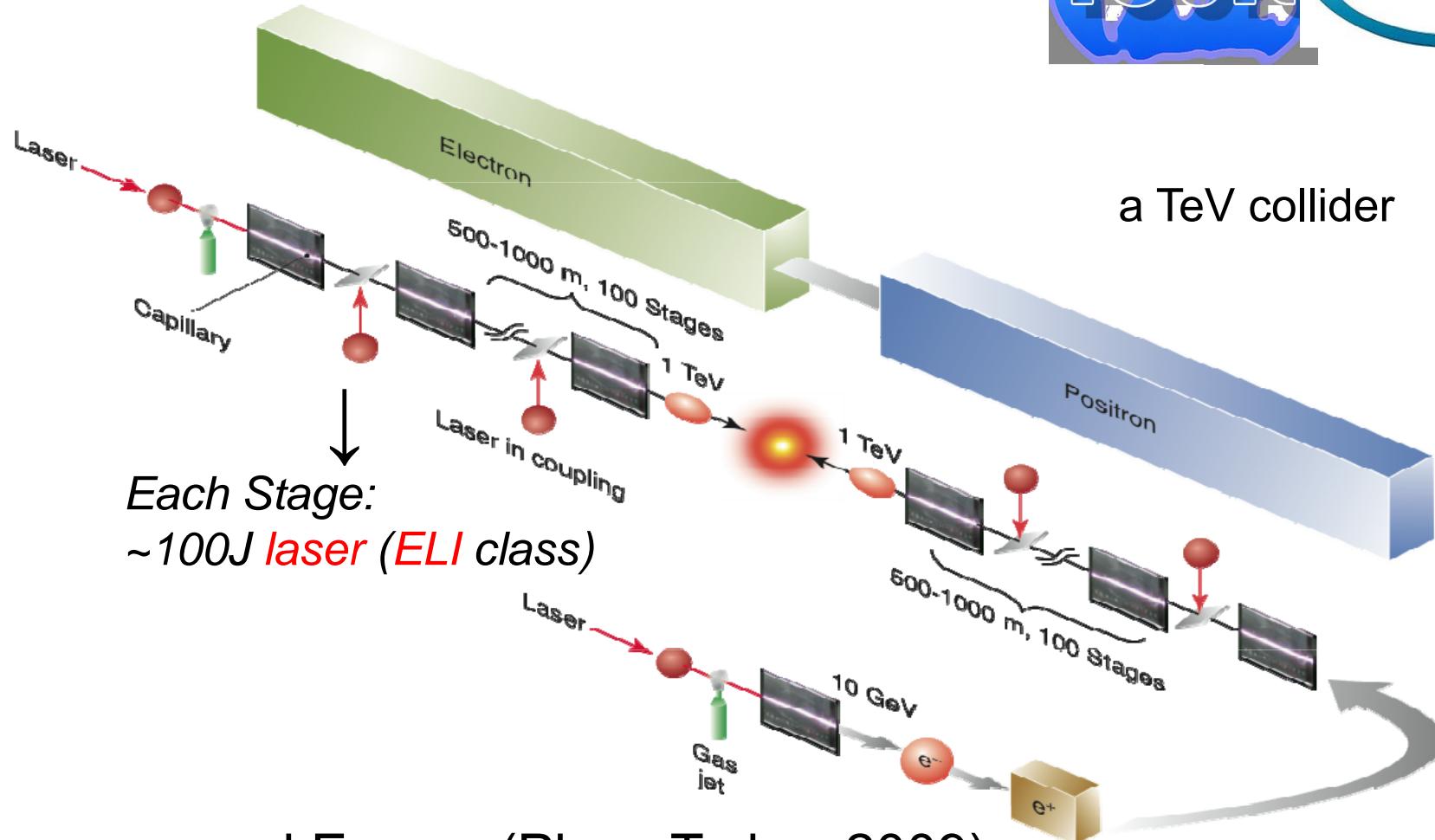


(Plasma physics vs.  
String theory?)

# Laser driven collider concept



LMF



Leemans and Esarey (Phys. Today, 2009)

ICFA-ICUIL JTF on **Laser** Acceleration (Darmstadt, 2010)



# ICFA-ICUIL Joint Task Force on laser acceleration (Darmstadt, 2010)



W. Leemans,  
Chair

Collider subgroup  
List of parameters  
(W. Chou)

Table 1  
Collider parameters

Case	1 TeV	10 TeV (Scenario I)	10 TeV (Scenario II)
Energy per beam (TeV)	0.5	5	5
Luminosity ( $10^{34} \text{ cm}^{-2}\text{s}^{-1}$ )	1.2	71.4	71.4
Electrons per bunch ( $\times 10^9$ )	4	4	1.3
Bunch repetition rate (kHz)	13	17	170
Horizontal emittance $\gamma \varepsilon_x$ (nm-rad)	700	200	200
Vertical emittance $\gamma \varepsilon_y$ (nm-rad)	700	200	200
$\beta^*$ (mm)	0.2	0.2	0.2
Horizontal beam size at IP $\sigma_x^*$ (nm)	12	2	2
Vertical beam size at IP $\sigma_y^*$ (nm)	12	2	2
Luminosity enhancement factor	1.04	1.35	1.2
Bunch length $\sigma_z$ ( $\mu\text{m}$ )	1	1	1
Beamstrahlung parameter $\Upsilon$	148	8980	2800
Beamstrahlung photons per electron $n_\gamma$	1.68	3.67	2.4
Beamstrahlung energy loss $\delta_E$ (%)	30.4	48	32
Accelerating gradient (GV/m)	10	10	10
Average beam power (MW)	4.2	54	170
Wall plug to beam efficiency (%)	10	10	10
One linac length (km)	0.1	1.0	0.3



# ICFA-ICUIL JTF Conclusions

(April, 2010; Darmstadt)

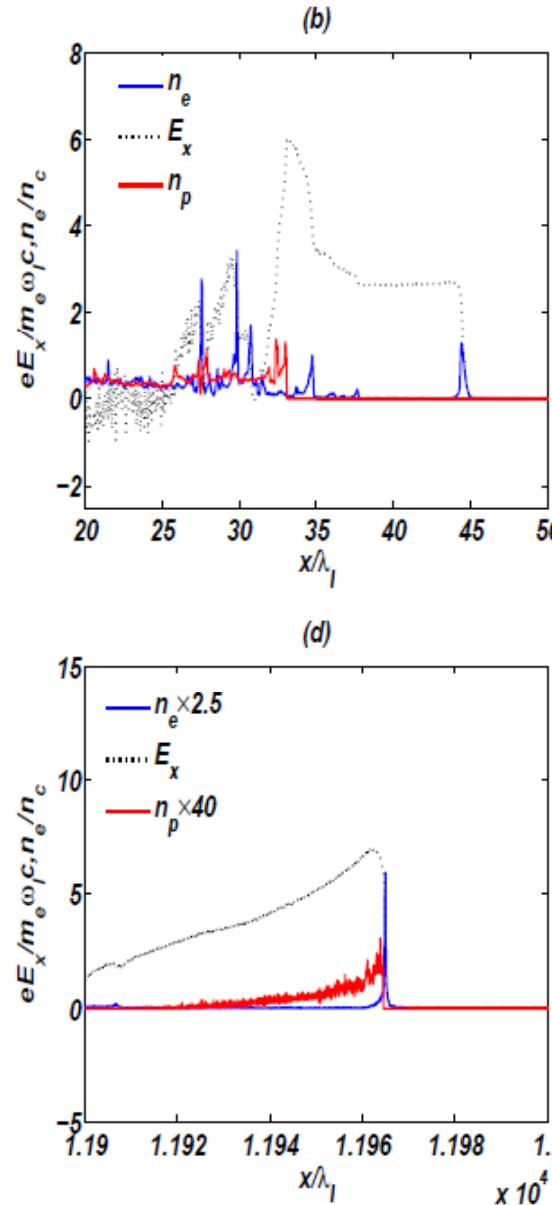
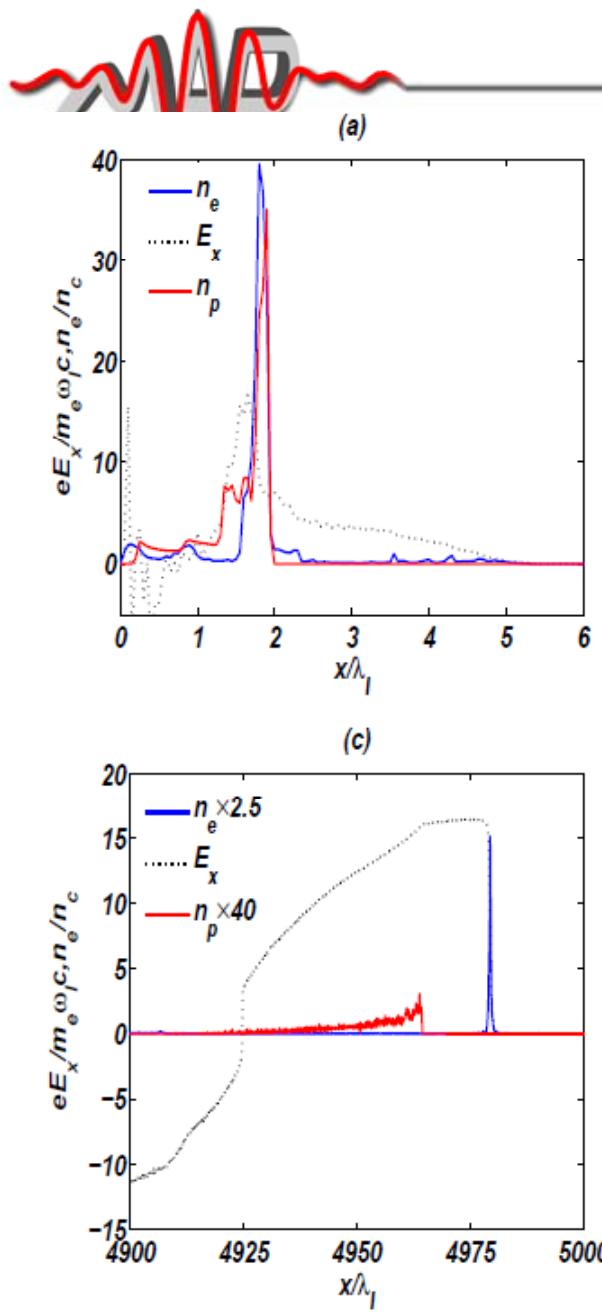


- *Science of LWFA* (US, Europe, Asia) matured to extend toward HEP accelerators
- *Laser technology lacking* suited for HEP accelerators: **laser efficiency, average power**
- *Technologies to rectify emerging and credible:*  
1. thin disk; 2. ceramic; 3. fiber **laser**,..
- *ICFA-ICUIL collaboration: important guide* of direction

→ Bridge Lab between HEP and Laser communities

“Symposium on Bridgelab for Laser Accelerator: Route toward Reality”  
L'orme (Paris), Jan. 14, 2011 (Organizers: Mourou, Tajima, ....)

# TeV ion acceleration in ELI regime



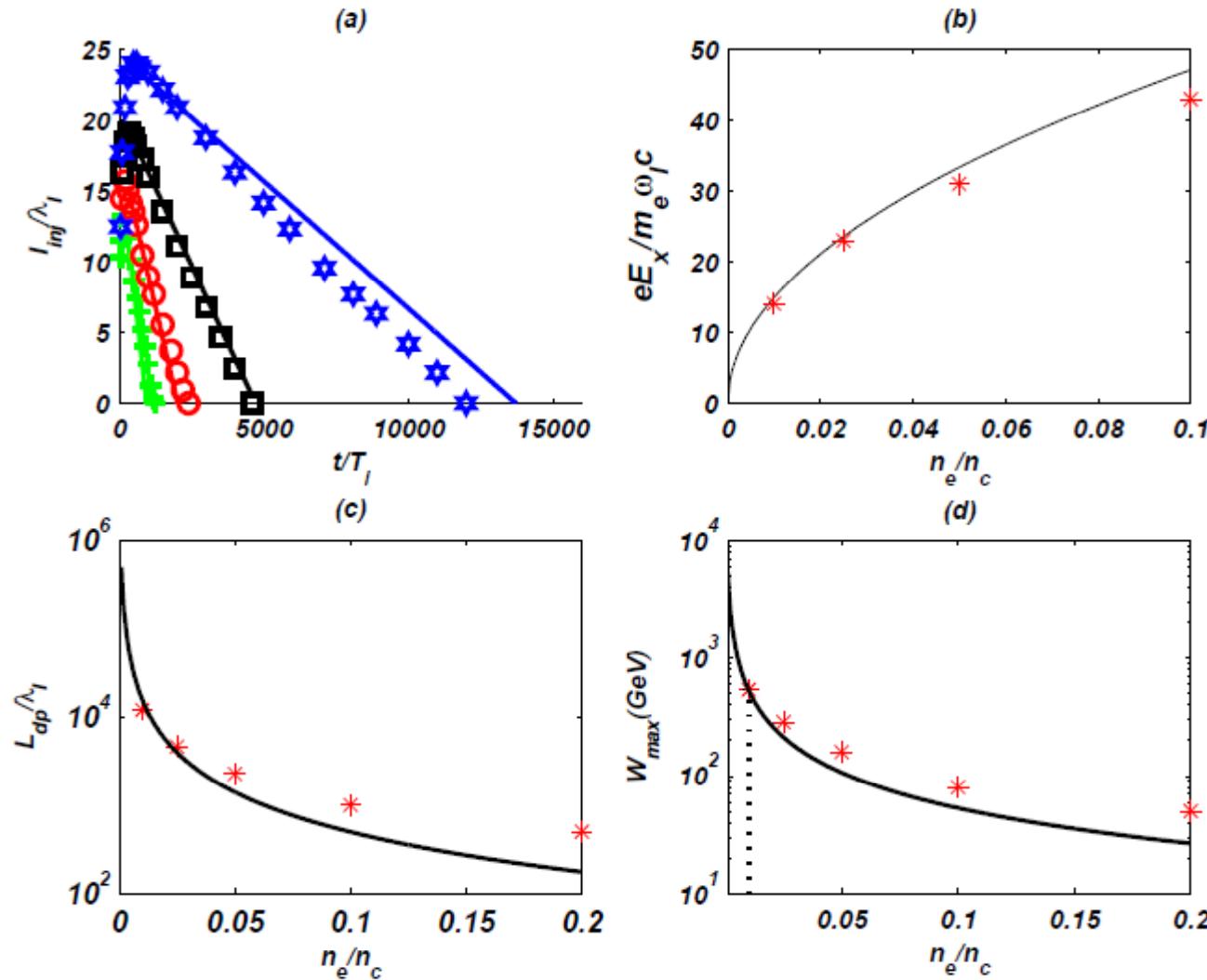
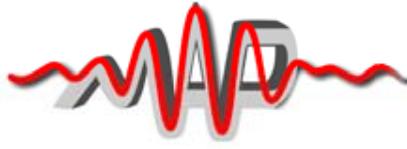
Snowplow LWFA  
of ions injected by RPA

$$\varepsilon_i = (1/6) a_0^2 (n_c/n_e) mc^2$$

0.5TeV over  
dephasing length of 1cm

Zheng et al., 2010

# TeV Ion Energy Scalings(**RPA x LWFA** )

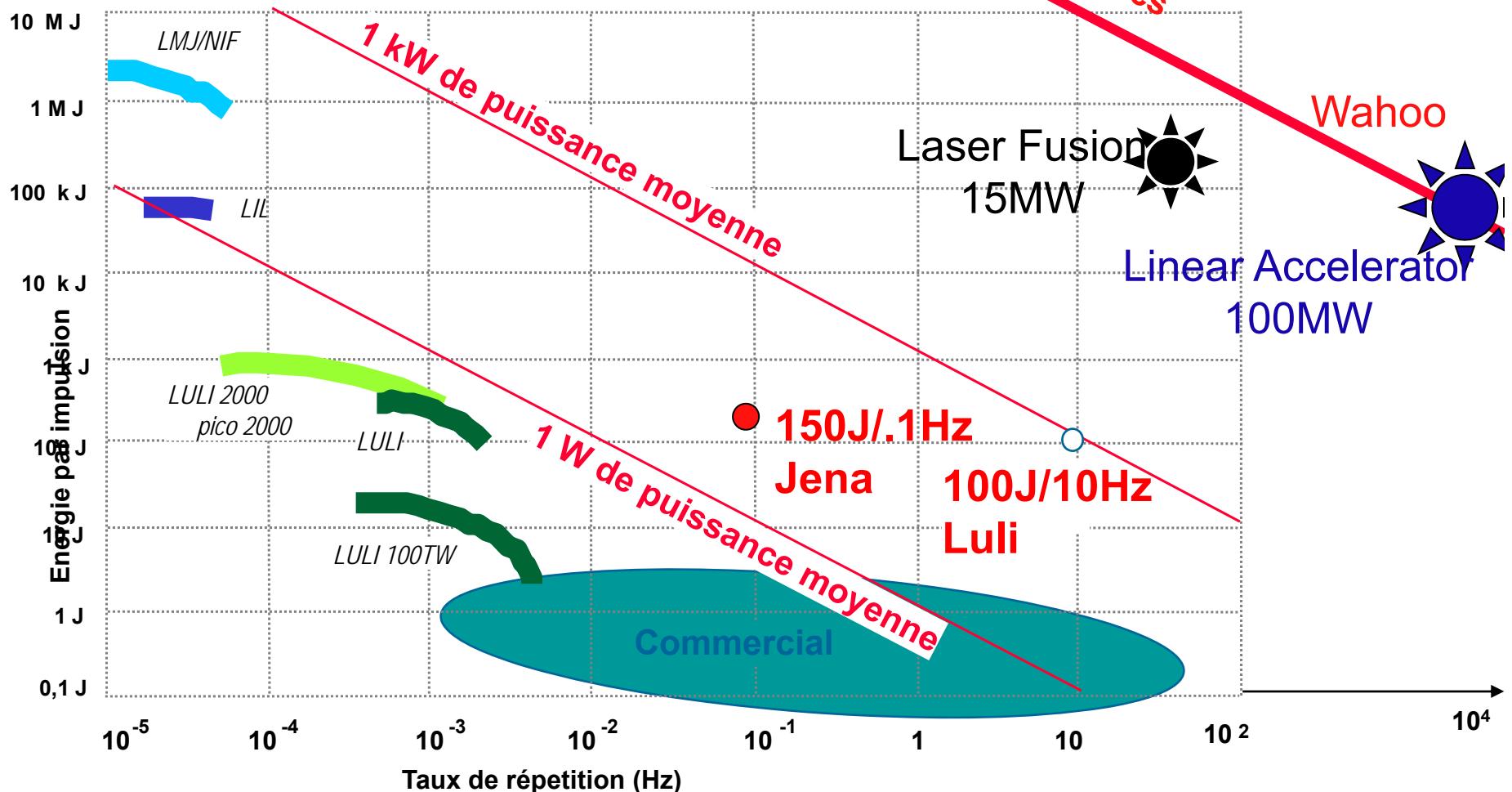


TeV over cm @  $10^{23}\text{W/cm}^2$

(Zheng et al. 2010)

# Etat de l'Art 2005 HEEAUP 2005

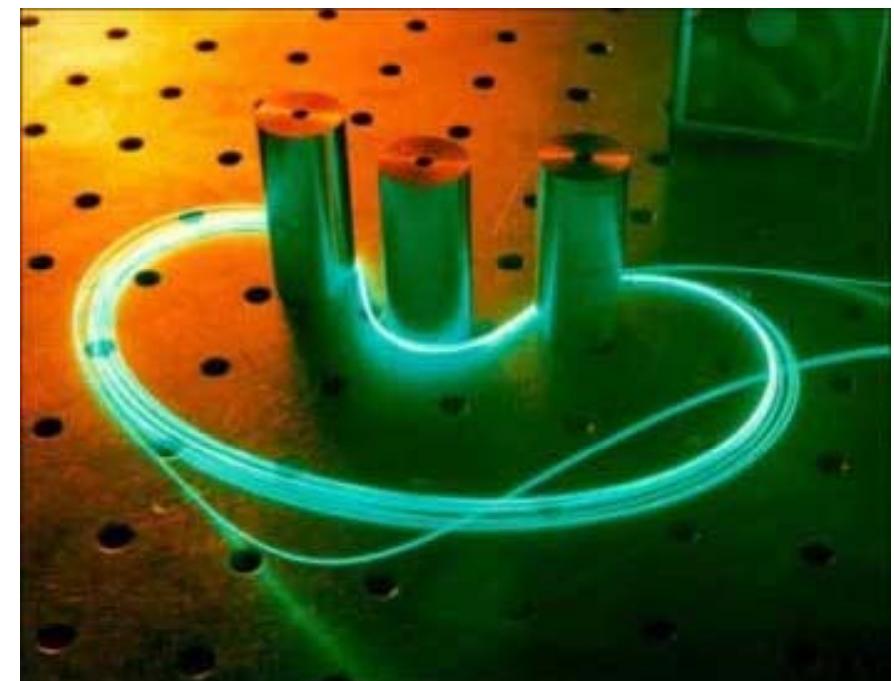
(Mourou, 2005)





# Fiber vs. Bulk Lasers

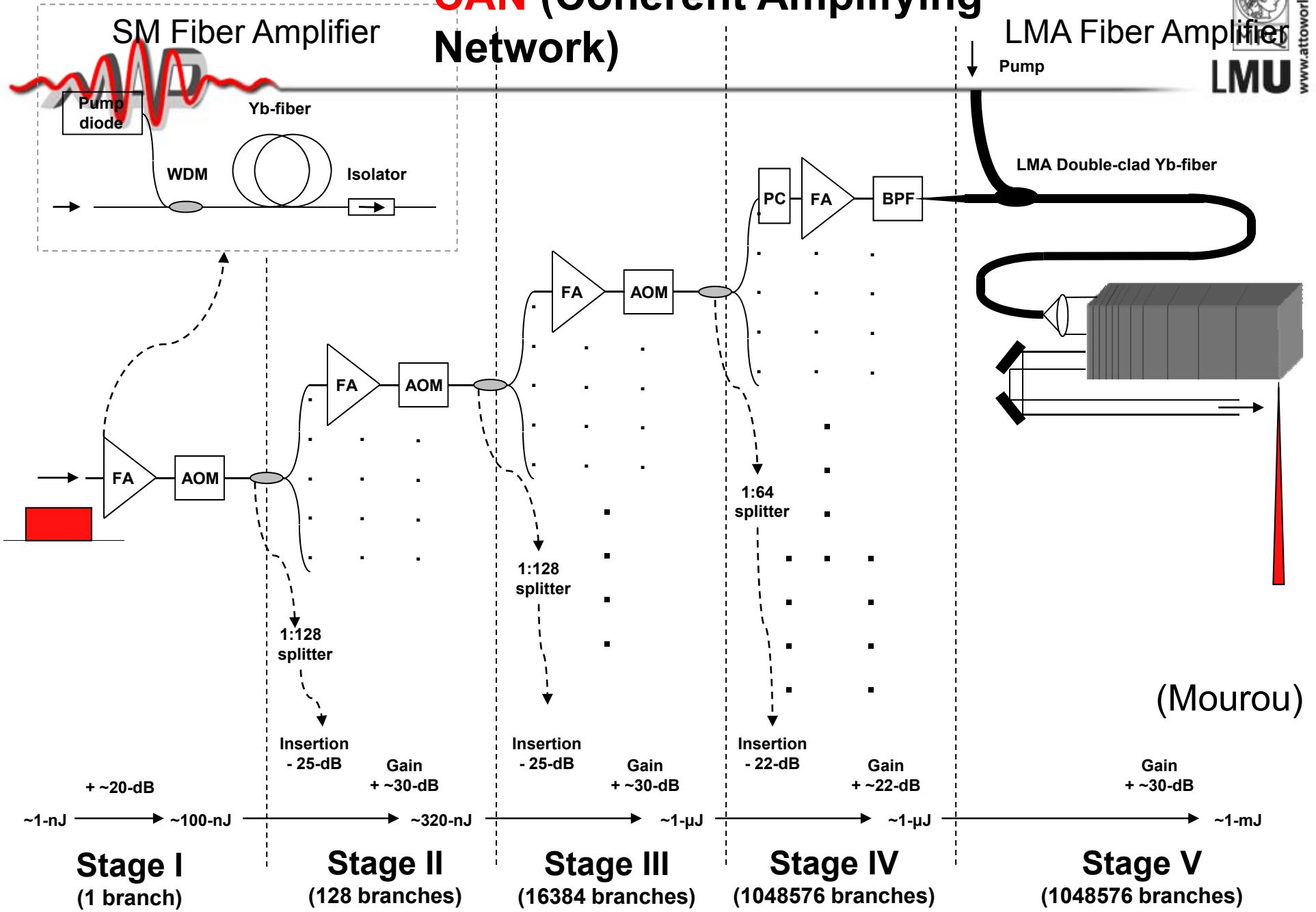
- High Gain *fiber* amplifiers allow ~ 50% total plug-to-optical output efficiency
- Single mode *fiber* amplifier reached multi-kW optical power.
- large bandwidth (100fs)
- immune against thermo-optical problems
- excellent beam quality
- efficient, diode-pumped operation
- high single pass gain
- can be mass-produced at low cost



(G. Mourou)

→ ICAN Consortium formed (Nov. 25, 2010)

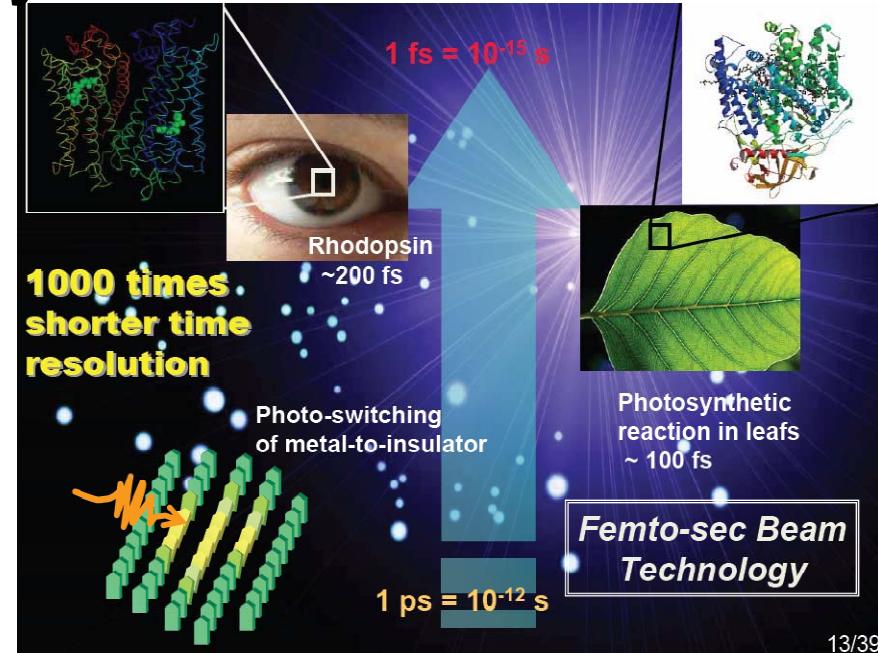
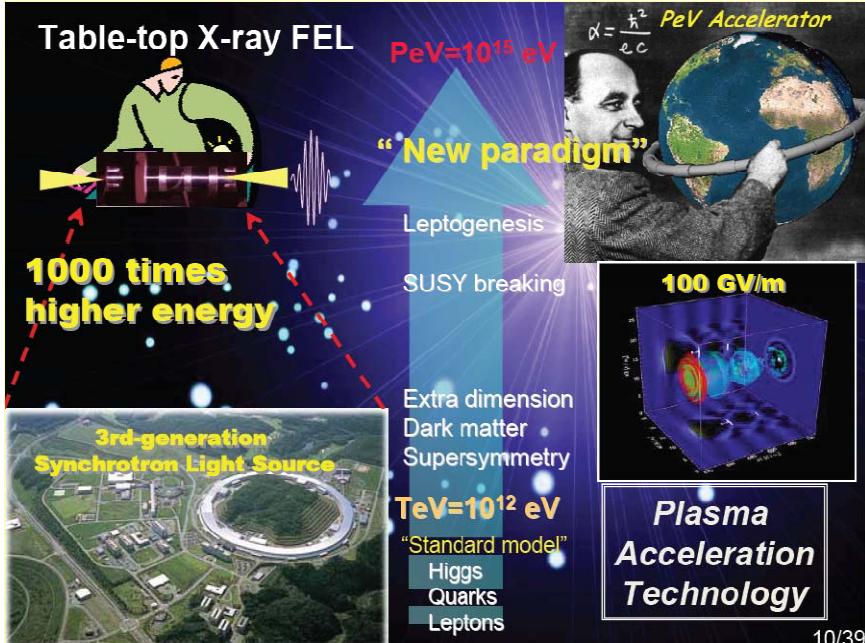
# CAN (Coherent Amplifying Network)



# Challenge Posed by DG Suzuki



Frontier science driven by advanced accelerator



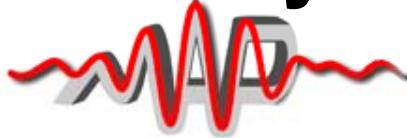
compact, ultrastrong a

*Can we meet the challenge?*

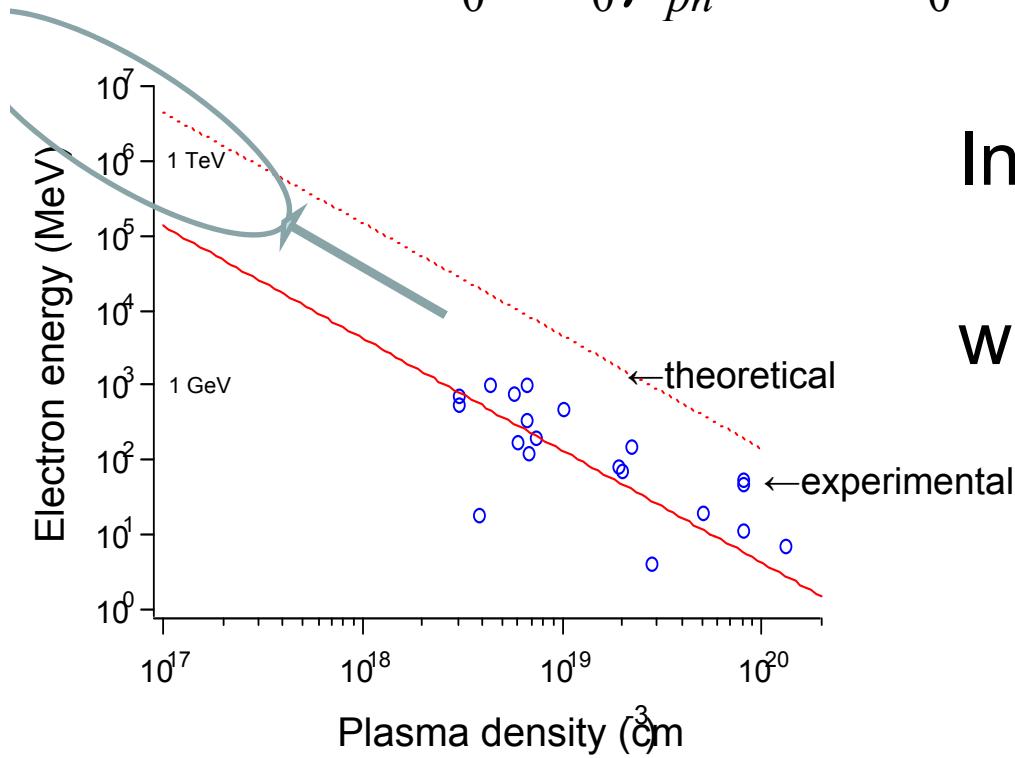
atto-, zeptosecond

A. Suzuki @KEK(2008)

# Theory of wakefield toward extreme energy



$$\Delta E \approx 2m_0c^2a_0^2\gamma_{ph}^2 = 2m_0c^2a_0^2\left(\frac{n_{cr}}{n_e}\right), \text{ (when 1D theory applies)}$$



$$L_d = \frac{2}{\pi} \lambda_p a_0^2 \left( \frac{n_{cr}}{n_e} \right),$$

dephasing length

$$L_p = \frac{1}{3\pi} \lambda_p a_0 \left( \frac{n_{cr}}{n_e} \right),$$

pump depletion length

In order to avoid wavebreak,

$$a_0 < \gamma_{ph}^{1/2},$$

where

$$\gamma_{ph} = (n_{cr}/n_e)^{1/2}$$

Adopt:  
**NIF laser (3MJ)**  
→ 0.7PeV  
(with Kando, Teshima)

# $\gamma$ -ray signal from primordial GRB

LETTERS

NATURE

(Abdo, et al, 2009)

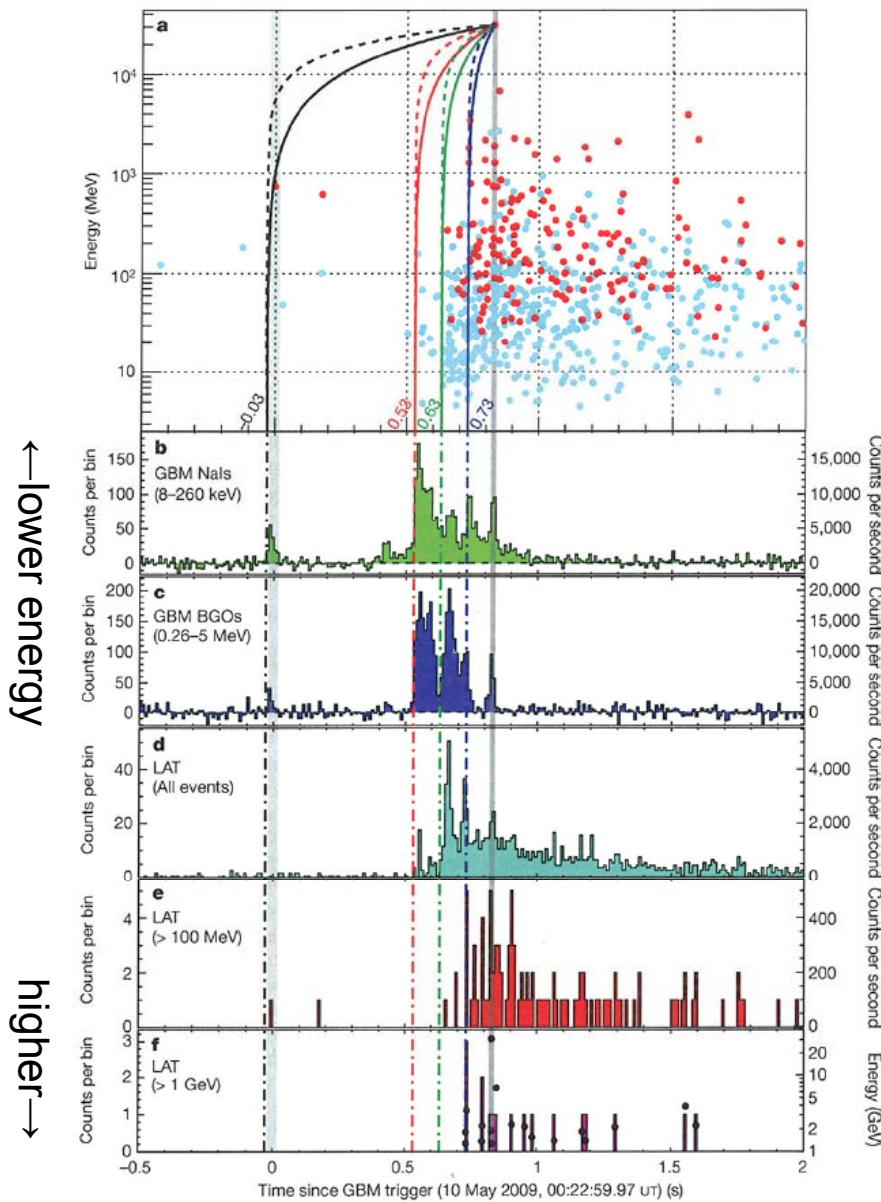
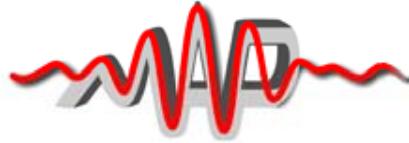


Figure 1 | Light curves of GRB 090510 at different energies. a, Energy lowest to highest energies. f also overlays energy versus arrival time for each

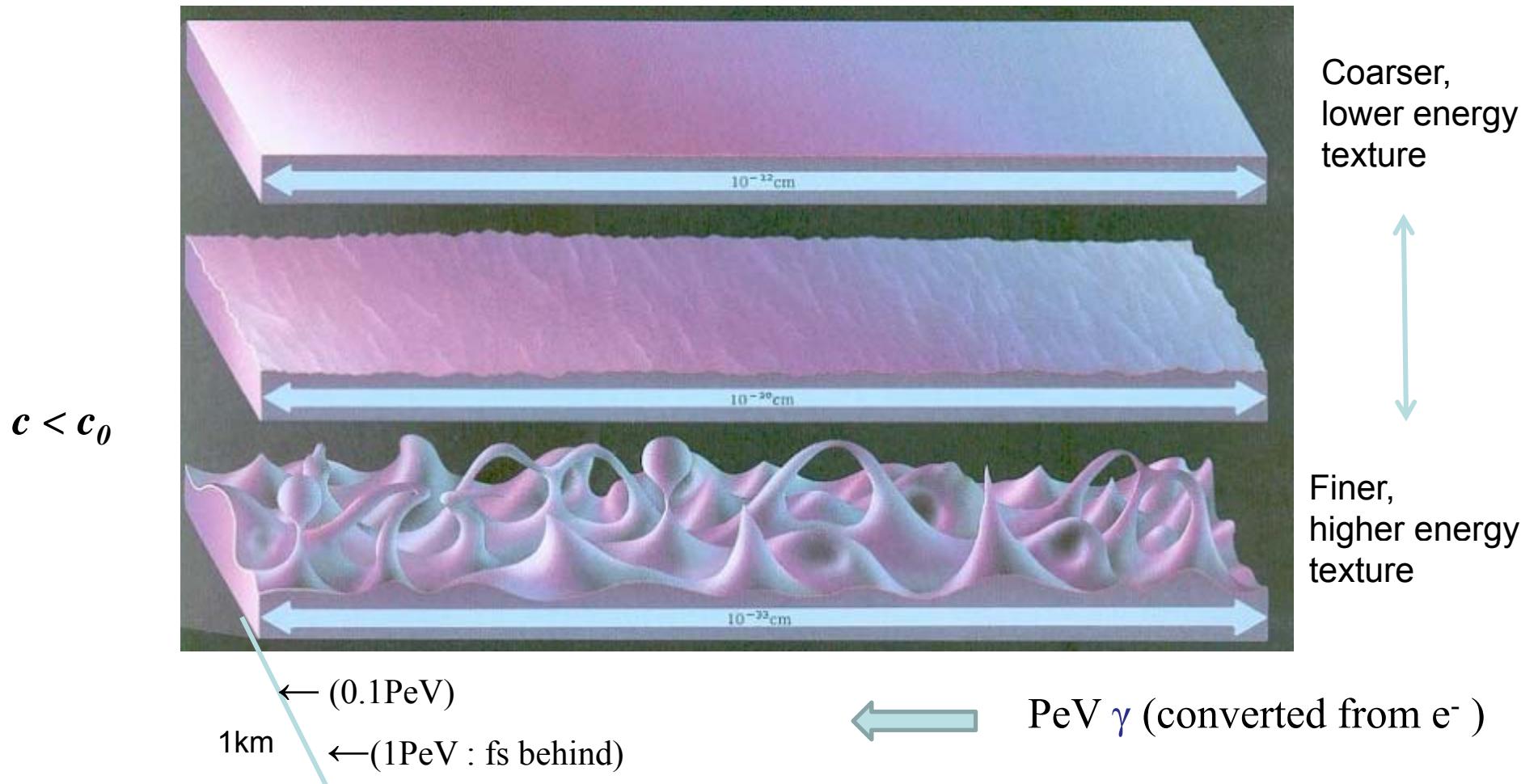
*Energy-dependent photon speed ?*  
**Observation of primordial Gamma Ray Bursts (GRB)**  
**(limit is pushed up close to Planck mass)**

**Lab PeV  $\gamma$  (from e-) can explore this with control**

# Feel vacuum texture: PeV energy $\gamma$

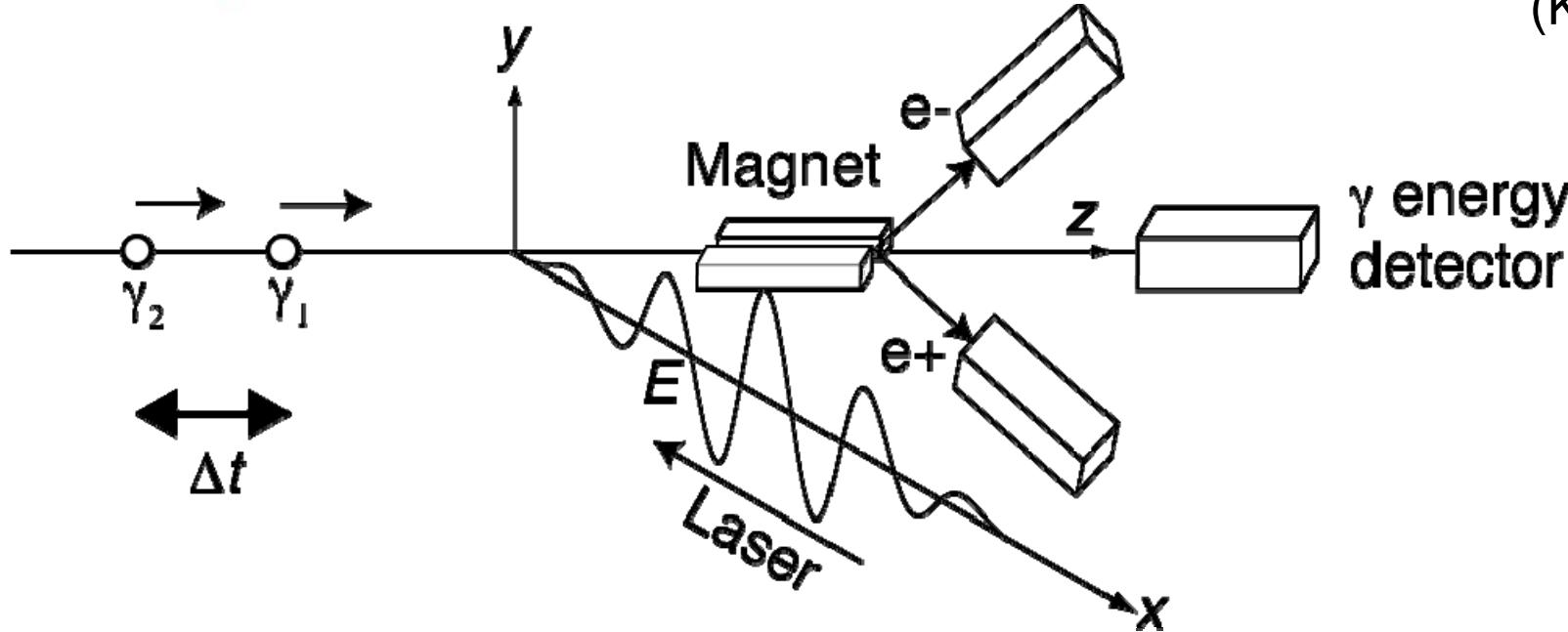


Laser acceleration → controlled laboratory test to see quantum gravity texture on photon propagation (Special Theory of Relativity:  $c_0$ )



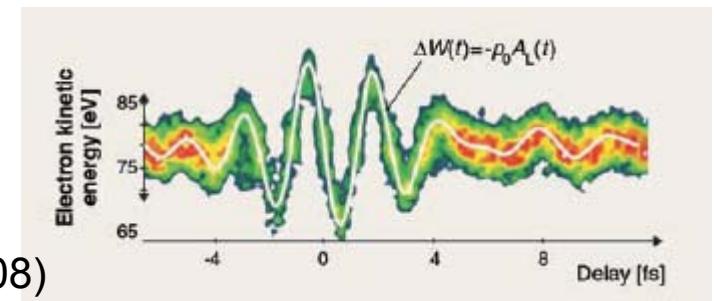
# Attosecond Resolution of PeV $\gamma$ Arrivals

(Kando, 2010)



High energy  $\gamma$ - induced Schwinger breakdown (Narozhny, 1968)  
CEP phase sensitive electron-positron acceleration  
Attosecond electron streaking  
 $\gamma$ - energy tagging possible

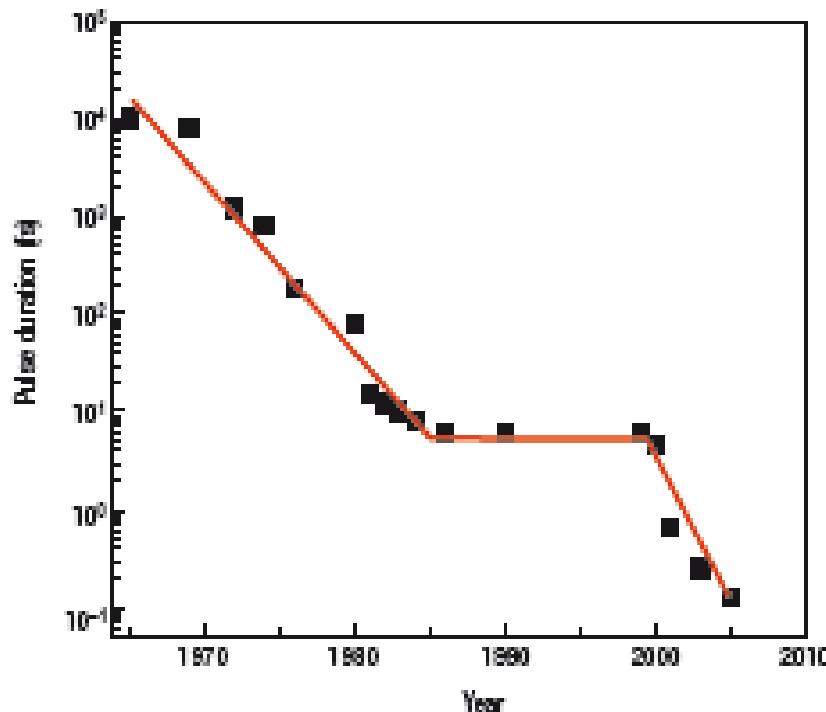
Goulielmakis(2008)



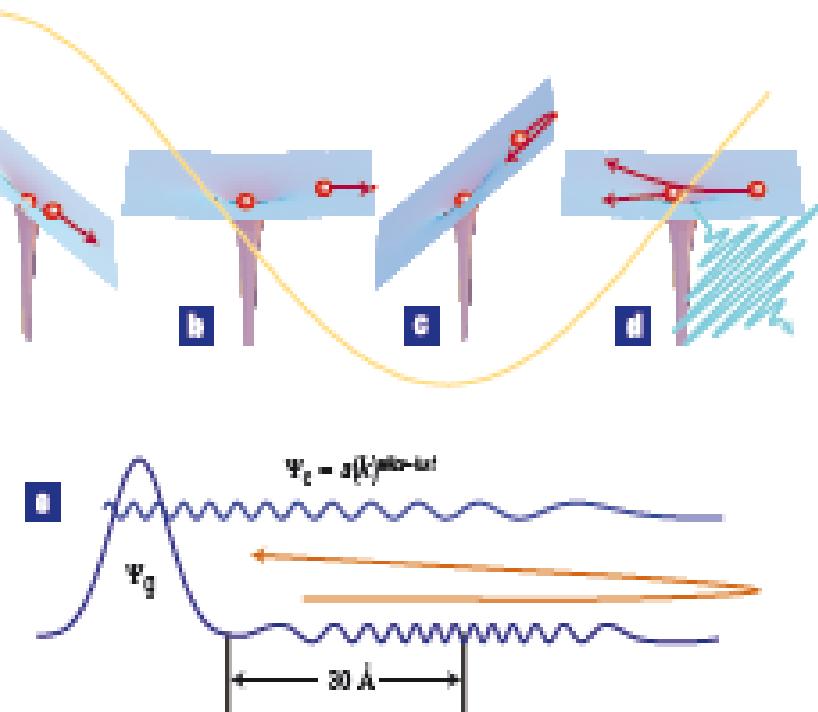


# Pulse Progress from fs to as

Corkum and Krausz (2007)



**Figure 1** Shorter and shorter. The minimum duration of laser pulses fall continually from the discovery of mode-locking in 1964 until 1986 when 6-fs pulses were generated. Each advance in technology opened new fields of science for measurement. Each advance in science strengthened the motivation for making even shorter laser pulses. However, at 6 fs (three periods of light), a radically different technology was needed. Its development took 15 years. Now attosecond technology is providing radically new tools for science and is yet again opening new fields for



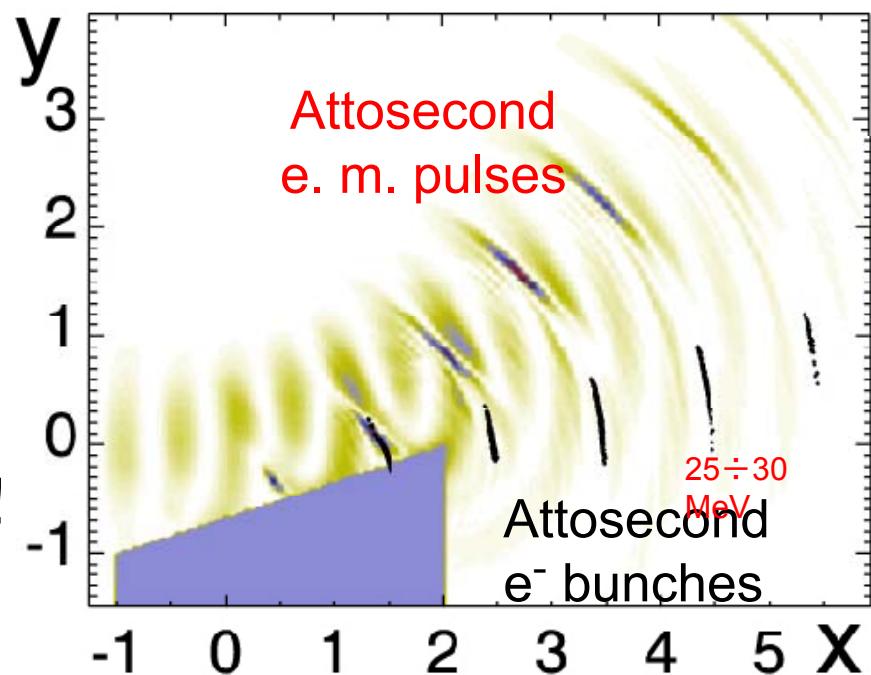
**Figure 2** Creating an attosecond pulse. a-d, An intense femtosecond near-infrared or visible (henceforth optical) pulse (shown in yellow) extracts an electron wavepacket from an atom or molecule. For ionization in such a strong field (a), Newton's equations of motion give a relatively good description of the response of the electron. Initially, the electron is pulled away from the atom (a, b), but after the field reverses, the electron is driven back (c) where it can 'recollide' during a small fraction of the laser oscillation cycle (d). The parent ion sees an attosecond electron pulse. This

# Electron ejection is synchronized with attosecond pulse generation

Escaped relativistic electrons

- compress the reflected radiation into attosecond pulses and
- inherit a peaked density distribution.
- Complete modulation of e.m. field occurs. This is relativistic microelectronics!

Efficiency of attosecond phenomena: ~15% converted to **attosecond pulses**, ~15% to electron bunches.



Naumova *et al.*, Phys. Rev. Lett. (2004)

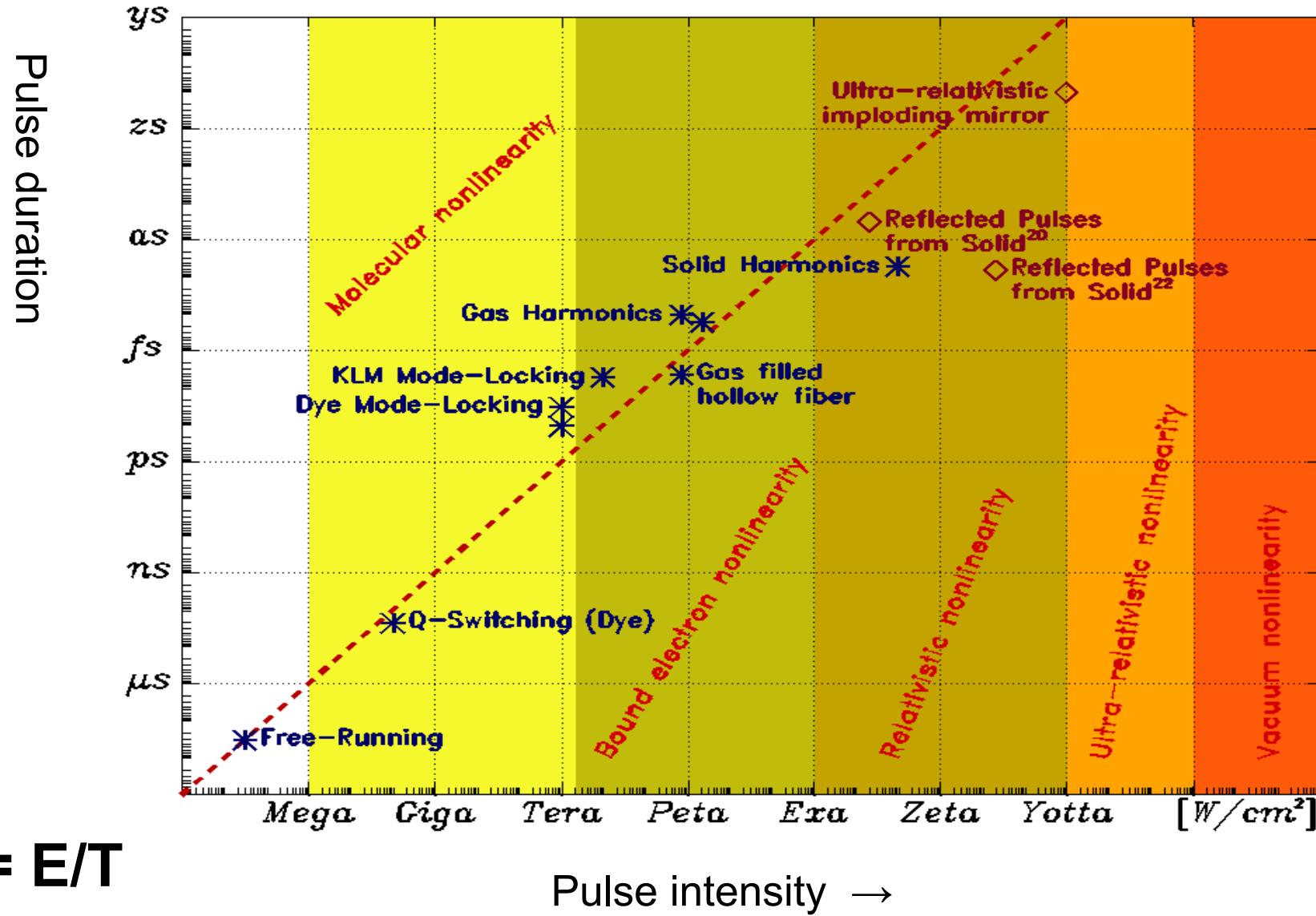
$a=10, 15\text{fs}, f/1,$   
 $n=25n_{cr}$

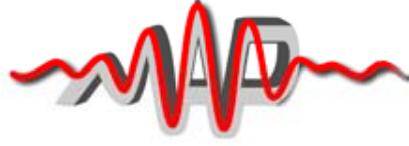
# The Mourou Conjecture



(← physics: “Matter is nonlinear”)

“The rigider nonlinearity, the more intense to manipulate it” )





# Yoctosecond coherent $\gamma$ -beam

The Mourou-Tajima Conjecture tells us

that ys coherent **gamma pulse**

may be generated by Yottawatt/cm<sup>2</sup>

Go beyond 10KeV coherent **X-rays**: solid is not dense enough

How can one make an Einstein mirror

that can backscatter **gamma rays** (several 100MeV)?

Suggest: the imploded ultra-relativistic shell

to increase the density by  $10 \times 10 \times 10$ .

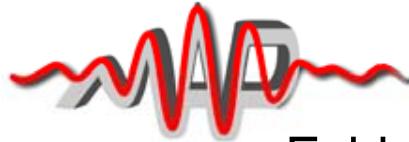
The relativistic flying object with  $\gamma = 100$ . The lab frame density enhanced by  $10^5$

Takes **large** energy ~MJ, Yottawatt/cm<sup>2</sup>

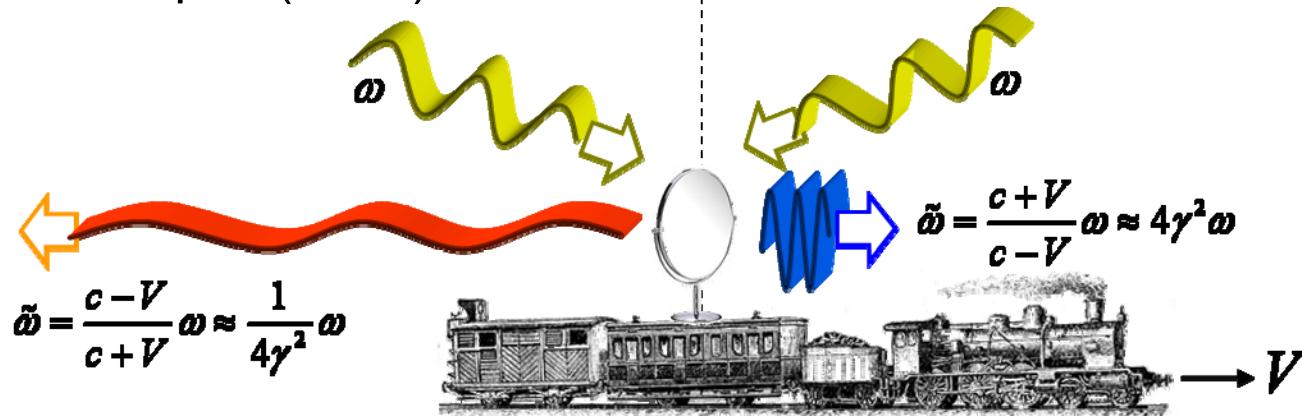
→ yielding 100ys coherent **gamma rays**

*Can we streak a nucleus with this gamma beam? :  
**Photonuclear Physics.** (Any other new physics?)*

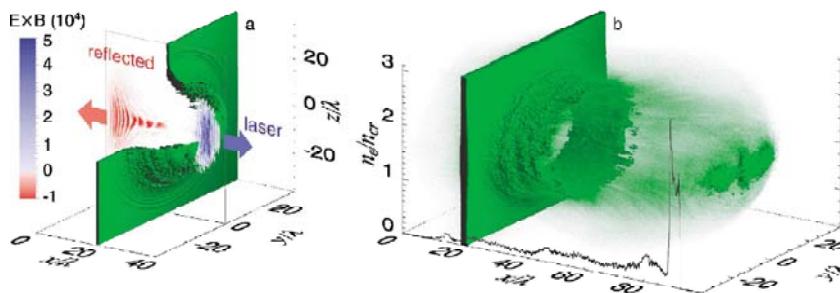
# Relativistic flying mirror and shorter pulses



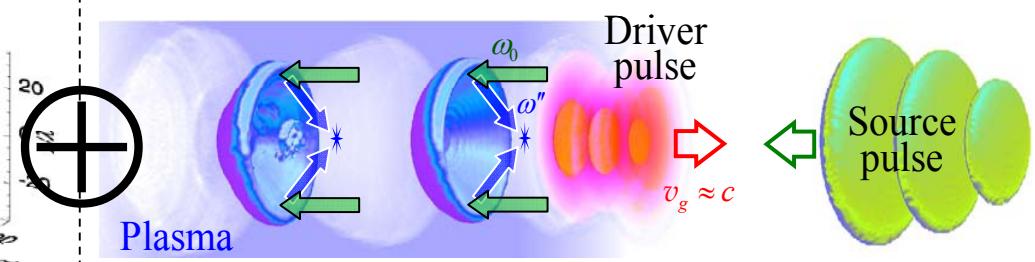
Esirkepov (2010)



## Laser Piston



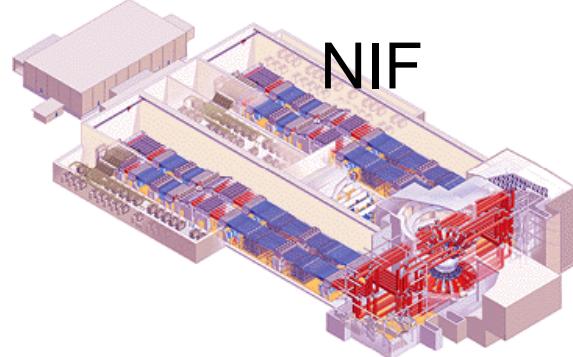
## Flying Mirror



24

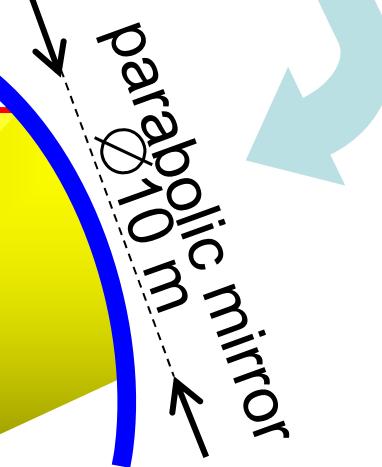
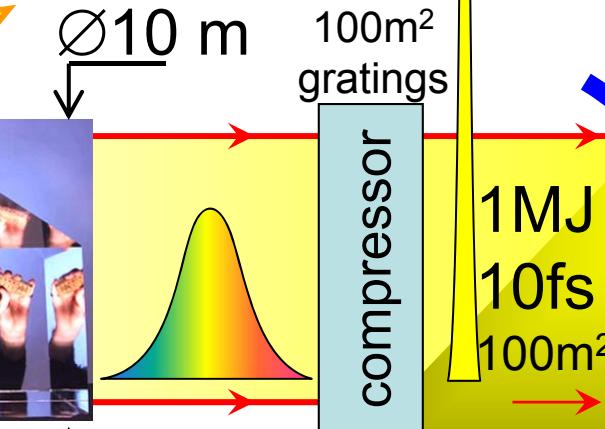
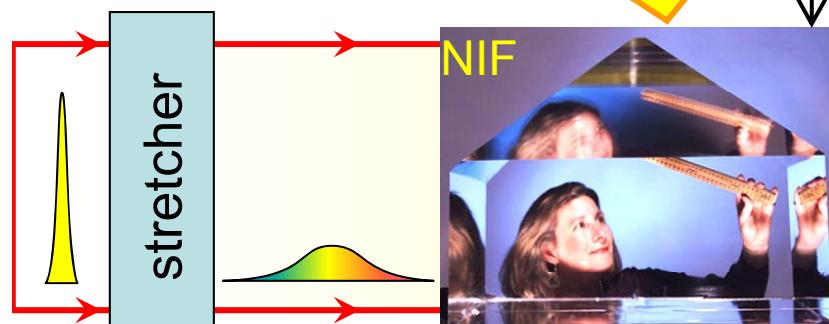
Tajima/Mourou/Moses(2010): use NIF ---ultra-relativistic imploding mirror → ys!

# MJ laser → Zettawatt → ultrashort pulse



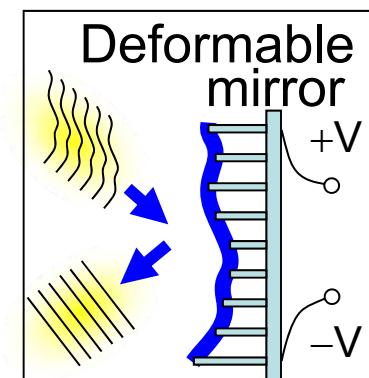
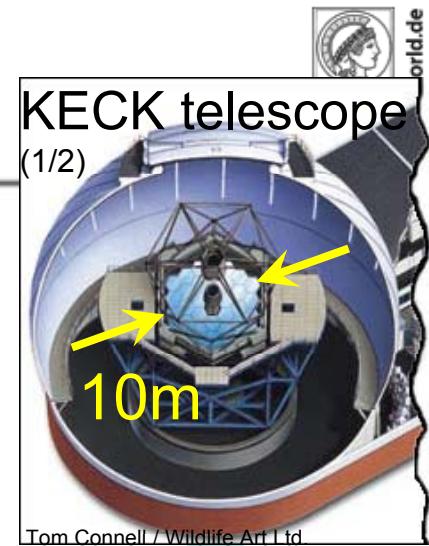
Tajima,Mourou (PR, 2002)

5MJ @ 10ns  
530nm



$10^{28} \text{ W/cm}^2 !$

∅1 micron



# Nuclear Wake?

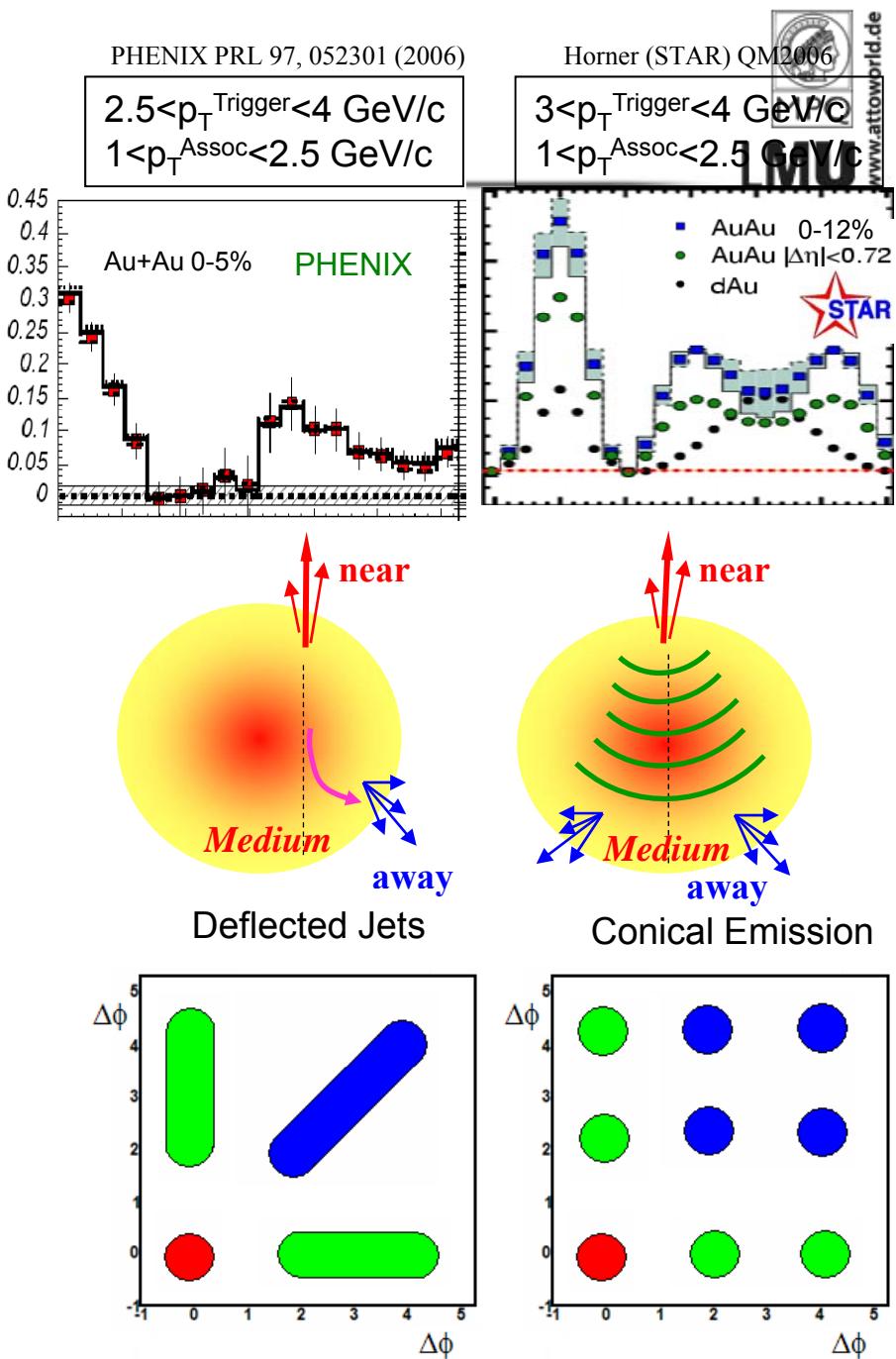


- Broadened and maybe double humped structure on the away-side in 2-particle correlations.
- Could be caused by:
  - Large angle gluon radiation (Vitev and Polsa and Salgado).
  - Deflected jets, due to flow (Armesto, Salgado and Wiedemann) and/or path length dependent energy loss (Chiu and Hwa).
  - Hydrodynamic conical flow from mach cone shock-waves (Stoecker, Casalderrey-Solanda, Shuryak and Teaney, Renk, Ruppert and Muller).
  - Cerenkov gluon radiation (Dremin, Koch).
- Jet quenching: collective deceleration by wakefield?
  - LWFA method, or Maldacena method?

8 Aug 2007

ISMD

Glynnis Ulery





High energy physics ← High field science, high intensity

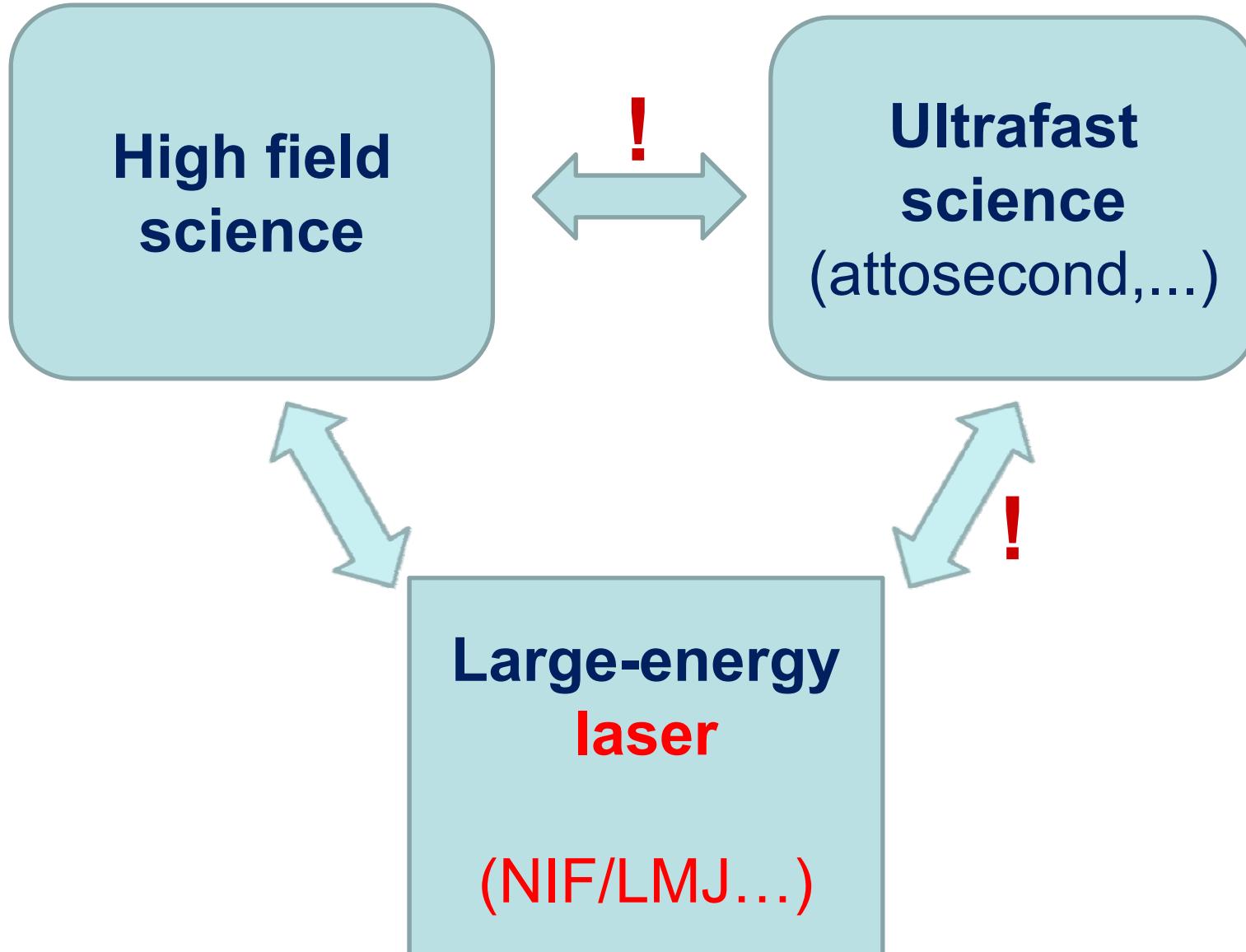
High field  
science



High energy  
Physics  
(fundamental  
physics)



**Ultrafast science** ← High field science, Large-energy **laser**





# Self-focusing in air to vacuum

Critical power for self-focusing in matter /plasma / vacuum:

$\chi_3$  nonlinearity

$$P_{cr} = \lambda^2/(2\pi n_0 n_2) \sim \text{GW}$$

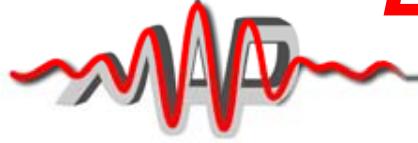
relativistic plasma nonlinearity

$$P_{cr} = mc^5/e^2(\omega/\omega_p)^2 \sim 17 (\omega/\omega_p)^2 \text{ GW}$$

vacuum nonlinearity

$$P_{cr} = (90/28) c E_S^2 \lambda^2 / \alpha \sim 10^{15} (\lambda/\lambda_{I_\mu})^2 \text{ GW}$$

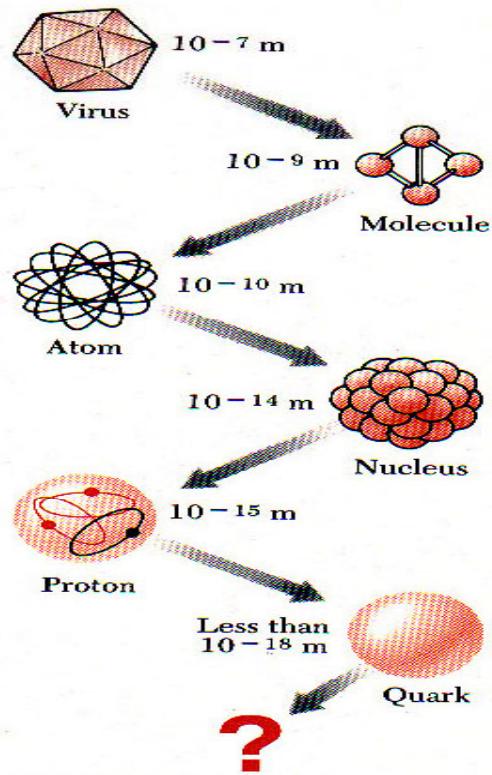
e.g. X-ray of 10keV,  $P_{cr} \sim 10\text{PW}$



# 'ELI Long-term Ambition' =

## *Studying the Atomic Structure to the Vacuum Structure*

(Mourou)



Vacuum structure

Does the atomic world repeat itself in vacuum?

Keldysh field

Keldysh parameter

$$\gamma_K = \omega \sqrt{2m\Phi}/eE,$$

Schwinger intensity / Keldysh intensity  $= \alpha^{-6} \sim 10^{14}$

Vacuum self-focusing /  $\chi_3$  self-focusing power  $\sim \alpha^{-6} \sim 10^{15}$

Vacuum parameter

$$\gamma_V = 1/a_0.$$

$$\gamma_V \geq 1.$$

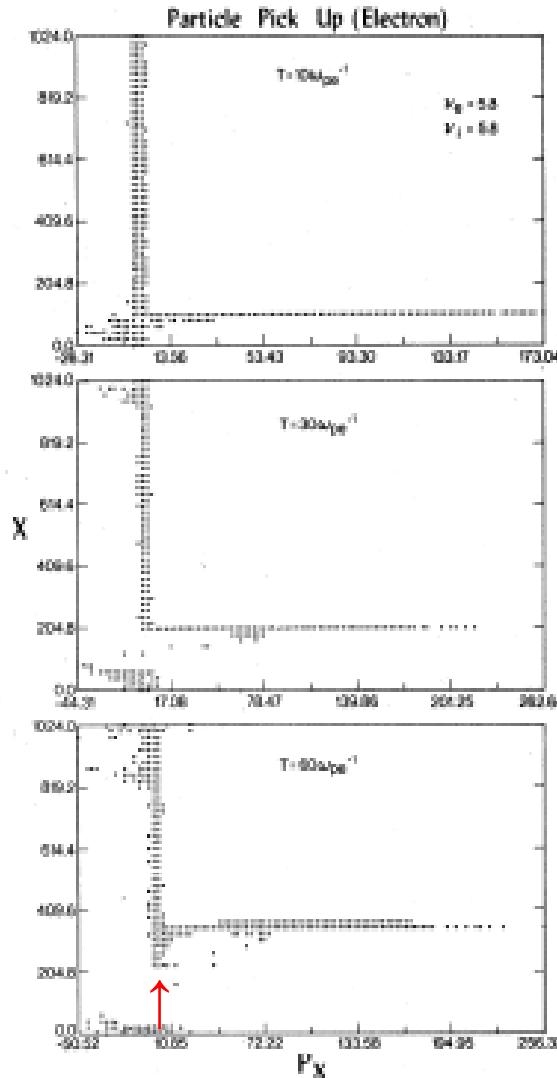
= Path toward superSchwinger

Schwinger field

# Self-focusing and laser acceleration in vacuum

3

## ULTRARELATIVISTIC ELECTROMA



Ashour-abdalla, et al. (1981)

on the piston, and pressure balance yields

$$2\langle n \rangle_w \langle p_x \rangle_w v_x = \frac{1}{2} [E_w^2 + B_w^2]/8\pi, \quad (10)$$

where  $v_x$  is the piston (wave-front) velocity, the subscript  $w$  indicates that the fields, momentum, and density are to be evaluated in the moving frame of the interface, and  $\langle n \rangle$  and  $\langle p_x \rangle$  are, respectively, the average density and momentum at the leading edge of the pulse (over a length  $c/\omega_{pe}$  in the laboratory frame). Lorentz transformations from the moving to the lab frame yield

$$2v_p(\langle n \rangle_L/\gamma_x)(\langle p_x \rangle_L/2\gamma_p) = \frac{1}{2} E_L^2/\gamma_x^2, \quad (11)$$

where the subscript  $L$  denotes laboratory frame variables and the relativistic factor  $\gamma_x = [1 - (v_x^2/c^2)]^{-1/2}$ .

Our simulations show that  $\langle p_x \rangle_L \approx \frac{1}{2} p_x^m$  and  $\langle n \rangle_L \approx \frac{1}{2} n_{max}$ , with  $n_{max} \approx 20n_0$  (twenty times the original density) at saturation, so that

$$p_x^m = 2(c/v_p)[\omega^2/\omega_{pe}^2(n_0)][\omega_{pe}^2(n_0)/\omega_{pe}^2(n_{max})]\nu^2 mc. \quad (12)$$

With  $\omega^2/\omega_{pe}^2(n_0) \approx 10$  and  $\omega_{pe}^2(n_0)/\omega_{pe}^2(n_{max}) \approx \frac{1}{20}$ , we obtain

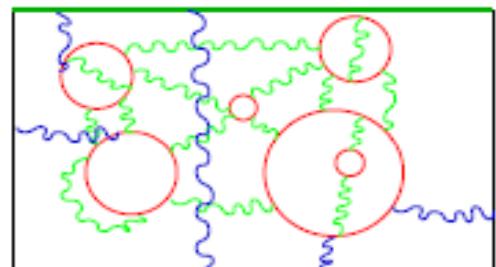
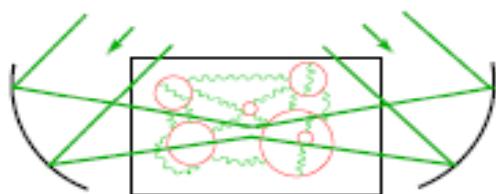
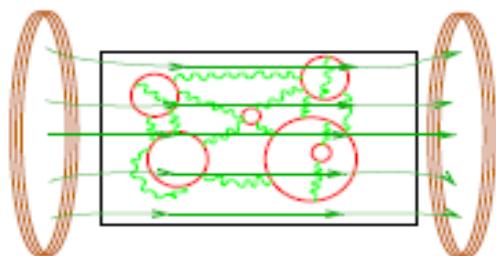
$$p_x^m \approx (c/v_p)\nu^2 mc \sim \nu^2 mc, \quad (13)$$

→ **Laser** acceleration in vaccum by self-focused **X-ray** crossed with **laser** ?

Can we repeat LWFA and plasma physics in vacuum?  
**Laser** vacuum acceleration with ‘snowplow’

# Why quantum vacuum physics?

## Vacuum nonlinearities



- Heisenberg-Euler/Casimir in mathematical physics
  - QFT in strong fields or with boundaries
  - functional determinants
- applied quantum vacuum physics
  - quantum fluctuations as a building block
  - dispersive forces in micro/nano machinery [DEKIEVET @ THISWORKSHOP]
- fundamental effect of QFT
  - ( $\sim$  Lamb shift,  $g - 2$ , ...)
- fundamental physics
  - search for new physics
  - new particles or forces

H. Gies (2008)

# Light Propagation in a $B$ field.

- ▷ quantum Maxwell equation for a "light probe"  $f^{\mu\nu}$

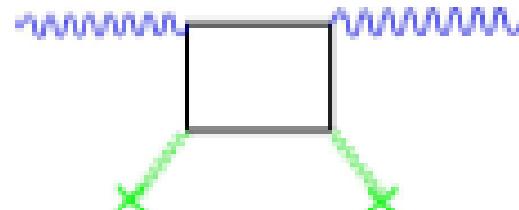
$$0 = \partial_\mu f^{\mu\nu} - \frac{8}{45} \frac{\alpha^2}{m^4} F_{\alpha\beta} F^{\mu\nu} \partial_\mu f^{\alpha\beta} - \frac{14}{45} \frac{\alpha^2}{m^4} \tilde{F}_{\alpha\beta} \tilde{F}^{\mu\nu} \partial_\mu f^{\alpha\beta}$$

↑ vacuum nonlinearity ↑

## Phase and group velocity

$$v_{||} \approx 1 - \frac{14}{45} \frac{\alpha^2}{m^4} B^2 \sin^2 \theta_B$$

$$v_{\perp} \approx 1 - \frac{8}{45} \frac{\alpha^2}{m^4} B^2 \sin^2 \theta_B$$



(Tol'sai)

(BAIER, BREITENLOHNER '67; NARODZINSKI)

(BIALYNICKA-BIRULA, BIAŁYNICKI-BIRULA '74)

(MOUROU '71)

⇒ magnetized quantum vacuum induces birefringence

[DiPiazza @ ThisWorkshop]

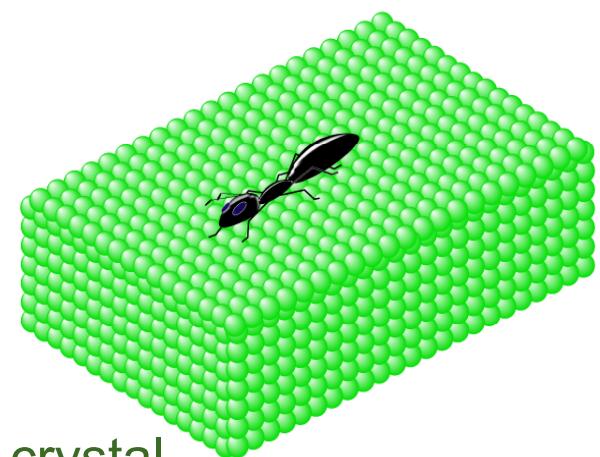
- ▷ detection schemes: PVLAS, BMV, Q&A, OSQAR, TR18-B7

H. Gies (2008)

Self-focusing:  $P_{cr} = (90/28) c E_S^2 \lambda^2 / \alpha \rightarrow P_{cr} \approx 10^{24} (\omega_l \text{ (at } 1\mu\text{)} / \omega)^2 \text{ W}$   
 (Mourou/Tajima/Bulanov (2006))



# What is vacuum?

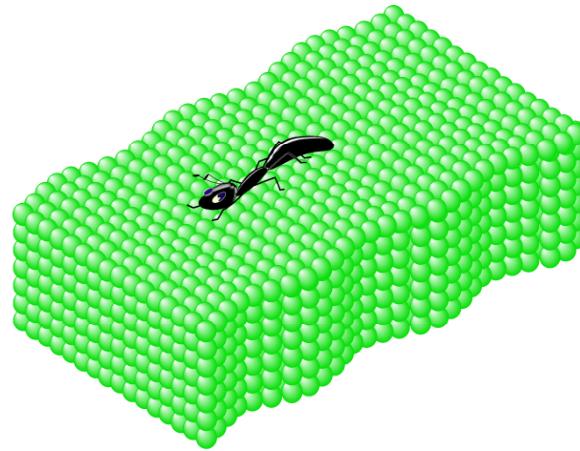


crystal

An observer (bug) in crystal looks at **vacuum**

vacuum

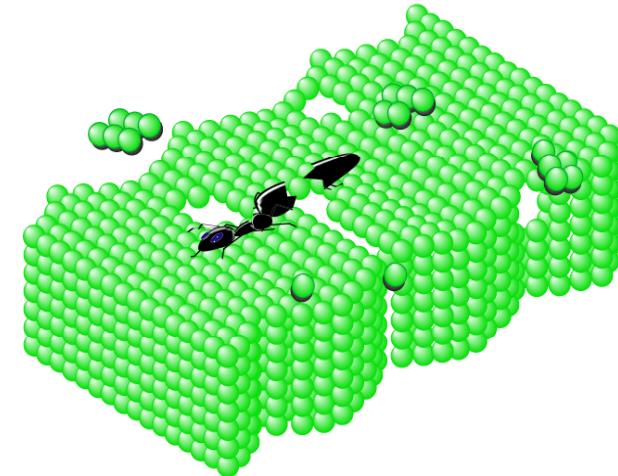
「真(true)空(nothing)」



Phonon : excitation of **vacuum**



Photon : distortion of **vacuum**  
「色=(即是)空」

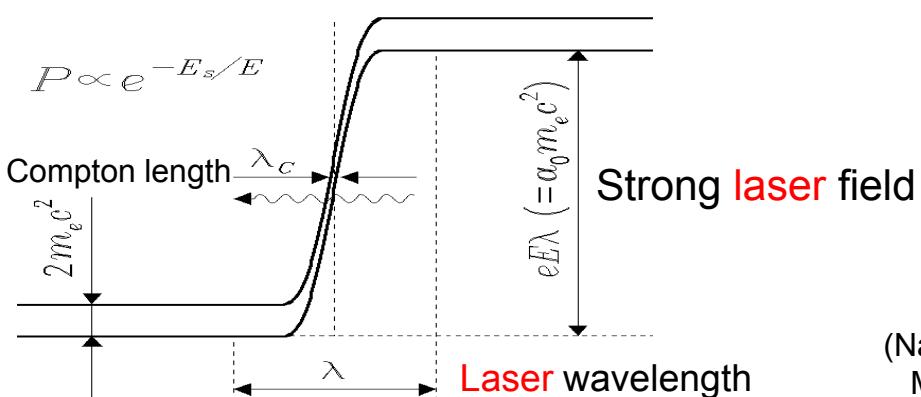


Strong field breaks **vacuum**



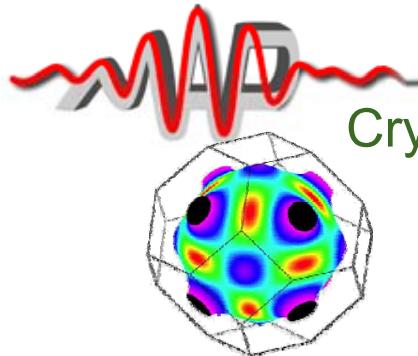
vacuum produces **e+e- pair**  
「空=(即是)色」

QED **vacuum breakdown**

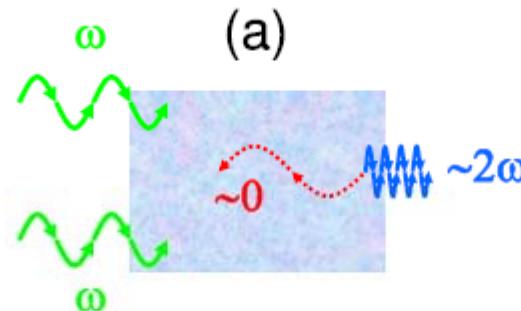


(Naumova  
Mourou)

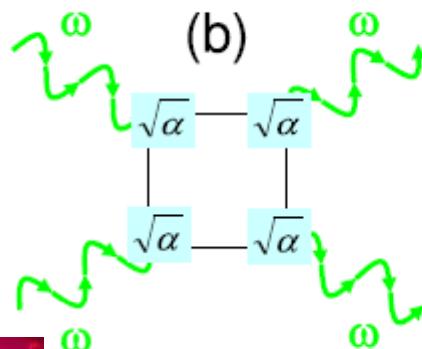
# Intense laser probes matter /vacuum nonlinearity



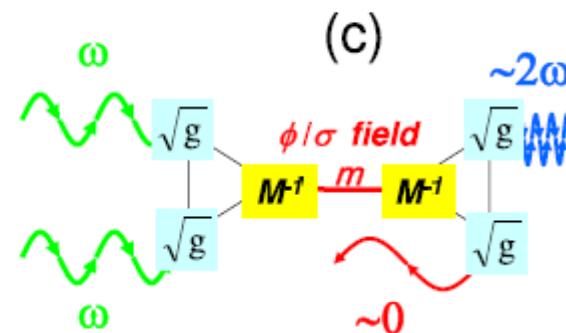
Crystal nonlinearity →  
second harmonic generation (Franken et al)



Learn from Nonlinear Optics of matter for vacuum:



QED nonlinearity



Vacuum nonlinearity by light- mass  
field (dark energy, axion,...)  
→ second harmonic

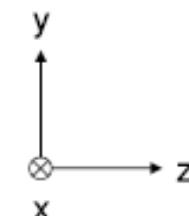
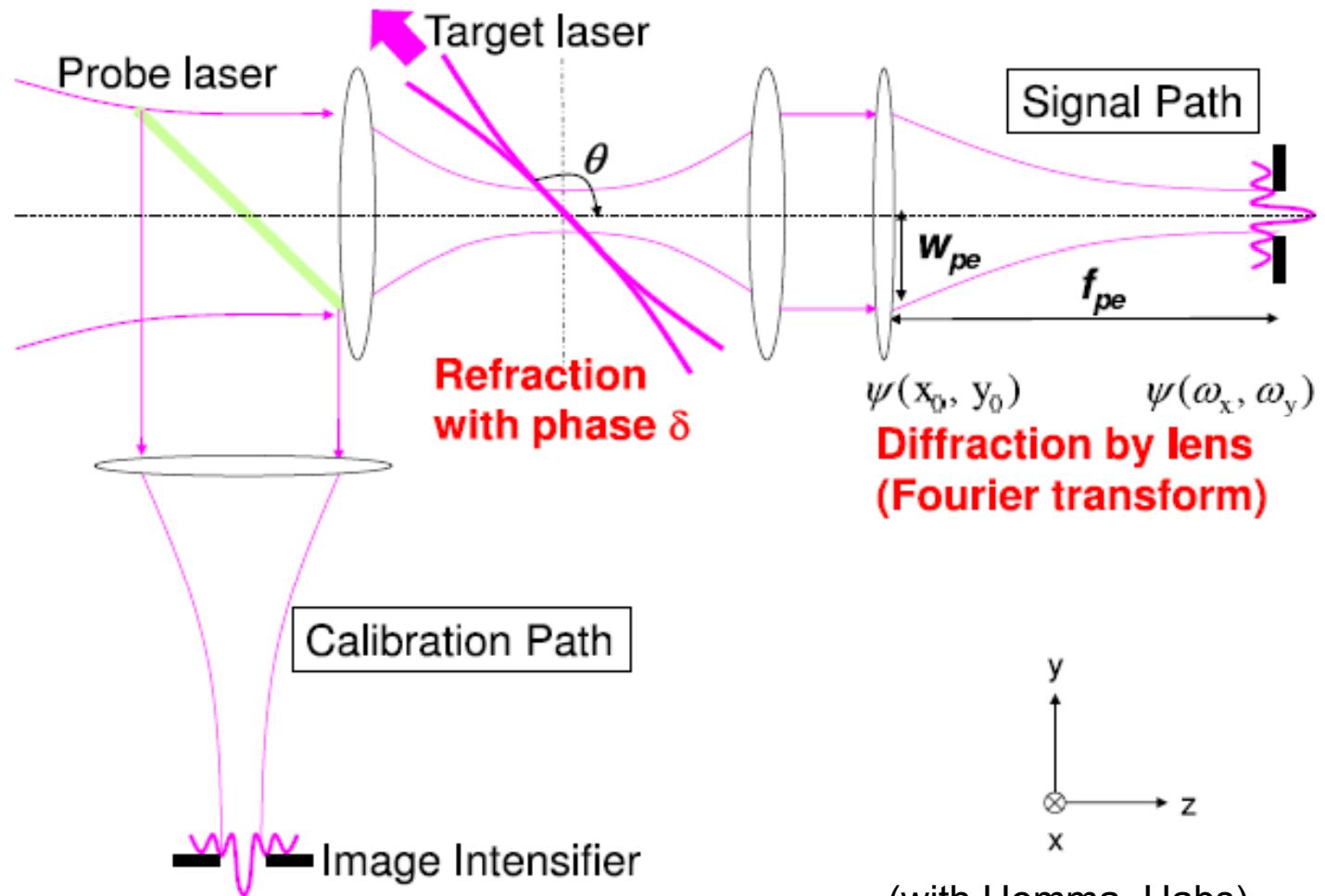


# QED vacuum probe by intense laser



Heisenberg-Euler Langrangian: tiny nonlinearity, never observed  
→ intense **laser** needed; sensitive probe, avoid blinding **laser**

**Phase contrast imaging**(refractive index →diffraction, noise reduction)

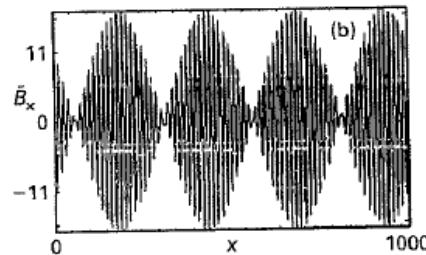
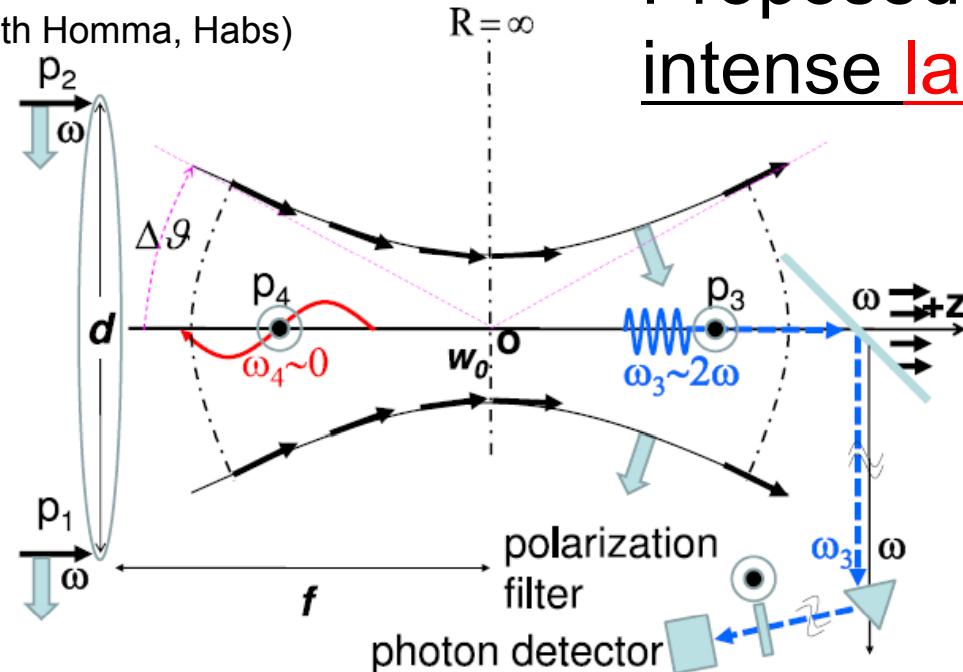


(with Homma, Habs)

# Learning from **laser** parametric scattering: low energy (meV - neV) fields (vacua)



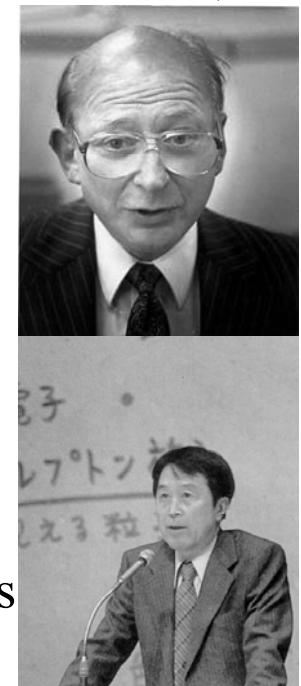
(with Homma, Habs)



cf. Brillouin forward scattering beat / optical parametric excitation = phonon mediating (Nambu-Goldston boson)

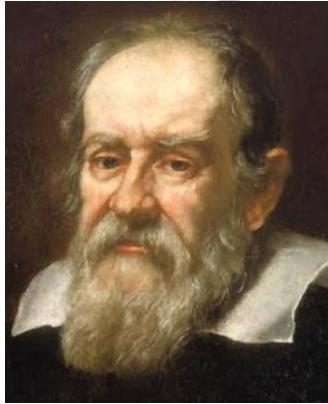
Mass of light fields(dark energy fields, axion-like fields) resonates with specific crossing angle of co-propagating **lasers**

Many orders of magnitude gain in resonant coupling and sensitivity over long interaction:  
Nonlinearity of **vacuum**  
 $\omega + \omega \rightarrow 2\omega$  (SHG a la Franken)

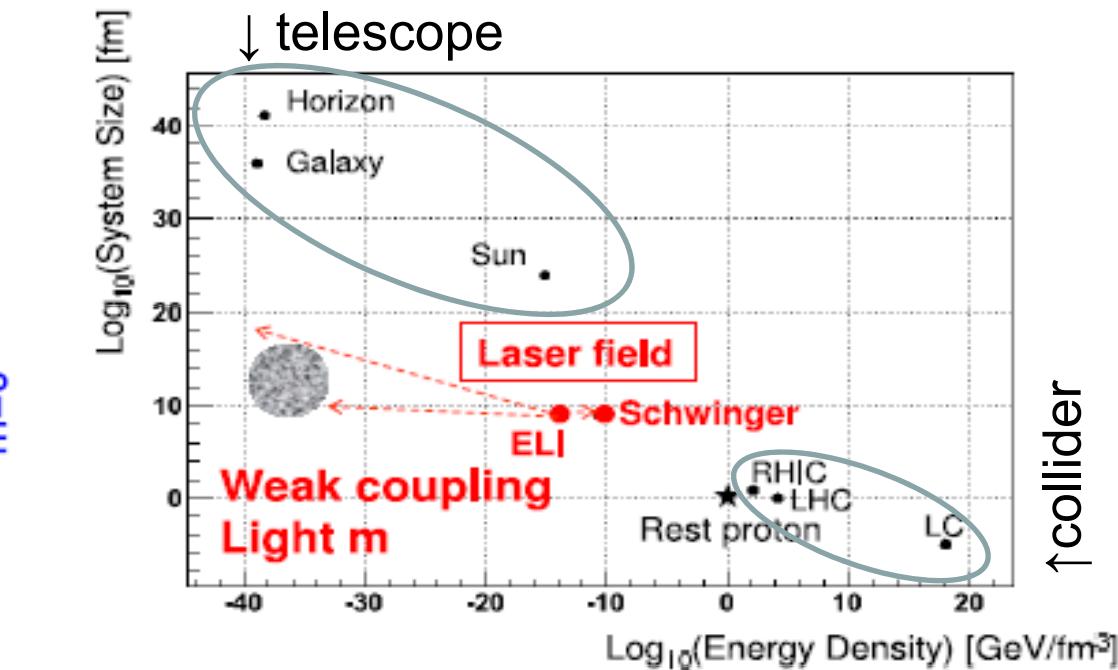




# Scope of High Field Science vs traditional approaches



Weak coupling  
 $m=0$



↑collider



Strong coupling  
Heavy m      High energy  
collider

(with Homma, Habs)



(Mourou, 2010)